



# The laser which came from the cold

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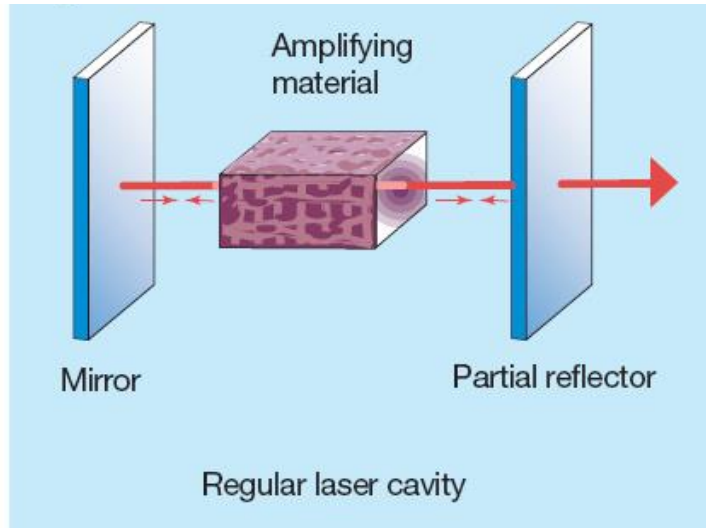
...has been done: @ INLN (post-doc, 2007 – 2009)  
@ Tübingen University, Germany (post-doc, 2010 – 2012)  
@ INLN (CR CNRS, since end 2012)

Most of it is contained in the following PhD thesis:

- Frank Michaud, Nice, 2008
- Nicolas Mercadier, Nice, 2010
- Alexander Schilke, Tübingen, 2013
- Quentin Baudouin, Nice, 2013

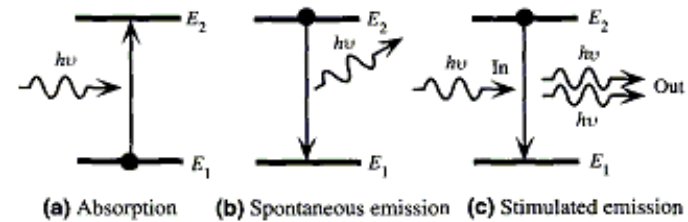
The work at INLN has been supervised by **Robin Kaiser**

More information at: <http://www.inln.cnrs.fr/activites/themesrecherche/atomes-froids>



Two ingredients for a **standard laser** :

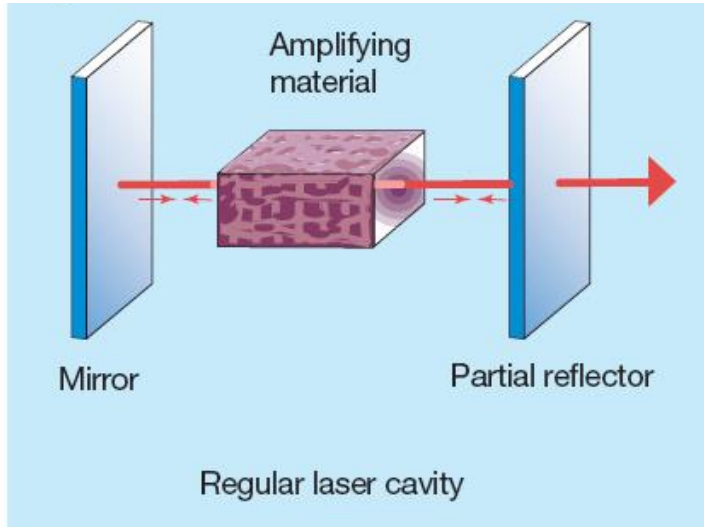
- 1) An amplifying material  
(Gain based on stimulated emission)



