

Pre-Born–Oppenheimer Molecular Structure Theory

Edit Mátyus

Institute of Chemistry and
MTA-ELTE Complex Chemical Systems Research Group
Eötvös Loránd University
Budapest, Hungary

Quantum chemistry and quantum mechanics

Born–Oppenheimer separation

Electronic structure theory

- In the external electric field of the atomic nuclei

- Kinetic energy operator:

$$-\sum_{i=1}^{n_e} \frac{1}{2m_e} \Delta_{r_i}$$

- Potential energy:

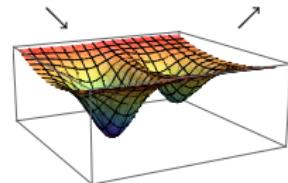
$$\sum_{i=1}^{n_e} \sum_{j>i}^{n_e} \frac{1}{r_{ij}} - \sum_{i=1}^{n_e} \sum_{n=1}^{n_n} \frac{Z_n}{r_{in}}$$

Nuclear motion theory

- Translational and rotational invariance for an isolated molecule

- Kinetic energy operator: curvilinear coordinates, $\hat{T}(\rho)$

- Potential energy: $(3n_n - 6)$ -dimensional representations



potential energy surface (PES)

Quantum chemistry and quantum mechanics

Born–Oppenheimer separation

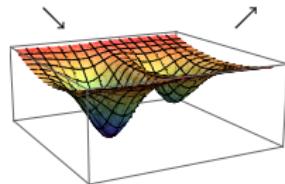
Electronic structure theory

- In the external electric field of the atomic nuclei
- Kinetic energy operator:

$$-\sum_{i=1}^{n_e} \frac{1}{2m_e} \Delta_{r_i}$$

- Potential energy:

$$\sum_{i=1}^{n_e} \sum_{j>i}^{n_e} \frac{1}{r_{ij}} - \sum_{i=1}^{n_e} \sum_{n=1}^{n_n} \frac{Z_n}{r_{in}}$$



potential energy surface (PES)

Nuclear motion theory

- Translational and rotational invariance for an isolated molecule
- Kinetic energy operator: curvilinear coordinates, $\hat{T}(\rho)$
- Potential energy: $(3n_n - 6)$ -dimensional representations

The tremendous success of the usual practice might perhaps be best regarded as a tribute to the insight and ingenuity of the practitioners for inventing an effective variant of quantum theory for chemistry.

B. T. Sutcliffe and R. G. Woolley, J. Chem. Phys. 137, 22A544 (2012).

Quantum chemistry and quantum mechanics

Born–Oppenheimer separation

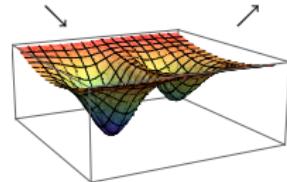
Electronic structure theory

- In the external electric field of the atomic nuclei
- Kinetic energy operator:

$$-\sum_{i=1}^{n_e} \frac{1}{2m_e} \Delta_{r_i}$$

- Potential energy:

$$\sum_{i=1}^{n_e} \sum_{j>i}^{n_e} \frac{1}{r_{ij}} - \sum_{i=1}^{n_e} \sum_{n=1}^{n_n} \frac{Z_n}{r_{in}}$$



potential energy surface (PES)

Nuclear motion theory

- Translational and rotational invariance for an isolated molecule
- Kinetic energy operator: curvilinear coordinates, $\hat{T}(\rho)$
- Potential energy: $(3n_n - 6)$ -dimensional representations

Few-particle quantum theory

Electrons & atomic nuclei

- Translational and rotational invariance for an isolated molecule
- Kinetic energy operator:

$$-\sum_{i=1}^{n_p+1} \frac{1}{2m_i} \Delta_{r_i} \text{ or } \hat{T}(\rho)$$

- Potential energy:

$$\sum_{i=1}^{n_p+1} \sum_{j>i}^{n_p+1} \frac{q_i q_j}{r_{ij}}$$

The tremendous success of the usual practice might perhaps be best regarded as a tribute to the insight and ingenuity of the practitioners for inventing an effective variant of quantum theory for chemistry.

B. T. Sutcliffe and R. G. Woolley, J. Chem. Phys. 137, 22A544 (2012).

Molecules as few-particle quantum systems

Very accurate numerical results

- numerical solution of three-particle systems

Bishop and co-workers, Frolov and co-workers, ...

- positronium-electron complexes and nuclear physics

Suzuki and Varga, ...

- high-accuracy calculations for molecules

Adamowicz and co-workers, ...

⇒ high-resolution spectroscopy (incl. rel. effects)

Molecules as few-particle quantum systems

Very accurate numerical results

- numerical solution of three-particle systems
Bishop and co-workers, Frolov and co-workers, ...
- positronium-electron complexes and nuclear physics
Suzuki and Varga, ...
- high-accuracy calculations for molecules
Adamowicz and co-workers, ...

⇒ high-resolution spectroscopy (incl. rel. effects)

Interesting conceptual questions

- The molecular structure conundrum
Weininger, J. Chem. Educ. 61, 939 (1984).
Woolley, Sutcliffe, Primas, Claverie and Diner, ..., Isr. J. Chem. vol. 19 (1980)
- a single $\langle r_{C-H} \rangle$ in an organic molecule
e.g., B. T. Sutcliffe and R. G. Woolley, Chem. Phys. Lett. 408, 445 (2005).
- linear or triangular H_3^+ ?
M. Cafiero and L. Adamowicz, Chem. Phys. Lett., 387, 136 (2004).
 $\langle \alpha_{HHH} \rangle = (0^\circ + 180^\circ + 0^\circ)/3 = 60^\circ$ or $\langle \alpha_{HHH} \rangle = (60^\circ + 60^\circ + 60^\circ)/3 = 60^\circ$

⇒ classical features in a quantum description

Molecules as few-particle quantum systems

⇒ pre-Born–Oppenheimer theory

Schrödinger equation

$$\hat{H}\Psi = E\Psi \quad \text{with} \quad \hat{H}(\mathbf{m}, \mathbf{q}; \mathbf{r}) = -\sum_{i=1}^{n_p+1} \frac{1}{2m_i} \Delta_{\mathbf{r}_i} + \sum_{i=1}^{n_p+1} \sum_{j>i}^{n_p+1} \frac{q_i q_j}{|\mathbf{r}_i - \mathbf{r}_j|} \quad (1)$$

⇒ mass and electric charge: m_i, q_i ($i = 1, 2, \dots, n_p + 1$) parameters

Pauli principle

Spin statistics for fermions and bosons

⇒ spin: s_i ($i = 1, 2, \dots, n_p + 1$) parameters

$3(n_p + 1)$ physical parameters: m_i, q_i, s_i ($i = 1, 2, \dots, n_p + 1$)

