

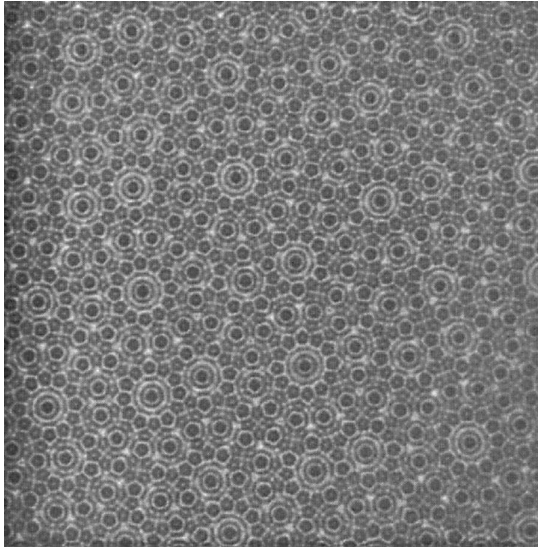
Existence of quasipatterns, solutions of the Bénard - Rayleigh convection problem

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Experiment of Faraday type. Kudrolli, Pier, Gollub 1998

Steady Bénard - Rayleigh system between two horizontal planes

$$\begin{aligned}V \cdot \nabla V + \nabla p &= \mathcal{P}(\theta \mathbf{e}_z + \mathcal{R}^{-1/2} \Delta V), \\V \cdot \nabla \theta &= \mathcal{R}^{-1/2} \Delta \theta + V \cdot \mathbf{e}_z, \\ \nabla \cdot V &= 0.\end{aligned}$$

Boundary Conditions: $v^{(z)} = \theta = 0$ in $z = 0, 1$.

either "rigid - rigid": $V^{(H)} = 0$ in $z = 0, 1$,

or "rigid - free": $V^{(H)} = 0$ in $z = 0$, $\frac{\partial V^{(H)}}{\partial z} = 0$ in $z = 1$,

or "free - rigid": $\frac{\partial V^{(H)}}{\partial z} = 0$ in $z = 0$, $V^{(H)} = 0$ in $z = 1$.

We do not consider the "free-free" case: $\frac{\partial V^{(H)}}{\partial z} = 0$ in $z = 0, 1$.

We choose to look for bifurcating solutions, **quasiperiodic in $\mathbf{x} \in \mathbb{R}^2$, invariant under rotations of angle π/q .**

$$u = \sum_{\mathbf{k} \in \Gamma} u^{(\mathbf{k})} e^{i\mathbf{k} \cdot \mathbf{x}}, \quad u^{(-\mathbf{k})} = \overline{u^{(\mathbf{k})}}$$

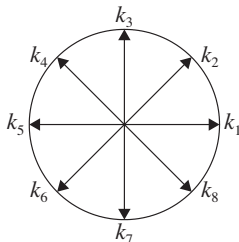
$$\Gamma = \left\{ \mathbf{k} = \sum_{j=1, \dots, 2q} m_j \mathbf{k}_j, \quad m \in \mathbb{N}^{2q}, \quad |\mathbf{k}_j| = k_c, \quad (\mathbf{k}_j, \mathbf{k}_{j+1}) = \pi/q \right\}$$

For $q = 1, 2, 3$ Γ is a lattice leading to a periodic pattern

See V.Yudovich et al (1963-67), W.Velte (1964-69),

K.Kirchgässner et al (1967-73), P.Rabinowitz (1968)

For $q \geq 4$ Γ is a quasilattice leading to a **quasipattern**



Example $q = 4$, the 8 wavevectors which form the basis of the quasilattice

Diophantine estimate

\mathbb{Q} vector space $\text{span}\{\mathbf{k}_j; j = 1, \dots, 2q\}$ has dimension d ,
 $d/2 = l_0 + 1 \leq q/2$ is the degree of the minimal Polynomial for
the algebraic integer $\omega = 2 \cos \pi/q$ (coef in \mathbb{Z} and first coef = 1).

$$\mathbf{k} = \sum_{j=1}^{2q} m_j \mathbf{k}_j = \frac{1}{\vartheta} \sum_{s=1}^d m_s^* \mathbf{k}_s^*, \quad \mathbf{m}^* = (m_1^*, \dots, m_d^*) \in \mathbb{Z}^d$$

$$N_{\mathbf{k}} = \sum_{s=1}^d |m_s^*|$$

for $q = 4, 5, 6$, $\omega = \sqrt{2}, \frac{1 + \sqrt{5}}{2}, \sqrt{3}$, $l_0 = 1$, $d = 4$

for $q = 7, 9$, $l_0 = 2$, $d = 6$,

for $q = 8, 10, 12$, $l_0 = 3$, $d = 8$, for $q = 11$, $l_0 = 4$, $d = 10 \dots$

For $q = 4, 5, \dots, 12$ then $\vartheta = 1$ and $\mathbf{k}_s^* = \mathbf{k}_s$, $s = 1, \dots, d$.