- 2) An optical cavity

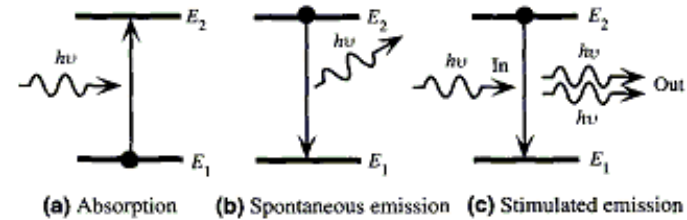
Roles of the optical cavity:

- To provide feedback
  - Chain reaction: intensity grows until gain saturation
- Fabry-Perot interferometer
  - Mode selection: spatial and temporal coherence properties



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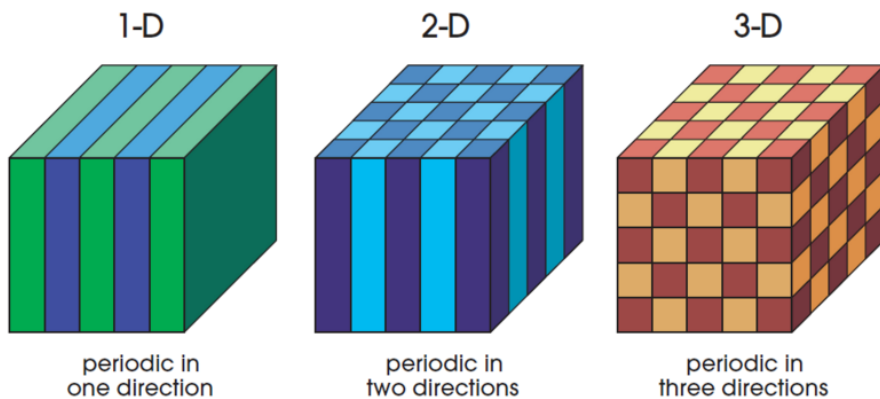
→ Chain reaction: intensity grows until gain saturation