⇒ **Variational solution with explicitly correlated Gaussians (ECGs)**

Adamowicz and co-workers, Suzuki and Varga

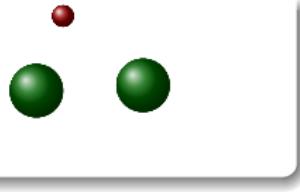
Pre-Born–Oppenheimer molecular structure theory

- Molecular structure in quantum mechanics
- Non-relativistic limit for small molecules

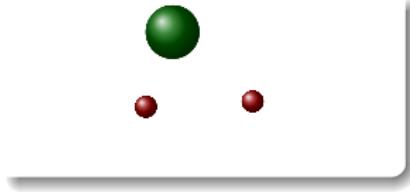
On the Born–Oppenheimer approximation

We describe the six degrees of freedom of the ground state of the **helium atom** (considered as 3-particle problem with the center-of-mass motion separated) as a problem of two interacting particles in an external Coulomb potential. However, in the case of the molecule H_2^+ we discuss the very same type of differential equation in an entirely different way, and split the 6 degrees of freedom into 1 vibrational mode, 2 rotational modes, and 3 electronic type degrees of freedom. This qualitatively different description does by no means follow from a purely mathematical discussion.

HANS PRIMAS: Chemistry, Quantum Mechanics and Reductionism
(Springer-Verlag, Berlin, Heidelberg and New York, 1981)



vs.



Molecular structure in quantum mechanics

Quantum chemistry: quantum mechanics + Born–Oppenheimer approximation

- equilibrium structure: local minimum structure of the PES
- we put in the molecular structure into this theoretical description

Molecular structure in quantum mechanics

Quantum chemistry: quantum mechanics + Born–Oppenheimer approximation

- equilibrium structure: local minimum structure of the PES
- we put in the molecular structure into this theoretical description

Quantum theory and chemistry?

How to recover the classical molecular structure from a purely quantum mechanical description?

- a non-trivial problem: "The molecular structure conundrum"
Weininger, J. Chem. Educ. 61, 939 (1984);
Woolley, Sutcliffe, Primas, Claverie and Diner, ..., Isr. J. Chem. vol. 19 (1980)

Let's calculate expectation values for the molecular structural parameters with the system's wave function...

- there is only a single $\langle r_{C-H} \rangle$ for any organic molecule
e.g., B. T. Sutcliffe and R. G. Woolley, Chem. Phys. Lett. 408, 445 (2005).
- given that $\langle \alpha_{HHH} \rangle = 60^\circ$ for the ground state wave function of H_3^+ ,
can we decide whether it is linear or triangular (or both or something else)?
M. Cafiero and L. Adamowicz, Chem. Phys. Lett., 387, 136 (2004).
 $\langle \alpha_{HHH} \rangle = (0^\circ + 180^\circ + 0^\circ)/3 = 60^\circ$ or $\langle \alpha_{HHH} \rangle = (60^\circ + 60^\circ + 60^\circ)/3 = 60^\circ$

⇒ problem: quantum mechanical indistinguishability of identical particles

Probabilistic interpretation of the wave function

Structural parameters

do not have sharp, dispersionless values, but they are characterized by **some probability density function**

1-particle density

$$\rho_{P,a}(\mathbf{R}_1) = \langle \Psi(\mathbf{r}) | \delta(\mathbf{r}_a - \mathbf{r}_P - \mathbf{R}_1) | \Psi(\mathbf{r}) \rangle \quad (2)$$

2-particle density

$$D_{P,ab}^{(2)}(\mathbf{R}_1, \mathbf{R}_2) = \langle \Psi(\mathbf{r}) | \delta(\mathbf{r}_a - \mathbf{r}_P - \mathbf{R}_1) \delta(\mathbf{r}_b - \mathbf{r}_P - \mathbf{R}_2) | \Psi(\mathbf{r}) \rangle \quad (3)$$

$$\Gamma_{P,ab}(\alpha) = \int_0^{\infty} dR_1 R_1^2 \int_0^{\infty} dR_2 R_2^2 D_{P,ab}^{(2)}(\mathbf{R}_1, \mathbf{R}_2) \quad \text{with} \quad \mathbf{R}_1^T \mathbf{R}_2 = R_1 R_2 \cos \alpha \quad (4)$$

On the emergence of molecular structure...

Coulomb Hamiltonian for $\{a^\pm, a^\pm, b^\mp\}$ systems

$$\hat{H}(m_a, m_b; \mathbf{r}) = -\frac{1}{2m_a} \Delta_{\mathbf{r}_1} - \frac{1}{2m_a} \Delta_{\mathbf{r}_2} - \frac{1}{2m_b} \Delta_{\mathbf{r}_3} + \frac{1}{r_{12}} - \frac{1}{r_{13}} - \frac{1}{r_{23}} \quad (5)$$

- charge-inversion symmetry
- mass-scale similarity

$$\hat{H}(\eta m_a, \eta m_b; \mathbf{r}) = \eta \hat{H}(m_a, m_b, \eta \mathbf{r}), \quad \eta \in \mathbb{R} \setminus \{0\} \quad (6)$$

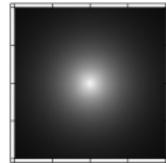
- bound for all m_a/m_b values [Martin et al., Phys. Rev. A 46, 3697 (1992)]
- let us consider the transition: $H^- \rightarrow H_2^+$

Particle density vs. mass ratio

$\rho_{0,a}(X, Y, 0)$ vs. m_a/m_b for $\{a^\pm, a^\pm, b^\mp\}$ -type systems

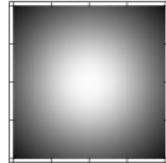
$N = 0, p = +1$ (COM is at the middle of each plot)

$$\frac{m_a}{m_b} : 0.000\,543$$

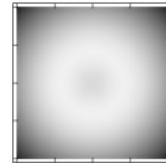


$\{e^-, e^-, p^+\}$

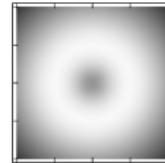
$$0.5$$



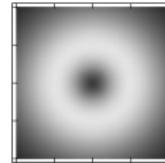
$$1$$



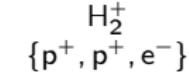
$$2$$



$$5$$



$$1840$$

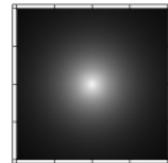


Particle density vs. mass ratio

$\rho_{0,a}(X, Y, 0)$ vs. m_a/m_b for $\{a^\pm, a^\pm, b^\mp\}$ -type systems

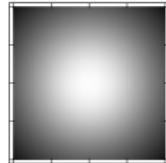
$N = 0, p = +1$ (COM is at the middle of each plot)