In all cases, there exists $c > 0$ such that

$$(|\mathbf{k}|^2 - k_c^2)^2 \geq c(1 + N_{\mathbf{k}}^2)^{-2l_0}, \text{ if } \mathbf{k} \neq \mathbf{k}_j, j = 1, \dots, 2q$$

Suitable formulation

Hilbert space for the 4-components vector field $u = (V, \theta)$:

$$\mathcal{K}_s = \left\{ u = (V, \theta)(\mathbf{x}, z) = \sum_{\mathbf{k} \in \Gamma} u_{\mathbf{k}}(z) e^{i\mathbf{k} \cdot \mathbf{x}}; \nabla \cdot V = 0, v_z|_{z=0,1} = 0, \right. \\ \left. \sum_{\mathbf{k} \in \Gamma} \left((1 + N_{\mathbf{k}}^2)^s \|u_{\mathbf{k}}\|_{L^2(0,1)}^2 \right) < \infty \right\}$$

For $s > d/2$, $\lambda := \mathcal{R}^{-1/2} > 0$, finding a solution $u \in \mathcal{K}_s$ of

$$\lambda u - \mathcal{A}u + \mathcal{B}(u, u) = 0,$$

is equivalent to find a classical solution of Bénard-Rayleigh system. In \mathcal{K}_s , \mathcal{A} is linear bounded, **selfadjoint**, \mathcal{B} is quadratic, bounded. Operators \mathcal{A} and \mathcal{B} **commute with \mathbf{R}_ϕ** defined by

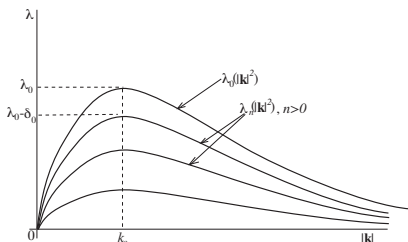
$$\mathbf{R}_\phi u = (R_\phi V(R_{-\phi} \mathbf{x}, z), \theta(R_{-\phi} \mathbf{x}, z))$$

We are interested by quasipatterns solutions of the problem which are invariant under **$\mathbf{R}_{\pi/q}$**

Study of the linear equation :

$$(\mathcal{A} - \lambda)u = G \in \mathcal{K}_s$$

comes back to the study made by V. Yudovich (1966) in the periodic case, for any wave number $|\mathbf{k}|$.



positive eigenvalues of \mathcal{A}

Spectrum of \mathcal{A} : real interval $[-\lambda_0, \lambda_0]$

$\lambda = \lambda_0$ largest e.v. $\text{Ker}(\mathcal{A} - \lambda_0)$ spanned by

$$\xi_j = \mathbf{R} \frac{\pi(j-1)}{q} \left(\widehat{U}_{\mathbf{k}_1}(z) e^{i\mathbf{k}_1 \cdot \mathbf{x}} \right), \quad j = 1, 2, \dots, 2q, \quad \text{with } |\mathbf{k}_1| = k_c$$

$$(\mathcal{A} - \lambda_0)u = G \in \mathcal{K}_s$$

Assume $\frac{d^2 \lambda_0}{d|\mathbf{k}|^2} |_{k_c} \neq 0$.

For G satisfying compatibility conditions $\langle G, \xi_j \rangle = 0, j = 1, \dots, 2q$

$$\|u_{\mathbf{k}}\|_0 \leq \frac{c(|\mathbf{k}|^2)}{(|\mathbf{k}|^2 - k_c^2)^2} \|G_{\mathbf{k}}\|_0, \mathbf{k} \in \Gamma \text{ except } \mathbf{k}_j, j = 1, 2, \dots, 2q$$

$c(|\mathbf{k}|^2)$ is analytic and $\mathcal{O}(|\mathbf{k}|^4)$ as $|\mathbf{k}|^2$ tends towards ∞

diophantine estimate: $\frac{1}{(|\mathbf{k}|^2 - k_c^2)^2} \leq C(1 + N_{\mathbf{k}}^2)^{2l_0}$ for $\mathbf{k} \neq \mathbf{k}_j$

Formal series - Approximate solution

$$(\mathcal{A} - \lambda_0)u = -\mu u + \mathcal{B}(u, u), \quad \lambda = \lambda_0 - \mu$$

$$u = \sum_{n \geq 1} \varepsilon^n u_n, \quad \mu = \sum_{n \geq 1} \varepsilon^n \mu_n, \quad \text{with } u_n \in \mathcal{K}_s, \quad \langle u_n, u_1 \rangle_0 = 0, \quad n \geq 2$$

$$(\mathcal{A} - \lambda_0)u_1 = 0,$$

$$(\mathcal{A} - \lambda_0)u_2 = -\mu_1 u_1 + \mathcal{B}(u_1, u_1)$$

$$(\mathcal{A} - \lambda_0)u_3 = -\mu_1 u_2 - \mu_2 u_1 + 2\mathcal{B}(u_1, u_2).$$

$$u_1 = \sum_{1 \leq j \leq 2q} \xi_j, \quad \text{invariant under rotation } \mathbf{R}_{\pi/q}, \quad \text{spans } \ker(\mathcal{A} - \lambda_0)$$

$$\langle \mathcal{B}(u_1, u_1), u_1 \rangle_0 = 0, \quad \text{implies } \mu_1 = 0.$$

$$\mu_2 \langle u_1, u_1 \rangle_0 = \langle 2\mathcal{B}(u_1, u_2), u_1 \rangle_0 = -\langle (\mathcal{A} - \lambda_0)u_2, u_2 \rangle_0 > 0, \dots$$



Each step involves the pseudo-inverse $(\widetilde{\mathcal{A} - \lambda_0})^{-1}$, implying only **Gevrey series** for $\sum_{n \geq 1} \varepsilon^n u_n$ and $\sum_{n \geq 1} \varepsilon^n \mu_n$.

