

# Lectures Références

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# Numerical Methods for PDE: Finite Differences and Finite Volumes

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2009

- 1 Finite Difference(FD) and Finite volume(FV) : Overview
- 2 Modelization and Simplified models of PDE.
- 3 Scalar Advection-Diffusion Eqation.
- 4 Approximation of a Scalar 1D ODE.
- 5 FD for 1D scalar poisson equation (elliptic).
- 6 FD for 1D scalar difusion equation (parabolic).
- 7 FD for 1D scalar advection-diffusion equation.
- 8 Scalar Nonlinear Conservation law : 1D (hyperbolic).
- 9 FV for scalar nonlinear Conservation law : 1D
- 10 Multi-Dimensional extensions

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# Scalar Conservation Laws

General Integral form :

$$\frac{d}{dt} \int_{x_l}^{x_r} \mathbf{u} dx = - (f(\mathbf{u}_r) - f(\mathbf{u}_l))$$

$\mathbf{u}$  is the unknown conserved quantity (scalar) and  $f$  is the flux function(given).

- Conservative EDP form :  $\partial_t \mathbf{u} + \partial_x [f(\mathbf{u})] = 0$
- Primitive form (Non conservative) : **regular** solutions of the conservation law are also solutions of the following equation :

$$\partial_t \mathbf{u} + f'(\mathbf{u}) \partial_x \mathbf{u} = 0$$

- Viscous profile. Question : can we find a solution  $\mathbf{u}^\varepsilon$  of the following system

$$\partial_t \mathbf{u}^\varepsilon + \partial_x f(\mathbf{u}^\varepsilon) = \varepsilon \partial_x^2 \mathbf{u}^\varepsilon$$

such that  $\mathbf{u} = \lim_{\varepsilon \rightarrow 0} \mathbf{u}^\varepsilon$ . This will help for numerical design.

# Examples

- Linear Advection equation :  $u$  the density and

$$f(u) = cu$$

where  $c$  is constant.

- Inviscid Burgers equation :  $u$  is the velocity and

$$f(u) = \frac{1}{2}u^2.$$

- Traffic flow :  $u$  is the density of cars and

$$f(u) = \alpha \left( u - \frac{u^2}{\beta} \right).$$

- Buckley-Leverett equation :  $u$  is the saturation of water and

$$f(u) = \frac{u^2}{u^2 + \alpha(1-u)^2}.$$

# Characteristic curves

The curves, parametrized as  $\mathbf{x} \equiv \mathbf{x}(t)$  and such that

$$\frac{d\mathbf{x}(t)}{dt} = f'(\mathbf{u}(\mathbf{x}(t), t))$$

are called Characteristics.

The solution  $\mathbf{u}$  is constant along the characteristics

$$\frac{d\mathbf{u}(\mathbf{x}(t), t)}{dt} = \frac{\partial \mathbf{u}}{\partial t} + \frac{d\mathbf{x}(t)}{dt} \frac{\partial \mathbf{u}}{\partial \mathbf{x}} = \frac{\partial \mathbf{u}}{\partial t} + f'(\mathbf{u}(\mathbf{x}(t), t)) \frac{\partial \mathbf{u}}{\partial \mathbf{x}} = 0$$

Therefore, Characteristics are straight lines :

$\mathbf{u}(\mathbf{x}(t), t) = \mathbf{u}(\mathbf{x}(t_0), t_0) = \mathbf{u}_0(\mathbf{x}(t_0))$  and

$$\mathbf{x}(t) = \mathbf{x}(t_0) + (t - t_0) f'(\mathbf{u}(\mathbf{x}(t_0)))$$

# Existence of smooth Solutions

When the function  $\xi \rightarrow f'(\mathbf{u}_0(\xi))$  is increasing and  $\mathbf{u}_0(\xi)$  is regular and bounded, There exists a unique solution of the Cauchy problem. This solution is defined as

$$\mathbf{u}(x, t) = \mathbf{u}_0(\xi(x, t)) \quad \text{where} \quad \xi(x, t) = x - t f'(\mathbf{u}_0(\xi(x, t)))$$

Existence is OK : by using Characteristic curves.