$$\frac{m_a}{m_b} : 0.000\,543$$

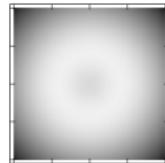


$\{\text{e}^-, \text{e}^-, \text{p}^+\}$

$$0.5$$

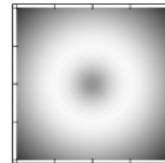


$$1$$

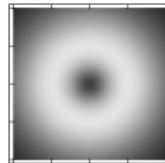


$\{\text{e}^-, \text{e}^-, \text{e}^+\}$

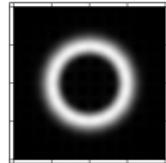
$$2$$



$$5$$



$$5$$



$\{\text{p}^+, \text{p}^+, \text{e}^-\}$

Superposition vs. rotating dumbbell



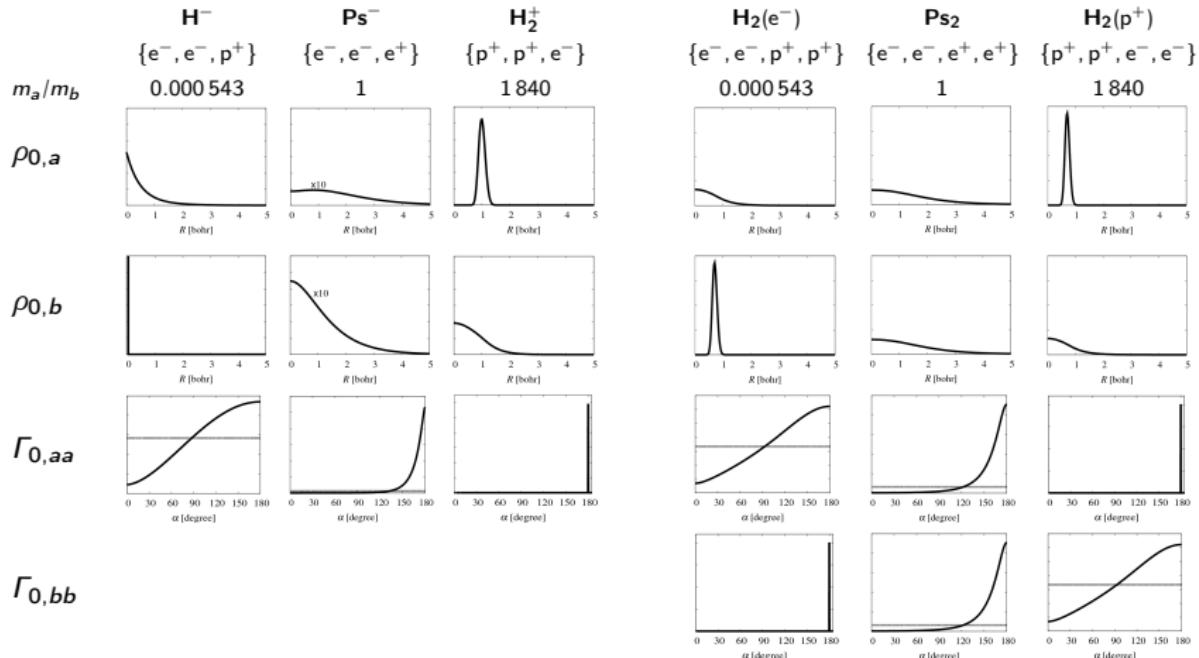
or



E. Mátýus, J. Hutter, U. Müller-Herold, and M. Reiher, Phys. Rev. A 83, 052512 (2011); J. Chem. Phys. 135, 204302 (2011).

Radial and angular densities

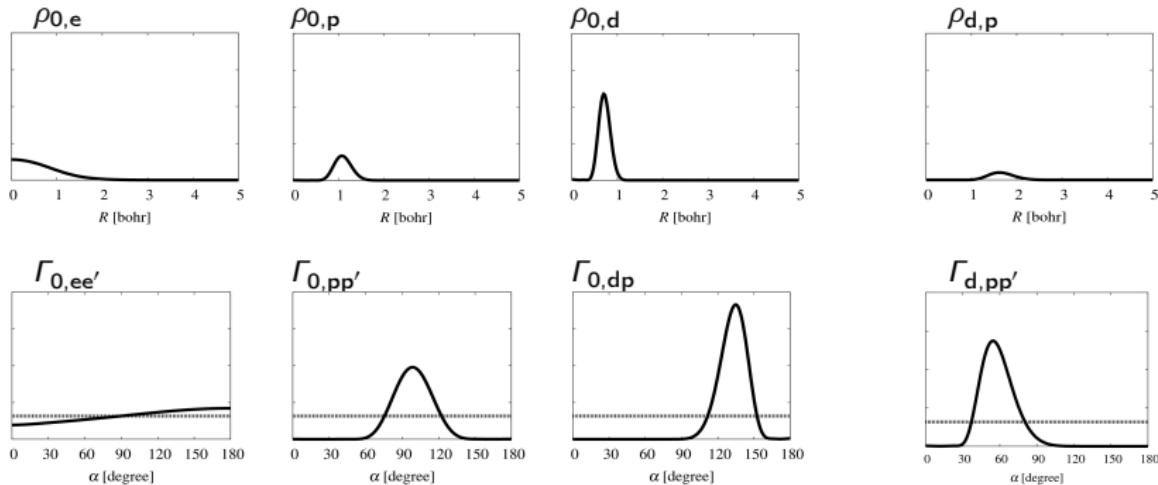
$\{a^\pm, a^\pm, b^\mp\}$ - and $\{a^\pm, a^\pm, b^\mp, b^\mp\}$ -type systems



⇒ “mass-polarization effects” [Fröman & Kinsey (1961) and Woolley (1976)]

Radial and angular densities

$$\text{H}_2\text{D}^+ = \{\text{e}^-, \text{e}^-, \text{p}^+, \text{p}^+, \text{d}^+\}$$



Remark: qualitative features of the radial and angular densities converged quickly

⇒ “molecular structure seen as correlation effects for the nuclei”

Claverie and Diner, Isr. J. Chem. 19, 54 (1980).

⇒ N.B. “traditional route”: introduce internal coordinates to describe nuclear motion...