Idea: We decompose the system, as usual in bifurcation problems. The range equation contains the small divisor problem, we hope to use a parameter $\tilde{\mu}$ able to move the whole spectrum of the linearized operator as this is used by Berti, Bolle, Procesi (2010).

$$\begin{aligned}u &= u_\varepsilon + h(\varepsilon, \mu') + \varepsilon^4 \tilde{v}, \quad \mu = \mu_\varepsilon + \varepsilon^3 \mu', \\u_\varepsilon &= \varepsilon u_1 + \varepsilon^2 u_2 + \varepsilon^3 u_3 + \varepsilon^4 u_4, \quad \tilde{v} \in \{u_1\}^\perp \cap \mathcal{K}_s \\ \mu_\varepsilon &= \varepsilon^2 \mu_2 + \varepsilon^3 \mu_3, \quad \tilde{\mu} = \varepsilon^3 \mu'\end{aligned}$$

$$\text{Range equ. } \mathfrak{L}_{\varepsilon, \tilde{\mu}} \tilde{v} + g(\varepsilon, \tilde{\mu}) - \varepsilon^4 \mathbf{Q}_0 \mathcal{B}(\tilde{v}, \tilde{v}) = 0,$$

$$\text{Bifurc. equ. } \tilde{\mu} - \varepsilon^4 \mu_4 + \mathcal{O}[\varepsilon^3(\varepsilon + \|\tilde{v}\|)^2] = 0$$

$$\mathfrak{L}_{\varepsilon, \tilde{\mu}} := \mathbf{Q}_0(\mathcal{A} - \lambda_0) + \tilde{\mu} \mathbb{I} + \mathcal{R}_{\varepsilon, \tilde{\mu}}$$

$\mathfrak{L}_{\varepsilon, \tilde{\mu}}$ is analytic while $g(\varepsilon, \tilde{\mu})$ is only C^2 in $(\varepsilon, \tilde{\mu})$.

Main difficulty: Solve the Range equation with respect to \tilde{v}

For $|\varepsilon| \leq \varepsilon_0, |\tilde{\mu}| \leq \varepsilon_0, v \in \mathcal{K}_s, s \geq 0$

$$\begin{aligned} \|\mathcal{R}_{\varepsilon, \tilde{\mu}} v\|_s &\leq c_s \varepsilon \|v\|_s, \\ \|\partial_{\tilde{\mu}} \mathcal{R}_{\varepsilon, \tilde{\mu}} v\|_s + \|\partial_{\tilde{\mu}^2} \mathcal{R}_{\varepsilon, \tilde{\mu}} v\|_s + \|\partial_{\varepsilon \tilde{\mu}}^2 \mathcal{R}_{\varepsilon, \tilde{\mu}} v\|_s &\leq c_s \varepsilon^2 \|v\|_s, \\ \|\partial_{\varepsilon} \mathcal{R}_{\varepsilon, \tilde{\mu}} v\|_s + \|\partial_{\varepsilon^2} \mathcal{R}_{\varepsilon, \tilde{\mu}} v\|_s &\leq c_s \|v\|_s \end{aligned}$$

For $s_0 > d/2$, $\|2\varepsilon^4 \mathbf{Q}_0 \mathcal{B}(V, v)\|_s \leq c_s \varepsilon^4 (\|V\|_{s_0} \|v\|_s + \|V\|_s \|v\|_{s_0})$.

For $|\tilde{\mu}| \leq |\varepsilon|$

$$\begin{aligned} \|\mathbf{g}(\varepsilon, \tilde{\mu})\|_s &\leq c_s \varepsilon^2, \|\partial_{\varepsilon, \tilde{\mu}} \mathbf{g}(\varepsilon, \tilde{\mu})\|_s \leq c_s \varepsilon^2, \\ \|\partial_{\tilde{\mu}^2}^2 \mathbf{g}(\varepsilon, \tilde{\mu})\|_s &\leq c_s, \|\partial_{\varepsilon^2}^2 \mathbf{g}(\varepsilon, \tilde{\mu})\|_s \leq c_s \varepsilon^2, \|\partial_{\varepsilon \tilde{\mu}}^2 \mathbf{g}(\varepsilon, \tilde{\mu})\|_s \leq c_s \varepsilon^2. \end{aligned}$$

We need to invert

$$\mathfrak{L}_{\varepsilon, \tilde{\mu}, V} = \mathfrak{L}_{\varepsilon, \tilde{\mu}} - 2\varepsilon^4 \mathbf{Q}_0 \mathcal{B}(V, \cdot)$$

$$\lambda_0 - \lambda_0(|\mathbf{k}|^2) \geq 0, \quad \lambda_0 - \lambda_j(|\mathbf{k}|^2) > \delta_0 > 0, \quad j = 1, 2, \dots$$

For $\|\mathbf{k}\| - k_c > \delta_1$, $|\lambda_0 - \lambda_0(|\mathbf{k}|^2)| > \delta_0/2$.