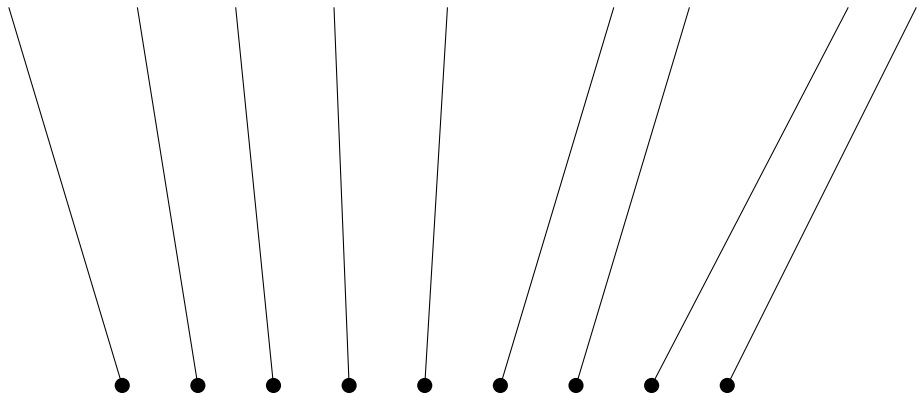
Uniqueness : Let  $\mathcal{F}(\xi) = \xi - x + t f'(\mathbf{u}_0(\xi))$ .

$$\frac{d}{d\xi} f'(\mathbf{u}_0(\xi(t))) > 0 \implies \mathcal{F}' = 1 + t \frac{d}{d\xi} f'(\mathbf{u}_0(\xi(t))) > 0$$

$$\mathcal{F}(-\infty) < 0, \quad \mathcal{F}(\infty) > 0, \implies \exists! \xi(x, t) \quad \mathcal{F}(\xi(x, t)) = 0$$

Therefore, we have the uniqueness of  $\mathbf{u}(x, t) = \mathbf{u}_0(\xi(x, t))$ .

# Existence of smooth Solutions



When the initial solution is smooth and Characteristic curves do not cross :  $f'(u_0(\xi))$  is an increasing function.



# Discontinuous Solutions : Shock formation

**Schock formation : Let  $\mathbf{u}_0$  and  $f$  regular functions.**

When  $\frac{d}{d\xi} f'(\mathbf{u}_0(\xi(t))) < 0$ , it exists a time  $t = t^* = -\frac{1}{\frac{d}{d\xi} f'(\mathbf{u}_0(\xi(t)))}$  such that  $Y(t) = \partial_x \mathbf{u}(x(t), t)$  is no longer bounded.

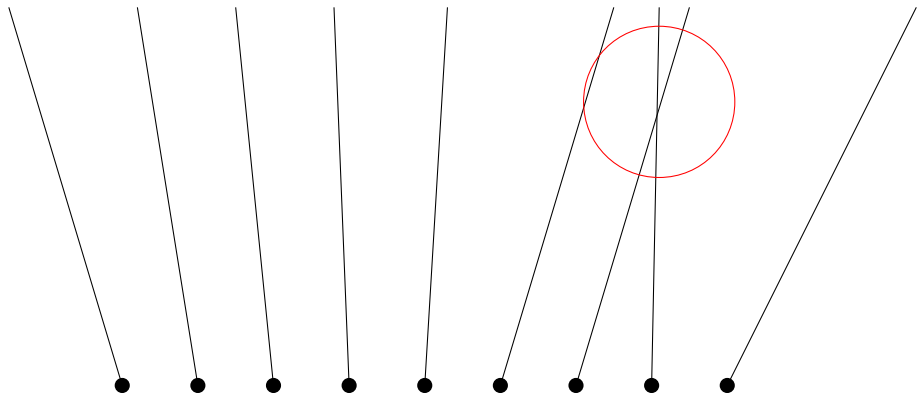
$$\begin{aligned} \frac{dY}{dt} &= \partial_x \partial_t \mathbf{u}(x(t), t) + \frac{dx(t)}{dt} \partial_x (\partial_x \mathbf{u}(x(t), t)) \\ &= -\partial_x (f'(\mathbf{u}) \partial_x \mathbf{u}) + f'(\mathbf{u}) \partial_x (\partial_x \mathbf{u}(x(t), t)) \\ &= -f''(\mathbf{u}(x(t), t)) Y^2 = -f''(\mathbf{u}_0(\xi)) Y^2 \end{aligned}$$