Pre-Born–Oppenheimer molecular structure theory

- Molecular structure in quantum mechanics
- **Non-relativistic limit for small molecules**

Molecules as few-particle quantum systems

⇒ pre-Born–Oppenheimer theory

Schrödinger equation

$$\hat{H}\psi = E\psi \quad \text{with} \quad \hat{H}(\mathbf{m}, \mathbf{q}; \mathbf{r}) = - \sum_{i=1}^{n_p+1} \frac{1}{2m_i} \Delta_{\mathbf{r}_i} + \sum_{i=1}^{n_p+1} \sum_{j>i}^{n_p+1} \frac{q_i q_j}{|\mathbf{r}_i - \mathbf{r}_j|} \quad (7)$$

⇒ mass and electric charge: m_i, q_i ($i = 1, 2, \dots, n_p + 1$) parameters

Pauli principle

Spin statistics for fermions and bosons

⇒ spin: s_i ($i = 1, 2, \dots, n_p + 1$) parameters

$3(n_p + 1)$ physical parameters: m_i, q_i, s_i ($i = 1, 2, \dots, n_p + 1$)

⇒ **Variational solution with explicitly correlated Gaussians (ECGs)**

Adamowicz and co-workers, Suzuki and Varga

Molecules as few-particle quantum systems

⇒ pre-Born–Oppenheimer theory

Schrödinger equation

$$\hat{H}\psi = E\psi \quad \text{with} \quad \hat{H}(\mathbf{m}, \mathbf{q}; \mathbf{r}) = - \sum_{i=1}^{n_p+1} \frac{1}{2m_i} \Delta_{\mathbf{r}_i} + \sum_{i=1}^{n_p+1} \sum_{j>i}^{n_p+1} \frac{q_i q_j}{|\mathbf{r}_i - \mathbf{r}_j|} \quad (7)$$

⇒ mass and electric charge: m_i, q_i ($i = 1, 2, \dots, n_p + 1$) parameters

Pauli principle

Spin statistics for fermions and bosons

⇒ spin: s_i ($i = 1, 2, \dots, n_p + 1$) parameters

$3(n_p + 1)$ physical parameters: m_i, q_i, s_i ($i = 1, 2, \dots, n_p + 1$)

⇒ **Variational solution with explicitly correlated Gaussians (ECGs)**

Adamowicz and co-workers, Suzuki and Varga

⇒ **Challenge: molecules with various quantum numbers**

Variational solution with explicitly correlated Gaussian functions

Key ingredients of the variational procedure

Input: number of particles, mass, electric charge, spin

Quantum numbers: total angular momentum (N, M_N), parity (p), spin ($S_a, M_{S_a}; S_b, M_{S_b}; \dots$)

- **Cartesian coordinates**
laboratory-fixed or translationally invariant (linear transformation)
- **Hamiltonian is simple in Cartesian coordinates**
- **symmetry-adapted explicitly correlated Gaussian functions**
 - global vector representation
a generalized angular momentum coupling scheme for a variational approach
[Y. Suzuki, J. Usukura, and K. Varga, J. Phys. B 31, 31 (1998)]
 - flexible parameterization
- **analytic matrix elements**
- **generalized eigenvalue problem**, direct solution
($\sim 1000 \times 1000$ real symmetric matrices)
- ◊ **variational principle**
optimization of the basis function parameters (stochastic and Powell's method)

Variational solution with explicitly correlated Gaussian functions

Key ingredients of the variational procedure

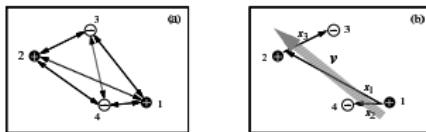
Input: number of particles, mass, electric charge, spin

Quantum numbers: total angular momentum (N, M_N), parity (p), spin ($S_a, M_{S_a}; S_b, M_{S_b}; \dots$)

- **Cartesian coordinates**
laboratory-fixed or translationally invariant (linear transformation)
- **Hamiltonian is simple** in Cartesian coordinates
- **symmetry-adapted explicitly correlated Gaussian functions**
 - global vector representation
a generalized angular momentum coupling scheme for a variational approach
[Y. Suzuki, J. Usukura, and K. Varga, J. Phys. B 31, 31 (1998)]
 - flexible parameterization
- **analytic matrix elements**
- **generalized eigenvalue problem**, direct solution
($\sim 1000 \times 1000$ real symmetric matrices)
- ◊ **variational principle**
optimization of the basis function parameters (stochastic and Powell's method)

Symmetry-adapted basis functions

Basis function

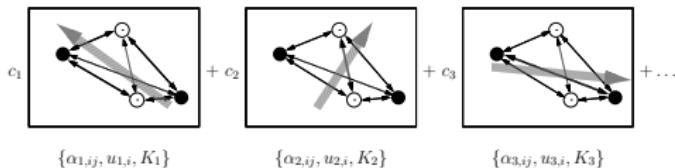


$$\phi^{(\lambda)}(\mathbf{r}; \mathbf{A}, \mathbf{u}, K) = |\mathbf{v}|^{2K+N} Y_N^{M_N}(\hat{\nu}) \exp\left(-\frac{1}{2} \sum_{i=1}^{n_p+1} \sum_{j>i}^{n_p+1} \alpha_{ij} (\mathbf{r}_i - \mathbf{r}_j)^2\right) \quad \text{with} \quad \mathbf{v} = \sum_{i=1}^{n_p+1} u_i \mathbf{r}_i$$

parameters: α_{ij} (\mathbf{A}), u_i (\mathbf{u}), K

$[\lambda = (N, p), p = (-1)^N]$

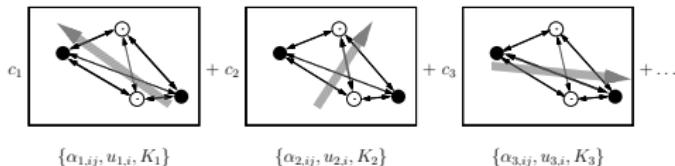
Basis-set expansion



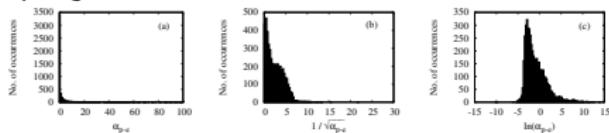
$$\Psi_0^{(\lambda, \varsigma)} = c_1 \hat{\mathcal{A}} \{\chi_1^{(\varsigma)} \phi^{(\lambda)}(\mathbf{r}; \{\alpha_{ij,1}, u_{i,1}, K_1\})\} + c_2 \hat{\mathcal{A}} \{\chi_2^{(\varsigma)} \phi^{(\lambda)}(\mathbf{r}; \{\alpha_{ij,2}, u_{i,2}, K_2\})\} + c_3 \hat{\mathcal{A}} \{\chi_3^{(\varsigma)} \phi^{(\lambda)}(\mathbf{r}; \{\alpha_{ij,3}, u_{i,3}, K_3\})\} + \dots$$

parameter set: $\{\alpha_{I,ij}(\mathbf{A}_I), u_{I,i}(\mathbf{u}_I), K_I, (I = 1, \dots, N_b)\}$