Projection π_0 : suppresses $\mathbf{k} \in \Gamma$ such that $\|\mathbf{k}\| - k_c > \delta_1$

Inverting $\mathfrak{L}_{\varepsilon, \tilde{\mu}, V}$ equiv to invert $\mathfrak{L}'_{\varepsilon, \tilde{\mu}, V}$:

$$\mathfrak{L}'_{\varepsilon, \tilde{\mu}, V} = \pi_0 \mathbf{Q}_0 (\mathcal{A} - \lambda_0) \mathbf{Q}_0 \pi_0 + \tilde{\mu} + \mathfrak{B}_\varepsilon + \varepsilon^2 \tilde{\mu} \mathfrak{C}_{\varepsilon, \tilde{\mu}} + \mathfrak{R}_{\varepsilon, \tilde{\mu}, V}$$

\mathfrak{B}_ε , $\mathfrak{C}_{\varepsilon, \tilde{\mu}}$ and $\mathfrak{R}_{\varepsilon, \tilde{\mu}, V}$ depend analytically on their arguments.

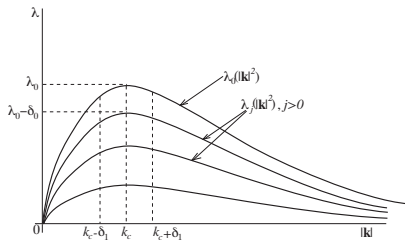
For $s \geq s_0 > d/2$

$$\|\mathfrak{B}_\varepsilon v\|_s \leq c\varepsilon \|v\|_s, \quad \|\mathfrak{C}_{\varepsilon, \tilde{\mu}} v\|_s + \|\partial_{\tilde{\mu}} \mathfrak{C}_{\varepsilon, \tilde{\mu}} v\|_s \leq c \|v\|_s$$

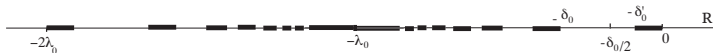
$$\|\mathfrak{R}_{\varepsilon, \tilde{\mu}, V} v\|_s \leq c\varepsilon^4 \{ \|V\|_{s_0} \|v\|_s + \|V\|_s \|v\|_{s_0} \},$$

$$\|\partial_{\tilde{\mu}} \mathfrak{R}_{\varepsilon, \tilde{\mu}, V} v\|_s \leq c\varepsilon^4 \{ \|V\|_{s_0} \|v\|_s + \|V\|_s \|v\|_{s_0} \}$$

Splitting by π_0 - continued



positive eigenvalues of \mathcal{A}



Spectrum of $\pi_0 \mathbf{Q}_0 (\mathcal{A} - \lambda_0) \mathbf{Q}_0 \pi_0$

Now, we truncate with the projection Π_N which cuts the \mathbf{k} such that $N_{\mathbf{k}} > N$.

This defines the space $E_N = \Pi_N \pi_0 \mathbf{Q}_0 \mathcal{K}_s$ still ∞ - dim, but with an isolated finite group of "small eigenvalues" perturbing $\lambda_0(|\mathbf{k}|^2) - \lambda_0, N_{\mathbf{k}} \leq N$.

$$\lambda_0 - \lambda_0(|\mathbf{k}|^2) \sim c(|\mathbf{k}|^2 - k_c^2)^2 \geq \frac{c}{(1 + N^2)^{2l_0}}$$

Lemma

Let V satisfy $\|V\|_{s_0} \leq 1$, and assume $(\varepsilon, \tilde{\mu}) \in [0, \varepsilon_0] \times [-\varepsilon, \varepsilon]$.
 Then for $N \leq M_\varepsilon$, where $M_\varepsilon = \left\lceil \frac{c_2}{\varepsilon^{1/4l_0}} \right\rceil$, we have the following estimate for $s_0 > d/2$

$$\begin{aligned} \|(\Pi_N \mathcal{L}'_{\varepsilon, \tilde{\mu}, V} \Pi_N)^{-1} v\|_s &\leq 2c(1 + N^2)^{2l_0} \{\|v\|_s + \|V\|_s \|v\|_{s_0}\}, \\ \|(\Pi_N \mathcal{L}_{\varepsilon, \tilde{\mu}, V} \Pi_N)^{-1} v\|_s &\leq 2cc'(1 + N^2)^{2l_0} \{\|v\|_s + \|V\|_s \|v\|_{s_0}\}. \end{aligned}$$

Hint: Classical perturbation theory, based on the smallness of $\varepsilon(1 + N^2)^{2l_0}$ for $N \leq M_\varepsilon$

For $M > 0$, $s_0 > d/2$ define

$$\mathcal{U}_M^{(N)} := \{u \in C^2([0, \varepsilon_1] \times [-\varepsilon, \varepsilon], E_N); u(0, \tilde{\mu}) = 0, \\ \|u\|_{s_0} \leq 1, \|\partial_{\varepsilon, \tilde{\mu}}^j u\|_{s_0} \leq M, j = 1, 2\}$$