$Y(t)$  satisfied a Riccati ODE, and the solution is :

$$Y(t) = \frac{Y(0)}{1 + tY(0)f''(\mathbf{u}_0(\xi))} = \frac{Y(0)}{1 + t\frac{d}{d\xi} f'(\mathbf{u}_0(\xi(t)))}$$

Therefore, when  $\frac{d}{d\xi} f'(\mathbf{u}_0(\xi(t))) < 0$  the solution cannot be regular after the time  $t = t^* = -\frac{1}{\frac{d}{d\xi} f'(\mathbf{u}_0(\xi(t)))} > 0$

# Discontinuous Solutions : Shock formation



The initial solution can be smooth but characteristic curves cross after a given time :  $f'(u_0(\xi))$  is a non increasing function.

# Weak Solutions : are often non unique

Cauchy Problem :

$$\begin{cases} \partial_t \mathbf{u} + \partial_x (f(\mathbf{u})) = 0 & \forall x \in \mathbb{R}, \quad \forall t > 0 \\ \mathbf{u}(x, t = 0) = \mathbf{u}_0(x) & \forall x \in \mathbb{R} \end{cases}$$

## Definition (Weak Solution)

$\mathbf{u}$  is a weak solution of the Cauchy problem if the following statements are satisfied :

- $\mathbf{u} \in L^1_{loc}(\mathbb{R} \times ]0, T[)$ .
- $\forall \phi(x, t)$  regular function with compact support in  $\mathbb{R} \times \mathbb{R}$  ( $C^1_0(\mathbb{R} \times \mathbb{R})$ ) :

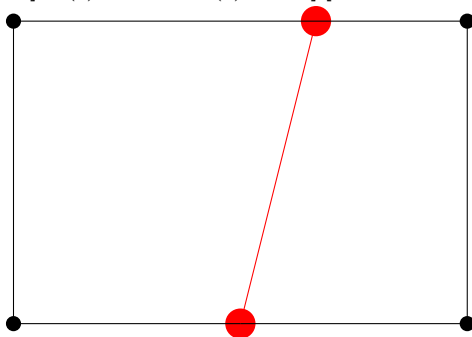
$$\int_0^\infty dt \int_{-\infty}^\infty dx (\mathbf{u} \partial_t \phi + f(\mathbf{u}) \partial_x \phi) = - \int_{-\infty}^\infty dx u_0(x) \phi(x, 0)$$

# Characterization of Discontinuities

Let consider the constant states  $\mathbf{u}_l$  and  $\mathbf{u}_r$  separated by a curve of discontinuity  $x = X(t)$

$$\mathcal{Q} = [X(t) - \delta x, X(t) + \delta x] \times [t, t + \delta t]$$

with  $X(t+s) \in ]X(t) - \delta x, X(t) + \delta x[$  pour  $0 \leq s \leq \delta t$