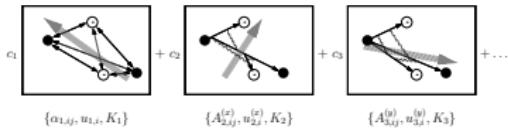
Parameterization of the basis functions



- variational principle
- fast eigenvalue estimator for a rank-1 update (known theorem)
- sampling-importance sampling



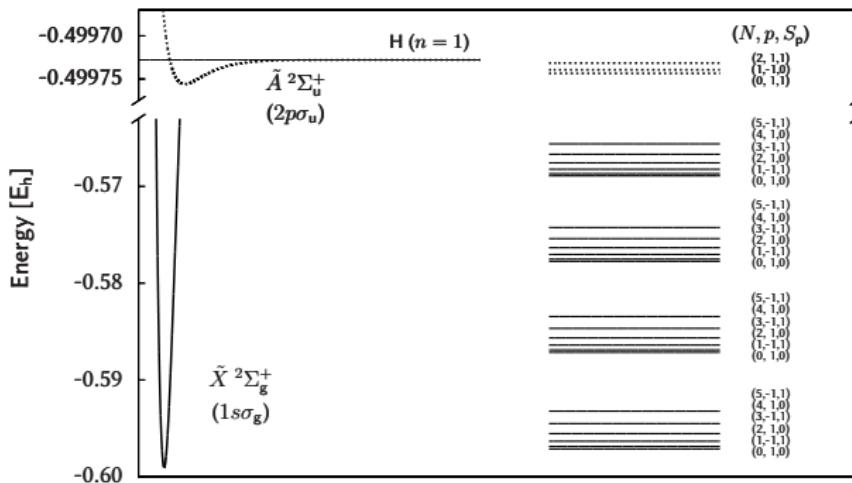
- fine-tuning: random walk or Powell's method
- multi-channel optimization



- parameter-transfer approach

$$\psi(\mathbf{A}') = \sum_{I=1}^{N_b} c_I \phi_I^{(\mathbf{A}')} [\mathcal{P}_I(\mathbf{A})]$$

Test case: $\text{H}_2^+ = \{\text{p}^+, \text{p}^+, \text{e}^-\}$



Energy levels and dissociation energies (examples)

| N | p | S_p | E / E_h | D / cm^{-1} | Assignment |
|-----|-----|-------|-------------------|----------------------|--|
| 2 | 1 | 1 | -0.499 731 516(7) | 0.807(1) | $\tilde{A} \ 2\Sigma_u^+, v = 0$; ortho |
| ... | | | | | |
| 0 | 1 | 0 | -0.597 139 060(4) | 21 379.290(2) | $\tilde{X} \ 2\Sigma_g^+, v = 0$; para |

$$\text{H}_2 = \{\text{p}^+, \text{p}^+, \text{e}^-, \text{e}^-\}$$

parameterization:

- sampling-importance sampling and Powell's method for fine-tuning
- parameter-transfer approach

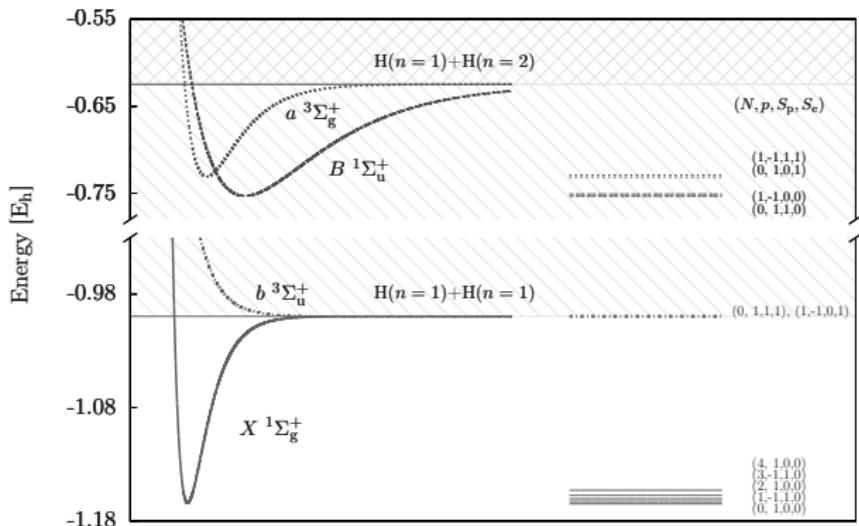
Calculated energy levels of H₂ with N = 0, 1, and 2
total (orbital plus rotational) angular momentum quantum number

| (N, p, S _p , S _e) ^a | E/E _h ^b | ΔE _{Ref} /μE _h ^c | Ref. | Assignment ^d |
|---|-------------------------------|---|------|---|
| (0, +1, 0, 0) | -1.164 025 030 | -0.000 6 | [1] | X ^{1Σ_g⁺} |
| (1, -1, 1, 0) | -1.163 485 171 | -0.001 4 | [1] | X ^{1Σ_g⁺} |
| (2, +1, 0, 0) | -1.162 410 408 | -0.001 9 | [1] | X ^{1Σ_g⁺} |
| (0, +1, 1, 0) | -0.753 027 186 | 0.135 4 | [2] | B ^{1Σ_u⁺} |
| (1, -1, 0, 0) | -0.752 850 233 | 0.834 2 | [2] | B ^{1Σ_u⁺} |
| (2, +1, 1, 0) | -0.752 498 022 | 0.918 8 | [2] | B ^{1Σ_u⁺} |
| (0, +1, 0, 1) | -0.730 825 193 | -0.006 9 | [3] | a ^{3Σ_g⁺} |
| (1, -1, 1, 1) | -0.730 521 418 | 0.008 0 | [3] | a ^{3Σ_g⁺} |
| (2, +1, 0, 1) | -0.729 916 268 | 0.047 9 | [3] | a ^{3Σ_g⁺} |
| (0, +1, 1, 1) | [-0.999 450 102] ^e | [-5.578] ^f | | b ^{3Σ_u⁺} |
| (1, -1, 0, 1) | [-0.999 445 835] ^e | [-9.844] ^f | | b ^{3Σ_u⁺} |
| (2, +1, 1, 1) | [-0.999 439 670] ^e | [-16.010] ^f | | b ^{3Σ_u⁺} |

[1]: Pachucki & Komasa, JCP 130, 164113 (2009), [2]: Wolniewicz, Orlikowski, and Staszewska, JMS 238, 118 (2006),