For $V \in \mathcal{U}_M^{(N)}$ let us denote $\mathfrak{L}_{\varepsilon, \tilde{\mu}}^{(N, V)} =: \Pi_N \mathfrak{L}'_{\varepsilon, \tilde{\mu}, V(\varepsilon, \tilde{\mu})} \Pi_N$,

$$\mathfrak{L}_{\varepsilon, \tilde{\mu}}^{(N, V)} = (\widetilde{\mathcal{A} - \lambda_0})_N + \tilde{\mu} \mathbb{I}d + \mathfrak{B}'_{\varepsilon}^{(N)} + \varepsilon^2 \mathfrak{C}'_{\varepsilon, \tilde{\mu}}^{(N)},$$

$$(\widetilde{\mathcal{A} - \lambda_0})_N =: \Pi_N \pi_0 \mathbf{Q}_0 (\mathcal{A} - \lambda_0) \mathbf{Q}_0 \pi_0 \Pi_N$$

Now define **the selfadjoint operator**

$$\mathfrak{L}_{\varepsilon, \tilde{\mu}}^{(N, V)} \mathfrak{L}_{\varepsilon, \tilde{\mu}}^{(N, V)*} = \tilde{\mu}^2 \mathbb{I}d + \widetilde{\mathfrak{B}}_{\varepsilon}^{(N)} + \widetilde{\mathfrak{C}}_{\varepsilon, \tilde{\mu}}^{(N)}$$

$$\mathfrak{L}_{\varepsilon, \tilde{\mu}}^{(N, V)} \mathfrak{L}_{\varepsilon, \tilde{\mu}}^{(N, V)*} = \tilde{\mu}^2 \mathbb{I}d + \widetilde{\mathfrak{B}}_{\varepsilon}^{(N)} + \widetilde{\mathfrak{C}}_{\varepsilon, \tilde{\mu}}^{(N)}$$

$$\widetilde{\mathfrak{B}}_{\varepsilon}^{(N)} = \widetilde{(\mathcal{A} - \lambda_0)_N^2} + \mathfrak{B}'_{\varepsilon}{}^{(N)} \widetilde{(\mathcal{A} - \lambda_0)_N} + \widetilde{(\mathcal{A} - \lambda_0)_N} \mathfrak{B}'_{\varepsilon}{}^{(N)*} + \mathfrak{B}'_{\varepsilon}{}^{(N)} \mathfrak{B}'_{\varepsilon}{}^{(N)*},$$

$$\begin{aligned} \widetilde{\mathfrak{C}}_{\varepsilon, \tilde{\mu}}^{(N)} &= \tilde{\mu} [2 \widetilde{(\mathcal{A} - \lambda_0)_N} + \mathfrak{B}'_{\varepsilon}{}^{(N)} + \mathfrak{B}'_{\varepsilon}{}^{(N)*}] + \\ &+ \varepsilon^2 [(\widetilde{(\mathcal{A} - \lambda_0)_N} + \mathfrak{B}'_{\varepsilon}{}^{(N)} + \tilde{\mu}) \mathfrak{C}'_{\varepsilon, \tilde{\mu}}{}^{(N)*} + \\ &+ \varepsilon^2 \mathfrak{C}'_{\varepsilon, \tilde{\mu}}{}^{(N)} [(\widetilde{(\mathcal{A} - \lambda_0)_N} + \mathfrak{B}'_{\varepsilon}{}^{(N)*} + \tilde{\mu}) + \varepsilon^4 \mathfrak{C}'_{\varepsilon, \tilde{\mu}}{}^{(N)} \mathfrak{C}'_{\varepsilon, \tilde{\mu}}{}^{(N)*}, \end{aligned}$$

Good set of $\tilde{\mu}$ (continued)

For $V \in \mathcal{U}_M^{(N)}$, then define the "good" set of $\tilde{\mu}$:

$$G_{\varepsilon, \gamma}^{(N)}(V) =: \{ \tilde{\mu} \in [-\varepsilon, \varepsilon]; \text{ for all } v \in E_N, \\ \| \Pi_N \mathfrak{L}'_{\varepsilon, \tilde{\mu}, V} \Pi_N^{-1} v \|_{s_0} \leq \frac{N^\tau}{\gamma} \| v \|_{s_0} \}.$$

Consequence: if $\tilde{\mu} \in G_{\varepsilon, \gamma}^{(N)}(V)$, then $\mathfrak{L}_{\varepsilon, \tilde{\mu}}^{(N, V)} \mathfrak{L}_{\varepsilon, \tilde{\mu}}^{(N, V)*}$ has all its eigenvalues $\geq (\frac{\gamma}{N^\tau})^2$ in E_N .

Notice that $\dim(E_N) = \mathcal{N} \leq bN^d$.