# Rankine-Hugoniot Jump condition

Integral form of the conservation law integrated over  $(t, t + \delta t)$  gives

$$\int_{X(t)-\delta x}^{X(t+\delta t)} \mathbf{u}_l dx + \int_{X(t+\delta t)}^{X(t)+\delta x} \mathbf{u}_r dx = \int_{X(t)-\delta x}^{X(t)} \mathbf{u}_l dx + \int_{X(t)}^{X(t)+\delta x} \mathbf{u}_r dx - \int_t^{t+\delta t} f(\mathbf{u}_r) dt + \int_t^{t+\delta t} f(\mathbf{u}_l) dt$$

$$\frac{X(t + \delta t) - X(t)}{\delta t} (\mathbf{u}_l - \mathbf{u}_r) = f(\mathbf{u}_l) - f(\mathbf{u}_r) \quad \forall \delta t > 0$$

$$\delta t \rightarrow 0 \implies f(\mathbf{u}_r) - f(\mathbf{u}_l) = s(t) (\mathbf{u}_r - \mathbf{u}_l) \quad (1)$$

where  $s(t) = X'(t)$  is the shock speed.

**Jump condition : define the wave speed  $s$**

$$f(\mathbf{u}_r) - f(\mathbf{u}_l) = s (\mathbf{u}_r - \mathbf{u}_l)$$

# Characterization of discontinuities : contact and shocks

$$m = f(\mathbf{u}_r) - s\mathbf{u}_r = f(\mathbf{u}_l) - s\mathbf{u}_l \iff [m] = 0$$

where  $m$  is the flux crossing the discontinuity.

- When  $m = 0$  : the associated wave is linearly degenerated and is named Contact Discontinuity (CD)
- When  $m \neq 0$  : the associated wave is linearly genuinely nonlinear and is named Shock.

For linear advection  $f(\mathbf{u}) = a\mathbf{u} \implies s = a$  and  $m = 0$ . Waves are Contact discontinuities.

Burgers Equation :  $f(\mathbf{u}) = \frac{\mathbf{u}^2}{2} \implies s = \frac{\mathbf{u}_r + \mathbf{u}_l}{2}$

$$m = \frac{\mathbf{u}_r^2}{2} - \mathbf{u}_r \frac{\mathbf{u}_r + \mathbf{u}_l}{2} = -\frac{\mathbf{u}_r \mathbf{u}_l}{2}$$

# Weak solutions piecewise regular

A piecewise function  $\mathbf{u}(\mathbf{x}, t)$  is a weak solution of the Cauchy problem if and only if :

- $\mathbf{u}(\mathbf{x}, t = 0) = \mathbf{u}_0(\mathbf{x})$
- Where  $\mathbf{u}$  is regular it satisfy the partial differential equation.
- across discontinuities the Rankine-Hugoniot relation is satisfied.

# Manipulating equations has impact on weak solutions

Let us consider functions  $\phi(X)$  and  $\Psi(X)$  such that

$$\phi'(\Psi(\mathbf{u})) = f'(\mathbf{u}) \quad \text{with} \quad \Psi'(\mathbf{u}) \neq 0$$

where  $\mathbf{u}$  is a solution (regular) of

$$\partial_t \mathbf{u} + \partial_x f(\mathbf{u}) = 0$$

Then  $v = \Psi(\mathbf{u})$  satisfy the following conservation law.

$$\partial_t v + \partial_x \phi(v) = 0$$

Jump condition for this new conservation law is not equivalent to the one associated to  $\partial_t \mathbf{u} + \partial_x f(\mathbf{u}) = 0$ . Indeed, in general,

$$s_u = \frac{f(\mathbf{u}_r) - f(\mathbf{u}_l)}{\mathbf{u}_r - \mathbf{u}_l} \neq s_v = \frac{\phi(\Psi(\mathbf{u}_r)) - \phi(\Psi(\mathbf{u}_l))}{\Psi(\mathbf{u}_r) - \Psi(\mathbf{u}_l)}$$



## Example : For the Burgers Equation

$$\partial_t \mathbf{u} + \partial_x \left( \frac{\mathbf{u}^2}{2} \right) = 0 \implies f(\mathbf{u}) = \frac{\mathbf{u}^2}{2}$$