[3]: Wolniewicz, MP 105, 1497 (2007).

$$\text{H}_2 = \{\text{p}^+, \text{p}^+, \text{e}^-, \text{e}^-\}$$



singlet, para

even $N, p = +1$: $X^1\Sigma_g^+$
odd $N, p = -1$: $B^1\Sigma_u^+$



triplet, ortho

even $N, p = +1$: $b^3\Sigma_u^+$
odd $N, p = -1$: $a^3\Sigma_g^+$



singlet, ortho

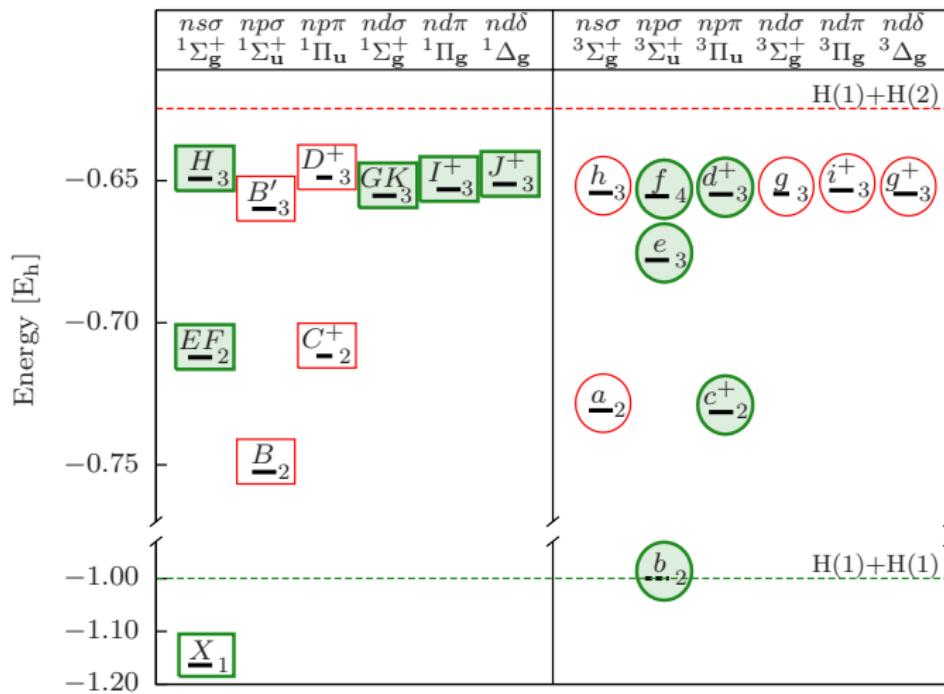
even $N, p = +1$: $B^1\Sigma_u^+$
odd $N, p = -1$: $X^1\Sigma_g^+$



triplet, para

even $N, p = +1$: $a^3\Sigma_g^+$
odd $N, p = -1$: $b^3\Sigma_u^+$

Orientation chart: electronic states of H₂



EM, J. Phys. Chem. A 117, 7195 (2013).

J. Brown and A. Carrington, Rotational Spectroscopy of Diatomic Molecules, Cambridge University Press, Cambridge (2003).

Towards the calculation of rovibronic resonances

Complex-coordinate rotation method

Complex coordinate rotation technique for the non-relativistic Coulomb Hamiltonian, $\mathbf{r} \rightarrow \mathbf{r}e^{i\theta}$:

$$\hat{H} = \hat{T} + \hat{V} \rightarrow \hat{H}(\theta) = e^{-2i\theta} \hat{T} + e^{-i\theta} \hat{V} \quad (8)$$

Calculation of eigenvalues:

$$\mathbf{H}(\theta)\mathbf{v}(\theta)_i = E_i(\theta)\mathbf{S}\mathbf{v}(\theta)_i \quad (9)$$

which is replaced by

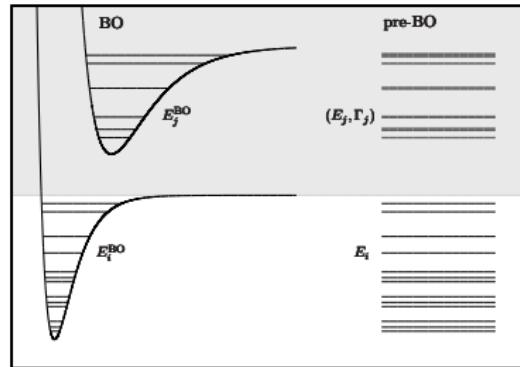
$$\mathbf{H}'(\theta)\mathbf{v}'(\theta)_i = E_i(\theta)\mathbf{v}'(\theta)_i \quad (10)$$

with

$$\begin{aligned} \mathbf{H}'(\theta) &= e^{-2i\theta} \mathbf{S}^{-1/2} \mathbf{T} \mathbf{S}^{-1/2} + e^{-i\theta} \mathbf{S}^{-1/2} \mathbf{V} \mathbf{S}^{-1/2} \\ &= \cos(2\theta) \mathbf{T}' + \cos(\theta) \mathbf{V}' - i(\sin(2\theta) \mathbf{T}' + \sin(\theta) \mathbf{V}') \end{aligned} \quad (11)$$

The complex symmetric eigenproblem is solved using LAPACK.

Motivation



Complex variational principle?

Ps⁻ = {e⁺, e⁻, e⁻} resonances

parameterization

- energy minimization
- random trials (sampling-importance sampling)
- it is easy to generate a reasonably good parameter set

Identified bound and resonance states for Ps⁻ (in E_h)