Bad set of $\tilde{\mu}$

For $V \in \mathcal{U}_M^{(N)}$ define the "bad" set of $\tilde{\mu}$

$$B_{\varepsilon, \gamma}^{(N)}(V) =: \{ \tilde{\mu} \in [-\varepsilon, \varepsilon]; \text{ there exists at least one eigenvalue } \sigma_j \text{ of } \mathfrak{L}_{\varepsilon, \tilde{\mu}}^{(N, V)} \mathfrak{L}_{\varepsilon, \tilde{\mu}}^{(N, V)*}, \text{ such that } 0 \leq \sigma_j \leq (\frac{\gamma}{N^\tau})^2 \}.$$

the eigenvalues of $\mathfrak{L}_{\varepsilon, \tilde{\mu}}^{(N, V)} \mathfrak{L}_{\varepsilon, \tilde{\mu}}^{(N, V)*}$ take the form

$$\sigma_j(\varepsilon, \tilde{\mu}) = \tilde{\mu}^2 + f_j(\varepsilon, \tilde{\mu}), \quad f_j \text{ is } C^2 \text{ in } \tilde{\mu}, \text{ and}$$

$$|f_j(\varepsilon_2, \tilde{\mu}_2) - f_j(\varepsilon_1, \tilde{\mu}_1)| \leq c(\delta'_0 + \varepsilon)(|\varepsilon_2 - \varepsilon_1| + |\tilde{\mu}_2 - \tilde{\mu}_1|).$$

Assumption For $\tilde{\mu} \in [-\varepsilon, \varepsilon]$, there exists $0 < k < 2$ with

$$|\partial_{\tilde{\mu}} f_j(\varepsilon, \tilde{\mu}_2) - \partial_{\tilde{\mu}} f_j(\varepsilon, \tilde{\mu}_1)| \leq k |\tilde{\mu}_2 - \tilde{\mu}_1|.$$

Lemma

Assume that $N > M_\varepsilon$, $d/2 < s_0$, $(\varepsilon, \tilde{\mu}) \in (0, \varepsilon_1] \times [-\varepsilon, \varepsilon]$, and $V \in \mathcal{U}_M^{(N)}$. Then there exists $C > 0$, such that the measure of $B_{\varepsilon, \gamma}^{(N)}(V)$ is bounded by $C\gamma/N^{\tau-d}$

Bad set of $\tilde{\mu}$ (proof of the bound)

Notice that in the Assumption above, **we accept to loose some regularity for the second derivative of σ_j with respect to $\tilde{\mu}$** . Up to now we have no mean to control this loss.

For a "bad" $\tilde{\mu}$, there exists j such that

$$0 \leq \tilde{\mu}^2 + f_j(\varepsilon, \tilde{\mu}) < \eta^2, \eta = \gamma/(N^\tau)$$

$$\text{define } \phi_\varepsilon(\tilde{\mu}) =: \tilde{\mu}^2 + f_j(\varepsilon, \tilde{\mu})$$

$$\partial_{\tilde{\mu}} \phi_\varepsilon(\tilde{\mu}) = 2\tilde{\mu} + \partial_{\tilde{\mu}} f_j(\varepsilon, \tilde{\mu})$$

increasing function of $\tilde{\mu}$, cancelling at a unique $\tilde{\mu} = \tilde{\mu}_m$.

Since $\phi_\varepsilon(\tilde{\mu}^\pm) = \eta^2$, then the convexity of ϕ_ε implies

$$\tilde{\mu}^+ - \tilde{\mu}^- \leq \frac{2\eta}{\sqrt{(1 - k/2)}}$$

Summing up for all eigenvalues, the measure of the set of bad $\tilde{\mu}$ is bounded by

$$\frac{2b\gamma}{\sqrt{(1 - k/2)N^{\tau-d}}}$$

Lemma

Let $d = 2(l_0 + 1)$ be the dimension of the \mathbb{Q} -vector space spanned by the wave vectors $k_j, j = 1, \dots, 2q$, and $\tau > d + 2 + 24l_0$. Let N be ≥ 1 . Assume moreover that $0 < \gamma \leq \tilde{\gamma} = \frac{c'}{c2^{2l_0+1}}$, $(\varepsilon, \tilde{\mu}, V) \in [0, \varepsilon_1] \times [-\varepsilon, \varepsilon] \times \mathcal{U}_M^{(N)}$ with $\tilde{\mu} \in G_{\varepsilon, \gamma}^{(N)}(V)$, ε_1 small enough. For $s_0 > \frac{d}{2}$, there exists $c' > 0$ independent of N and γ , such that for any $v \in E_N$, we have

$$\|(\Pi_N \mathcal{L}'_{\varepsilon, \tilde{\mu}, V(\varepsilon, \tilde{\mu})} \Pi_N)^{-1} v\|_{s_0} \leq c' \frac{N^\tau}{\gamma} \|v\|_{s_0}$$

and the same estimate holds for $(\Pi_N \mathcal{L}_{\varepsilon, \tilde{\mu}, V(\varepsilon, \tilde{\mu})} \Pi_N)^{-1}$.

We need an estimate in all \mathcal{K}_s , with an exponent on N independent of s . We may proceed as for the Swift-Hohenberg PDE, in adapting Bourgain 1995, Craig 2000, bert-Bolle 2010.