Then let us define  $\Psi(X)$  and  $\phi(X)$  as

$$\Psi(X) = X^2 \quad \text{and} \quad \phi(X) = \frac{2}{3} X^{\frac{3}{2}}$$

Then we find that  $\phi'(X) = \sqrt{X}$  and  $\Psi'(X) = 2X$  and finally, for positive  $\mathbf{u} > 0$ ,

$$\begin{aligned} \phi'(\Psi(\mathbf{u})) &= \mathbf{u} \quad \text{and} \quad \Psi'(\mathbf{u}) \neq 0 \\ \implies \partial_t (\mathbf{u}^2) + \partial_x \left( \frac{2}{3} \mathbf{u}^3 \right) &= 0 \end{aligned}$$

But

$$\frac{\mathbf{u}_r + \mathbf{u}_l}{2} \neq \frac{2}{3} \frac{(\mathbf{u}_l^2 + \mathbf{u}_l \mathbf{u}_r + \mathbf{u}_r^2)}{\mathbf{u}_l + \mathbf{u}_r}$$

# Weak solutions : Non-uniqueness

## Definition of the Riemann Problem

A Riemann problem is a Cauchy problem with the following initial data

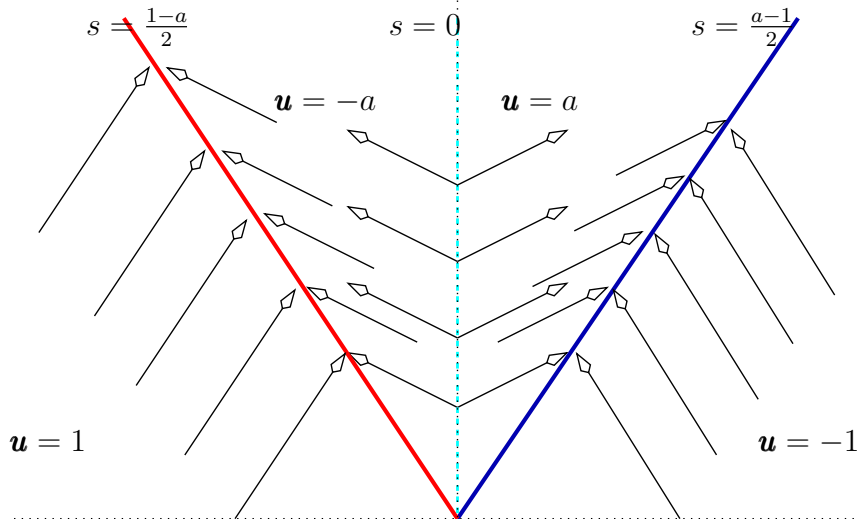
$$\mathbf{u}_0(x) = \begin{cases} \mathbf{u}_l & \text{when } x < 0 \\ \mathbf{u}_r & \text{when } x \geq 0 \end{cases}$$

where  $\mathbf{u}_l$  and  $\mathbf{u}_r$  are given constants.

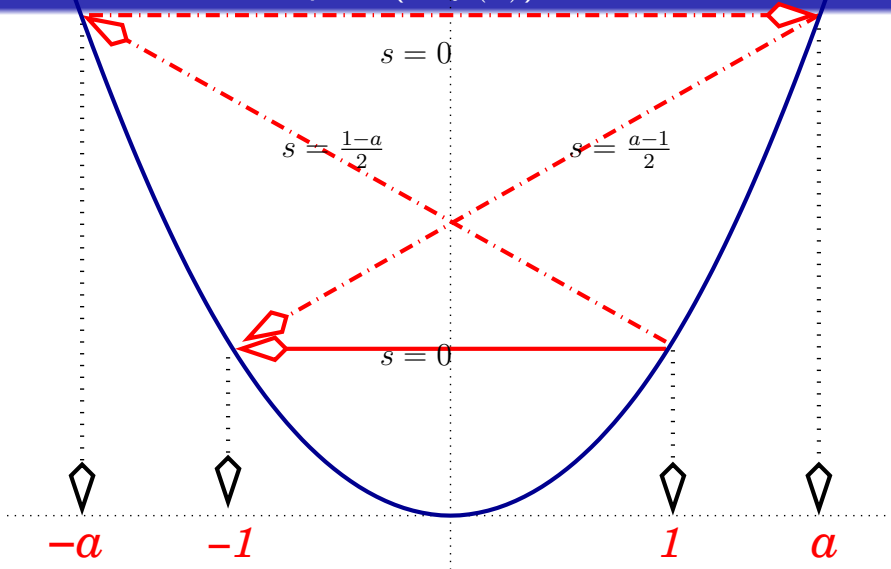
Let us consider a Riemann problem for the Burgers equation with  $\mathbf{u}_l = 1$  and  $\mathbf{u}_r = -1$ . Therefore, for any  $a \geq 1$ , solutions of the following form are weak solutions of the Riemann problem.