| (N, p, S ₋) ^b | Re(ε) ^c | Γ/2 ^c | Re(ε _{Ref}) ^c | Γ _{Ref} /2 ^c | Ref. |
|--------------------------------------|-----------------------------|--------------------------|------------------------------------|----------------------------------|------|
| (0, +1, 0) | -0.262 005 070 ^d | 0 ^d | -0.262 005 070 | 0 | [1] |
| (0, +1, 0) | -0.076 030 455 | 2.152 · 10 ⁻⁵ | -0.076 030 442 | 2.151 7 · 10 ⁻⁵ | [2] |
| (0, +1, 0) | -0.063 649 173 | 4.369 · 10 ⁻⁶ | -0.063 649 175 | 4.339 3 · 10 ⁻⁶ | [2] |
| (0, +1, 0) | -0.062 609 | 2.5 · 10 ⁻⁵ | -0.062 550 | 5.0 · 10 ⁻⁷ | [3] |
| (0, +1, 0) | -0.035 341 850 | 3.730 · 10 ⁻⁵ | -0.035 341 885 | 3.732 9 · 10 ⁻⁵ | [2] |
| (0, +1, 0) | -0.029 845 700 | 2.781 · 10 ⁻⁵ | -0.029 846 146 | 2.635 6 · 10 ⁻⁵ | [2] |
| (0, +1, 0) | -0.028 271 | 1.8 · 10 ⁻⁵ | -0.028 200 | 7.5 · 10 ⁻⁶ | [3] |
| (0, +1, 0) | -0.020 199 000 | 8.800 · 10 ⁻⁵ | -0.020 213 921 | 6.502 6 · 10 ⁻⁵ | [2] |
| (0, +1, 1) | -0.063 537 352 | 2.132 · 10 ⁻⁹ | -0.063 537 354 | 1.570 0 · 10 ⁻⁹ | [2] |
| (0, +1, 1) | -0.062 591 | 2.6 · 10 ⁻⁷ | -0.062 550 | 2.5 · 10 ⁻¹⁰ | [3] |
| (0, +1, 1) | -0.029 369 870 | 1.300 · 10 ⁻⁷ | -0.029 370 687 | 9.395 0 · 10 ⁻⁸ | [2] |
| (0, +1, 1) | -0.028 21 | 1.9 · 10 ⁻⁵ | -0.028 05 | 5.0 · 10 ⁻⁸ | [3] |
| (0, +1, 1) | -0.017 070 800 | 6.710 · 10 ⁻⁶ | -0.017 101 172 | 3.560 9 · 10 ⁻⁷ | [3] |

$$E(\text{Ps}(1)) = -1/4 = -0.25 \text{ E}_h, E(\text{Ps}(2)) = -1/16 = -0.0625 \text{ E}_h \text{ and } E(\text{Ps}(3)) = -1/36 = -0.027 \frac{7}{9} \text{ E}_h$$

[1]: Korobov, PRA 61, 064503 (2000), [2]: Li & Shakeshaft, PRA 71, 052505 (2005), [3]: Papp et al., Phys. Lett. A 304, 36 (2002).

Ps₂ = {e⁺, e⁺, e⁻, e⁻} resonances

parameterization:

- energy minimization
- sampling-importance sampling

Identified bound and resonance states for Ps₂ (in E_h)

| (N, p, c) ^b | (S_-, S_+) ^c | $\text{Re}(\mathcal{E})$ ^d | $\Gamma/2$ ^d | $\text{Re}(\mathcal{E}_{\text{Ref}})$ ^d | $\Gamma_{\text{Ref}}/2$ ^d | Ref. |
|--------------------------|---------------------------|---------------------------------------|-------------------------|--|--------------------------------------|------|
| (0, +1, +1) | (0, 0) | -0.516 003 789 741 ^e | 0 ^e | -0.516 003 790 416 | 0 | [1] |
| (0, +1, +1) | (0, 0) | -0.329 38 | $3.03 \cdot 10^{-3}$ | -0.329 4 | $3.1 \cdot 10^{-3}$ | [2] |
| (0, +1, +1) | (0, 0) | -0.291 7 | $2.5 \cdot 10^{-3}$ | -0.292 4 | $1.95 \cdot 10^{-3}$ | [2] |
| (0, +1, -1) | (0, 0) | -0.314 677 072 ^e | 0 ^e | -0.314 673 3 | 0 | [2] |
| (0, +1, -1) | (0, 0) | -0.289 789 3 | $7.7 \cdot 10^{-5}$ | -0.289 76 | $7 \cdot 10^{-5}$ | [2] |
| (0, +1, -1) | (0, 0) | -0.279 25 | $2.3 \cdot 10^{-4}$ | -0.279 13 | $1 \cdot 10^{-4}$ | [2] |
| (0, +1, +1) | (1, 1) | -0.277 2 | $5.4 \cdot 10^{-4}$ | -0.276 55 | $1.55 \cdot 10^{-4}$ | [2] |
| (0, +1, -1) | (1, 1) | -0.309 0 | $5.7 \cdot 10^{-3}$ | -0.308 14 | $1.2 \cdot 10^{-4}$ | [2] |
| (0, +1, -1) | (1, 1) | -0.273 3 | $2.3 \cdot 10^{-3}$ | -0.273 6 | $8.5 \cdot 10^{-4}$ | [2] |
| (0, +1, ±1) | (1, 0)/(0, 1) | -0.330 287 505 ^e | 0 ^e | -0.330 276 81 | 0 | [2] |
| (0, +1, ±1) | (1, 0)/(0, 1) | -0.294 3 | $3.1 \cdot 10^{-3}$ | -0.293 9 | $2.15 \cdot 10^{-3}$ | [2] |
| (0, +1, ±1) | (1, 0)/(0, 1) | -0.282 | $2 \cdot 10^{-3}$ | -0.282 2 | $8.5 \cdot 10^{-4}$ | [2] |

$$E(\text{Ps}(1) + \text{Ps}(1)) = -1/2 = -0.5 \text{ E}_h \text{ and } E(\text{Ps}(1) + \text{Ps}(2)) = -5/16 = -0.3125 \text{ E}_h$$

[1]: Bubin & Adamowicz, PRA 74, 052502 (2006), [2]: Suzuki & Usukura, Nucl. Inst. Meth. Phys. Res. B 221, 195 (2004).

Hunt for $\text{H}_2 = \{\text{p}^+, \text{p}^+, \text{e}^-, \text{e}^-\}$ resonances

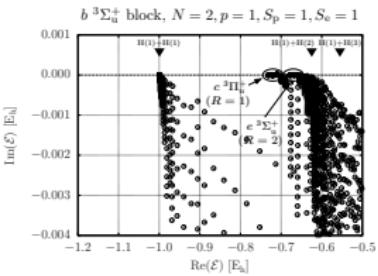
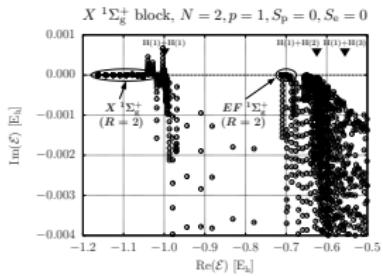
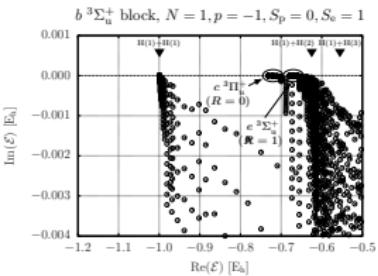
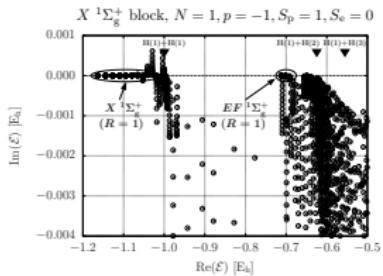
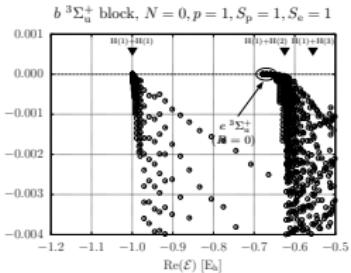
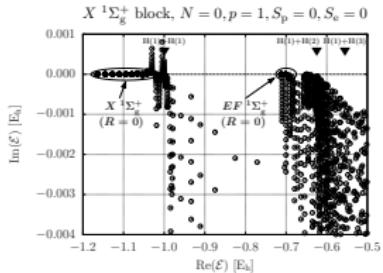
difficult to find a good parameter set, earlier strategies fail:

- energy minimization, sampling-importance sampling
- ✓ idea: parameter-transfer approach

Identified resonance states of H_2 embedded in the $b\ 3\Sigma_u^+$ continuum (in E_{h})

| (N, p, S_p, S_e) ^a | $\text{Re}(\mathcal{E})$ ^b | $\Gamma/2$ ^b | $E_{\text{Ref,exp}}$ ^c | $E_{\text{Ref,theo}}$ ^d | Assignment ^e |
|---------------------------------|--|-------------------------|-----------------------------------|------------------------------------|-------------------------------------|
| $(0, +1, 1, 1)$ | $[-0.999\ 450\ 1]$ ^f [...] | | | $[-0.999\ 455\ 7]$ | $\text{H}(1)+\text{H}(1)$ continuum |
| $(0, +1, 1, 1)$ | $-0.677\ 947\ 1$ | $1 \cdot 10^{-7}$ | $-0.677\ 946\ 1$ | $-0.677\ 942\ 7^*$ | $e\ 3\Sigma_u^+, R = 0, v = 0$ |
| $(0, +1, 1, 1)$ | $-0.668\ 549\ 3$ | $9 \cdot 10^{-7}$ | $-0.668\ 547\ 8$ | $-0.668\ 541\ 0^*$ | $e\ 3\Sigma_u^+, R = 0, v = 1$ |
| $(1, -1, 0, 1)$ | $[-0.999\ 445\ 8]$ ^f [...] | | | $[-0.999\ 455\ 7]$ | $\text{H}(1)+\text{H}(1)$ continuum |
| $(1, -1, 0, 1)$ | $-0.731\ 434\ 0$ | $5 \cdot 10^{-7}$ | $-0.731\ 438\ 8$ | $-0.731\ 469\ 1^{**}$ | $c\ 3\Pi_u^+, R = 0, v = 0$ |
| $(1, -1, 0, 1)$ | $-0.720\ 717\ 5$ | $2 \cdot 10^{-7}$ | $-0.720\ 782\ 6$ | | $c\ 3\Pi_u^+, R = 0, v = 1$ |
| $(1, -1, 0, 1)$ | | [...] | | | |
| $(1, -1, 0, 1)$ | $-0.677\ 705\ 5$ | $2 \cdot 10^{-7}$ | $-0.677\ 704\ 1$ | $-0.677\ 698\ 2^*$ | $e\ 3\Sigma_u^+, R = 1, v = 0$ |
| $(1, -1, 0, 1)$ | $-0.668\ 319\ 5$ | $1 \cdot 10^{-6}$ | $-0.668\ 319\ 7$ | $-0.668\ 309\ 8^*$ | $e\ 3\Sigma_u^+, R = 1, v = 1$ |
| $(2, +1, 1, 1)$ | $[-0.999\ 439\ 7]$ ^f [...] | | | $[-0.999\ 455\ 7]$ | $\text{H}(1)+\text{H}(1)$ continuum |
| $(2, +1, 1, 1)$ | $-0.730\ 888\ 2$ | $9 \cdot 10^{-7}$ | $-0.730\ 888\ 7$ | | $c\ 3\Pi_u^+, R = 1, v = 0$ |
| $(2, +1, 1, 1)$ | $-0.720\ 219\ 0$ | $< 2 \cdot 10^{-7}$ | $-0.720\ 258\ 0$ | | $c\ 3\Pi_u^+, R = 1, v = 1$ |
| $(2, +1, 1, 1)$ | | [...] | | | |
| $(2, +1, 1, 1)$ | $-0.677\ 222\ 9$ | $2 \cdot 10^{-8}$ | $-0.677\ 222\ 2$ | | $e\ 3\Sigma_u^+, R = 2, v = 0$ |
| $(2, +1, 1, 1)$ | $-0.667\ 863\ 2$ | $7 \cdot 10^{-7}$ | $-0.667\ 865\ 3$ | | $e\ 3\Sigma_u^+, R = 2, v = 1$ |

* Kołos & Rychlewski (1990), ** Kołos & Rychlewski (1977), $E_{\text{Ref,exp}} = E_0 + T_{\text{exp}}$ with $E_0 = -1.164\ 025\ 030\ \text{E}_h$ (g.s.).



Summary and outlook

Pre-Born–Oppenheimer Molecular Structure Theory

● Quantum structure of molecules

On the Emergence of Molecular Structure

EM, J. Hutter, U. Müller-Herold, and M. Reiher, Phys. Rev. A 83, 052512 (2011).

Extracting Elements of Molecular Structure from the All-Particle Wave Function

EM, J. Hutter, U. Müller-Herold, and M. Reiher, J. Chem. Phys. 135, 204302 (2011).

● Approaching the non-relativistic limit for various quantum numbers

On the Calculation of Resonances in pre-Born–Oppenheimer Molecular Structure Theory

EM, J. Phys. Chem. A 117, 7195 (2013).

Molecular Structure Calculations: a Unified Quantum Mechanical Description of Electrons and Nuclei using Explicitly Correlated Gaussian Functions and the Global Vector Representation

EM and M. Reiher, J. Chem. Phys. 137, 024104 (2012).

● Towards a relativistic four-component pre-Born–Oppenheimer theory

Relativistic Kinetic-Balance Condition for Explicitly Correlated Basis Functions

B. Simmen, EM, and M. Reiher, J. Phys. B 48, 245004 (2015).

Electric Transition Dipole Moment in pre-Born–Oppenheimer Molecular Structure Theory

B. Simmen, EM, and M. Reiher, J. Chem. Phys. 141, 154105 (2014).

Elimination of the Translational Kinetic Energy Contamination in pre-Born–Oppenheimer Calculations

B. Simmen, EM, and M. Reiher, Mol. Phys. 111, 2086 (2013).