Separation property of the singular set

Singular set in \mathbb{Z}^d :

$S(N) := \{\mathbf{z} \in \Gamma(N); (\lambda_0 - \lambda_0(|\mathbf{k}(\mathbf{z})|^2) < \rho, \mathbf{k}(\mathbf{z}) \in \Gamma(N)\}$ with

$$\mathbf{k}(\mathbf{z}) = \vartheta^{-1} \sum_{s=1}^d z_s \mathbf{k}_s^*, \quad \mathbf{z} = (z_1, \dots, z_d) \in \mathbb{Z}^d$$

$$\Gamma(N) := \{\mathbf{z} \in \mathbb{Z}^d; 0 \leq |\mathbf{z}| \leq N, \mathbf{k}(\mathbf{z}) \in \Gamma \setminus \{\mathbf{k}_j, j = 1, \dots, 2q\}\}.$$

Useful lemma (uses Bourgain 1995, Craig 2000, Berti-Bolle 2010)

There exists $\rho_0 > 0$ independent of N such that if $\rho \in]0, \rho_0]$ then

$S(N) = \bigcup_{\alpha \in \mathcal{A}} \Omega_\alpha$ is a **union of disjoint clusters** Ω_α satisfying :

- (H1), for all $\alpha \in \mathcal{A}$, $M_\alpha \leq 2m_\alpha$ where $M_\alpha = \max_{\mathbf{z} \in \Omega_\alpha} |\mathbf{z}|$ and $m_\alpha = \min_{\mathbf{z} \in \Omega_\alpha} |\mathbf{z}|$;
- (H2), there exists $\delta = \delta(d) \in]0, 1[$ independent of N such that if $\alpha, \beta \in \mathcal{A}, \alpha \neq \beta$ then

$$\text{dist}(\Omega_\alpha, \Omega_\beta) := \min_{\mathbf{z} \in \Omega_\alpha, \mathbf{z}' \in \Omega_\beta} |\mathbf{z} - \mathbf{z}'| \geq \frac{(M_\alpha + M_\beta)^\delta}{2}$$

Estimate of $(\Pi_N \mathcal{L}'_{\varepsilon, \tilde{\mu}, V} \Pi_N)^{-1}$ in $\Pi_N \pi_0 \mathbf{Q}_0 \mathcal{K}_s$ for all $s \in [s_0, \bar{s}]$

Lemma

Let $d = 2(l_0 + 1)$ be the dimension of the \mathbb{Q} -vector space spanned by the wave vectors $k_j, j = 1, \dots, 2q$, and $\tau > d + 2 + 24l_0$ as in previous Lemma. Assume moreover that $0 < \gamma \leq \tilde{\gamma} = \frac{c'}{c^{2l_0+1}}$, and $(\varepsilon, \tilde{\mu}, V) \in [0, \varepsilon_1] \times [-\varepsilon, \varepsilon] \times \mathcal{U}_M^{(N)}$, $\tilde{\mu} \in \mathcal{G}_{\varepsilon, \gamma}^{(N)}(V) = \cap_{K \leq N} \mathcal{G}_{\varepsilon, \gamma}^{(K)}(V)$, ε_1 small enough. There exists $s_0(d, \delta, \tau) > \frac{d}{2}$ where δ is the number introduced in separation property (H2), and let $\bar{s} > s_0$. There exists $m(d, \delta, \tau)$ such that for all $s \in [s_0, \bar{s}]$ there exists $K(s) > 0$ such that for any $h \in \Pi_N \pi_0 \mathbf{Q}_0 \mathcal{K}_{0,s}$, we have

$$\|(\Pi_N \mathcal{L}'_{\varepsilon, \tilde{\mu}, V(\varepsilon, \tilde{\mu})} \Pi_N)^{-1} h\|_s \leq K(s) \frac{N^m}{\gamma} (\|h\|_s + \|V(\varepsilon, \tilde{\mu})\|_s \|h\|_{s_0}),$$

and the same estimate holds for $(\Pi_N \mathcal{L}_{\varepsilon, \tilde{\mu}, V(\varepsilon, \tilde{\mu})} \Pi_N)^{-1}$.

Main ingredient for applying the Nash-Moser iteration process.

Resolution of the Range equation

Uses Nash-Moser method, following Berti-Bolle-Procesi 2010 with a complement for having a solution $v \in \mathcal{K}_{s_0}$ of the range equation, which is C^2 .

$$\mathcal{F}(\epsilon, \tilde{\mu}, v) =: \mathfrak{L}_{\epsilon, \tilde{\mu}} v + g(\epsilon, \tilde{\mu}) - \epsilon^4 \mathbf{Q}_0 \mathcal{B}(v, v) = 0,$$

Theorem

Let s_0 and $\tilde{\gamma}$ be as above. Then for all $0 < \gamma < \tilde{\gamma}$ there exist $\epsilon_2(\gamma) \in [0, \epsilon_0]$ and a C^2 -map $V : (0, \epsilon_2(\gamma)) \times [-\epsilon, \epsilon] \rightarrow \Pi_N \pi_0 \mathbf{Q}_0 \mathcal{K}_{s_0}$, such that $V(0, \tilde{\mu}) = 0$, $\|V\|_{s_0} \leq 1$, $\|\partial_{\tilde{\mu}} V\|_{s_0} \leq M$, $\|\partial_{\tilde{\mu}}^2 V\|_{s_0} \leq M$, and if $\epsilon \in (0, \epsilon_2(\gamma))$, $\tilde{\mu} \in ([-\epsilon, \epsilon] \setminus C_{\epsilon, \gamma})$, the function $V(\epsilon, \tilde{\mu})$ is solution of the range equation $\mathcal{F}(\epsilon, \tilde{\mu}, v) = 0$. Here $C_{\epsilon, \gamma}$ is a subset of $[-\epsilon, \epsilon]$ which is Hölder continuous in ϵ , and has Lebesgue-measure less than $C\gamma\epsilon^6$ for some constant $C > 0$ independent of ϵ and γ .