$$\mathbf{u}(x, t) = \begin{cases} 1 & \text{when } \frac{x}{t} < \frac{1-a}{2} \\ -a & \text{when } \frac{1-a}{2} < \frac{x}{t} < 0 \\ a & \text{when } 0 < \frac{x}{t} < \frac{a-1}{2} \\ -1 & \text{when } \frac{x}{t} > \frac{a-1}{2} \end{cases}$$

# Representation in the space $(x,t)$



# Representation in the space $(u, f(u))$



# What is Wrong! How to cure it!

Some of the characteristics have their initial values on the discontinuity where the function is not defined in a strong sense.

## Condition for uniqueness of the weak solution

Characteristics will not have their initial values localized on a discontinuity.

For scalar equation this means that, if  $s$  is the local velocity of a discontinuity between  $\mathbf{u}_l$  (on the left) and  $\mathbf{u}_r$  (on the right), This gives the **Entropy Condition**

$$f'(\mathbf{u}_l) \geq s \geq f'(\mathbf{u}_r)$$

A simplify version is the **Olenik Condition**

$$\frac{f(\mathbf{u}) - f(\mathbf{u}_l)}{\mathbf{u} - \mathbf{u}_l} \geq s \geq \frac{f(\mathbf{u}) - f(\mathbf{u}_r)}{\mathbf{u} - \mathbf{u}_r}, \quad \forall \mathbf{u} \in (\mathbf{u}_l, \mathbf{u}_r)$$

# Mathematical notion of entropy : for uniqueness

A convex function  $\eta(\mathbf{u})$  is an entropy of the conservation law if there exist a flux function  $\Psi(\mathbf{u})$  such as

$$\eta'(\mathbf{u})f'(\mathbf{u}) = \Psi'(\mathbf{u}) \quad pp. \quad (2)$$

Krushkov Family “entropy/flux” :

$$\eta(\mathbf{u}) = |\mathbf{u} - k|, \quad \Psi(\mathbf{u}) = \text{sgn}(\mathbf{u} - k) (f(\mathbf{u}) - f(k)) \quad \forall k \in \mathbb{R}$$

## Weak Entropy solution

A weak solution is called “Weak Entropy solution” if for any entropy/flux  $(\eta, \Psi)$  we have :

$$\partial_t \eta(\mathbf{u}) + \partial_x \Psi(\mathbf{u}) \leq 0 \quad (3)$$

jump conditions :  $\Psi(\mathbf{u}_r) - \Psi(\mathbf{u}_l) \leq s \left[ \eta(\mathbf{u}_r) - \eta(\mathbf{u}_l) \right]$

# Uniqueness of entropy solution.

## Theorem

*If the piecewise continue functions  $\mathbf{u}$  and  $v$  are weak entropy solutions of scalar conservation. Then*

$$\|\mathbf{u}(\cdot, t) - v(\cdot, t)\|_{L^1(\mathbb{R})} \leq \|\mathbf{u}_0(\cdot) - v_0(\cdot)\|_{L^1(\mathbb{R})} \quad \forall t \geq 0$$

As consequence, there is a unique piecewise continue weak entropy solutions of a scalar conservation law with the initial condition  $\mathbf{u}_0(\cdot)$ .

# Uniqueness of entropy solution.

$$\text{Total variation : } TV(\mathbf{u}(\cdot, t)) = \limsup_{\varepsilon \rightarrow 0} \left( \frac{1}{\varepsilon} \int_{\mathbb{R}} |\mathbf{u}(x + \varepsilon, t) - \mathbf{u}(x, t)| dx \right)$$

## Theorem

If  $\mathbf{u}_0 \in L^1(\mathbb{R}) \cap L^\infty(\mathbb{R}) \cap BV(\mathbb{R})$  and  $f \in C^1(\mathbb{R})$ . Then there is a unique entropy solution  $\mathbf{u}(x, t) \in L^\infty(\mathbb{R} \times \mathbb{R}^{+,*}) \cap C^1(0, T; L^1(\mathbb{R}))$  of the Cauchy problem, and

$$\|\mathbf{u}(\cdot, t)\|_{L^\infty(\mathbb{R})} \leq \|\mathbf{u}_0(\cdot)\|_{L^\infty(\mathbb{R})} \quad (4)$$

Moreover,  $\mathbf{u}(\cdot, t) \in BV(\mathbb{R}) \quad \forall t > 0$  and

$$TV(\mathbf{u}(\cdot, t)) \leq TV(\mathbf{u}_0(\cdot)) \quad (5)$$



# Riemann Problem

The Riemann problem consists of the “Hyperbolic” equation with the initial data defined as

$$\mathbf{u}_0(x) = \begin{cases} \mathbf{u}_l & \text{when } x < 0 \\ \mathbf{u}_r & \text{when } x > 0 \end{cases}$$

where  $\mathbf{u}_l$  and  $\mathbf{u}_r$  constants and the solution non trivial :  $\mathbf{u}_l \neq \mathbf{u}_r$ .

## The Riemann problem is fundamental

To understanding the mathematical theory of hyperbolic problems  
To derive Godunov type finite volume schemes.

If  $\mathbf{u}(x, t)$  is a solution of the Riemann problem, then  $\forall \alpha > 0$ ,  $\mathbf{u}(x\alpha, t\alpha)$  is also a solution of the Riemann problem. Therefore  $\mathbf{u}(x, t) = \mathbf{u}(\xi)$  with  $\xi = \frac{x}{t}$ .

# Riemann Problem : Rarefaction wave

The aim here is to define a smooth solution of the Riemann problem. For any  $t > 0$  this solution must satisfy the equation

$$\partial_t \mathbf{u} + f'(\mathbf{u}) \partial_x \mathbf{u} = 0 \iff \frac{1}{t} \left( -\xi + f'(\mathbf{u}) \right) \partial_\xi \mathbf{u} = 0$$

Therefore

- either  $\partial_\xi \mathbf{u} = 0$  and  $\mathbf{u}(\xi)$  is constant.
- or  $\partial_\xi \mathbf{u} \neq 0$  and

$$\mathcal{F}(\xi, \mathbf{u}) = f'(\mathbf{u}(\xi)) - \xi = 0$$

Moreover, if  $f \in C^2(\mathbb{R})$

$$f''(\mathbf{u}(\xi)) \partial_\xi \mathbf{u} = 1 \implies f \text{ is nonlinear}$$

# Riemann Problem : Rarefaction wave

$$\partial_{\xi} \mathbf{u} = 0 \quad \text{or} \quad f'(\mathbf{u}(\xi)) = \xi$$