Hint: Proof adapted from Berti, Bolle, Procesi 2010.

$$\begin{aligned}\tilde{\mu} - \epsilon^4 \mu_4 + \mathcal{O}[\epsilon^3(\epsilon + \|\tilde{v}\|)^2] &= 0 \\ \tilde{v} &= V(\epsilon, \tilde{\mu}) - h(\epsilon, \tilde{\mu})\end{aligned}$$

We solve this bifurcation equation with respect to $\tilde{\mu}$:

$$\tilde{\mu} = \epsilon^4 \mu_4 + \epsilon^5 \tilde{h}(\epsilon), \text{ "curve" (H) , } \tilde{h} \in C^1$$

We need to satisfy that $(\epsilon, \tilde{\mu})$ lies in the good set, defined in the Theorem above.

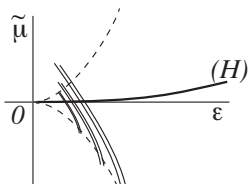
Definition

For N and V fixed, a set of "bad intervals" is defined by

$$BS_N(V) = \{(\varepsilon, \tilde{\mu}) \in [0, \varepsilon_2] \times [-\varepsilon^3, \varepsilon^3]; \tilde{\mu} \in I_\varepsilon^{(N)}\},$$

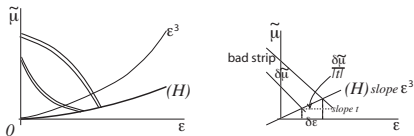
where $I_\varepsilon^{(N)}$ is one of the intervals $(\tilde{\mu}_j^-(\varepsilon), \tilde{\mu}_j^+(\varepsilon))$, or with one of the bounds replaced by $\pm\varepsilon^3$, as defined above.

$BS_N(V)$ is a union of thin Hölder strips in the plane $(\varepsilon, \tilde{\mu})$. For the proof of the range theorem, we choose $\tilde{\mu}$ outside of $\bigcup_{n \in \mathbb{N}} BS_{N_n}(V_{n-1})$ where $N_n = [N_0(\gamma)]^{2^n}$, and V_n are the successive points in the Newton iteration process.



Resolution of the Bifurcation equation- Transversality condition

We need a "transversality condition" to obtain a bound for the bad set of ε corresponding to the intersection of $\cup_{\varepsilon \in (0, \varepsilon_2)} C_{\varepsilon, \gamma}$ with (H) :
 The slopes $t(\varepsilon)$ of the curves $\tilde{\mu}_j^\pm(\varepsilon)$, are such that there exists $c > 0$ independent of (N, ε) , with $c\varepsilon^2 < |t(\varepsilon)|$.



Sketch of the "bad set" in the plane $(\varepsilon, \tilde{\mu})$ and its intersection by the "line" (H) given by the bifurcation equation.

The drawing on the right side explains the bound for the measure of $\delta\varepsilon \leq \delta\tilde{\mu}/c\varepsilon^2 \leq C\varepsilon^4$.

Notice that, as for the Swift-Hohenberg equation, we can weaken this transversality condition, so that the true hypothesis is that the curves $\tilde{\mu}_j^\pm(\varepsilon)$ are not flat.

Theorem

Let $q \geq 4$ be an integer. Assume that Hypothesis $\lambda_0'' \neq 0$ holds and that transversality condition above is verified. Moreover assume the convexity condition above, on small eigenvalues σ_j . Then, there exists $s_0 > d/2$, $\varepsilon_0 > 0$, such that, for any $s \geq s_0$, there exists a 1-dimensional set $\bar{\Lambda}_\varepsilon$ centered on μ_4 , with the following property : for any $|\varepsilon| < \varepsilon_0$, belonging to a set of asymptotically full measure as $\varepsilon \rightarrow 0$ there exist $\bar{\mu}_\varepsilon \in \bar{\Lambda}_\varepsilon$, such that the steady Bénard - Rayleigh system admits a quasipattern solution (u, λ) , C^1 in ε , $u \in \mathcal{K}_s$, $\lambda = \lambda_0 - \mu_2 \varepsilon^2 - \mu_3 \varepsilon^3 - \varepsilon^4 \bar{\mu}_\varepsilon$ invariant under rotations of angle π/q of the form

$$u = \varepsilon u_1 + \varepsilon^2 u_2 + \varepsilon^3 u_3 + \varepsilon^4 u_4 + \mathcal{O}(\varepsilon^5),$$

where $\mu_2 > 0$, and coefficients μ_2, μ_4, u_j occurring in formulae above, are the ones defined in the truncated asymptotic expansion of the solution.

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