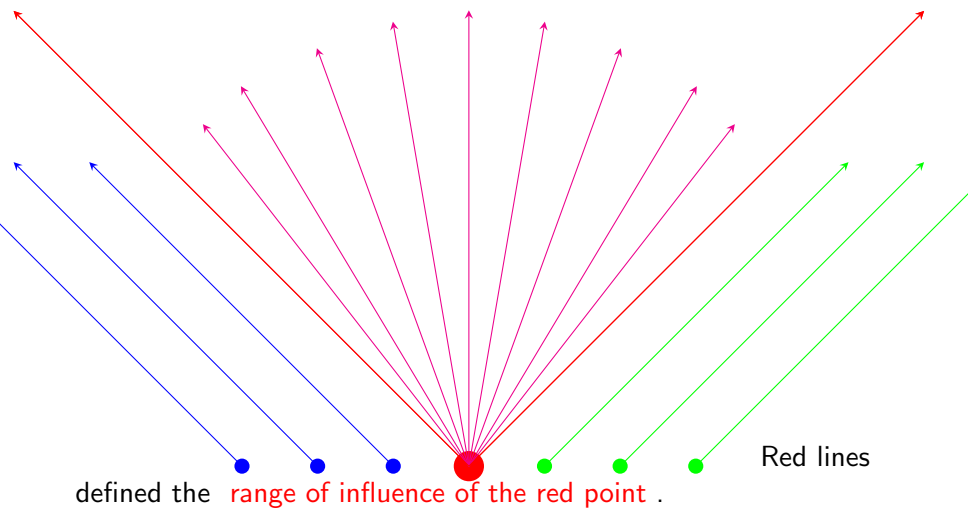
## Caractérisation of a single Rarefaction wave

$$\forall \xi \in \mathbb{R}, \forall \zeta \in \mathbb{R} \quad : \quad \xi < \zeta \implies f'(\mathbf{u}(\xi)) \leq f'(\mathbf{u}(\zeta))$$

The necessary condition to have a single rarefaction wave as solution of the Riemann problem is that

- $f$  is nonlinear
- $f'(\mathbf{u}_l) \leq f'(\mathbf{u}_r)$
- $f'(\mathbf{u})$  is a monoton function, at least in the interval  $(\mathbf{u}_l, \mathbf{u}_r)$

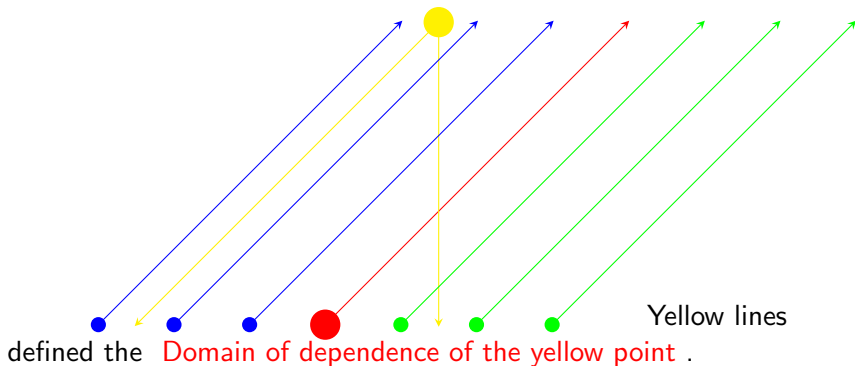
# Riemann Problem : Single Rarefaction wave



# Riemann Problem : Single Contact Discontinuity ( $m = 0$ )

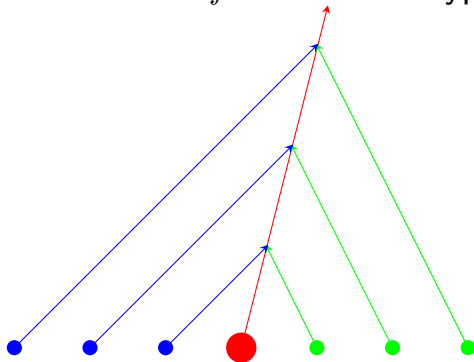
$$m = f(\mathbf{u}_r) - s\mathbf{u}_r = f(\mathbf{u}_l) - s\mathbf{u}_l = 0 \implies f(\mathbf{u}) \equiv s\mathbf{u}$$

$f$  is linear type



# Riemann Problem : Shock wave ( $m \neq 0$ )

$$m = f(\mathbf{u}_r) - s\mathbf{u}_r = f(\mathbf{u}_l) - s\mathbf{u}_l \neq 0$$
$$\Rightarrow f \text{ is nonlinear type}$$



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