~~- Fabry-Perot interferometer~~

~~→ Mode selection: spatial and temporal coherence properties~~

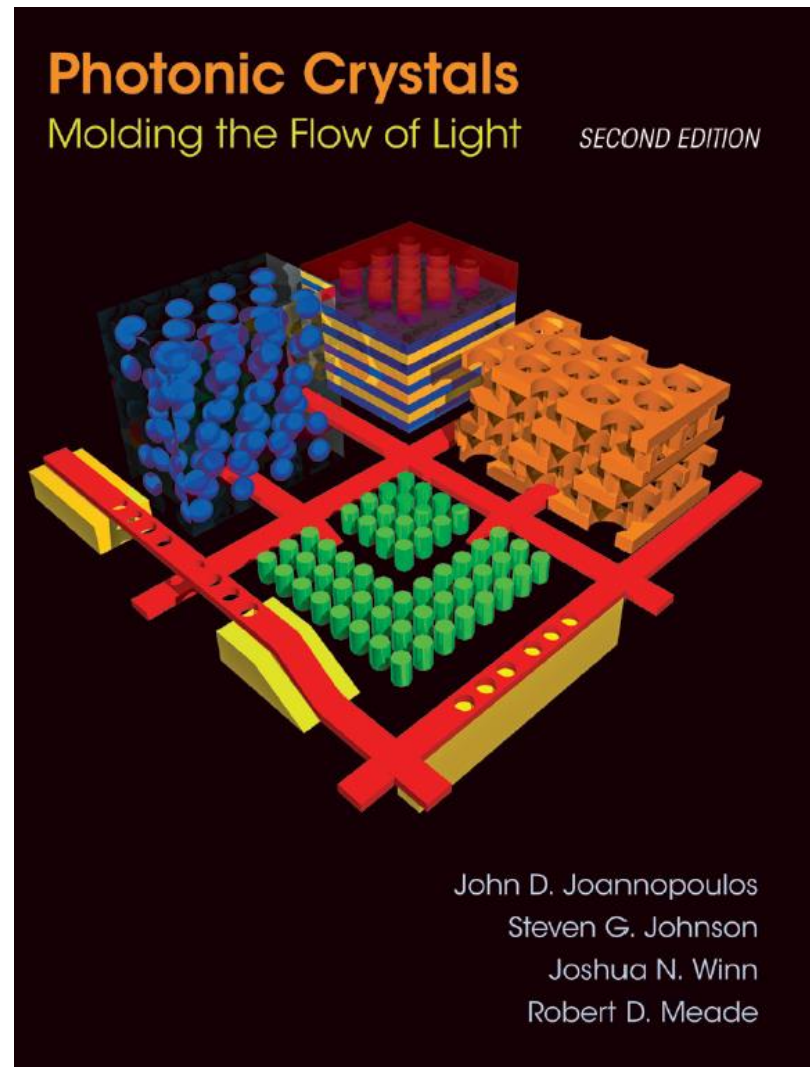


## First possibility: use a **periodic medium**

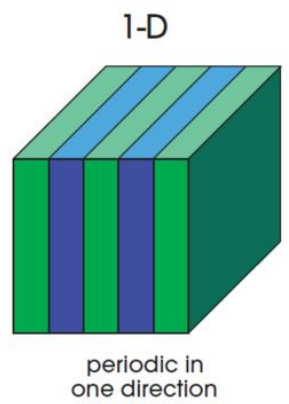


Photonic crystals can confine light in 1D, 2D or 3D.  
Can be combined with light emitters (e.g. quantum dots) or amplifiers.

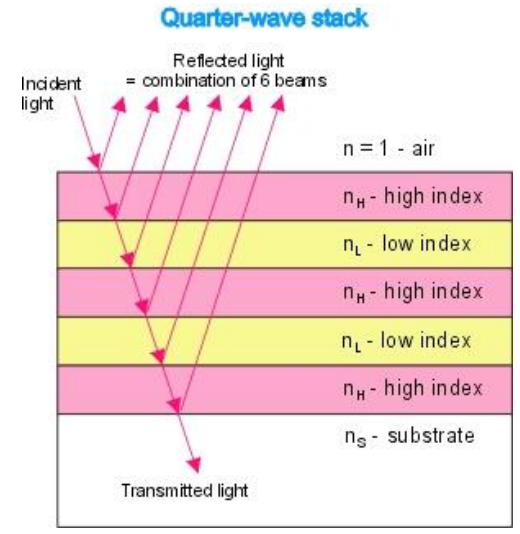
→ “photonic crystal lasers” / “nanolasers”



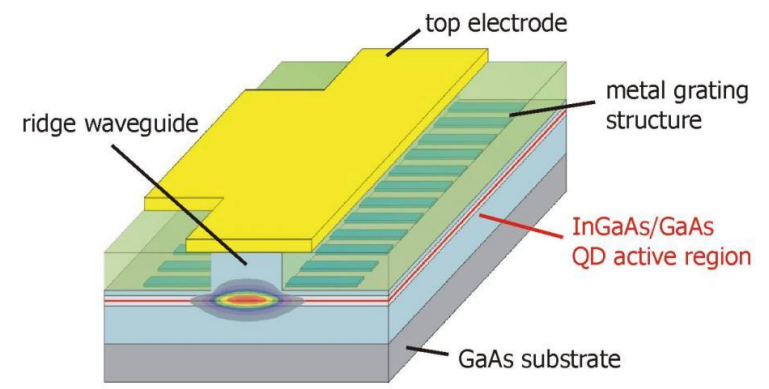
## First possibility: use a **periodic medium** (1D case)



Light propagation in a 1D periodic medium is known since Rayleigh.  
→ Bragg mirrors



Active medium (gain) + 1D modulation: known since the 70s...  
→ “distributed feedback laser” (DFB).



Kogelnik & Shank, Appl. Phys. Lett. **18**, 152 (1971).

## Second possibility: use a **diffusive (disordered) medium**

Many scatterers at random positions

→ Multiple scattering

→ “Radiation trapping”



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Many scatterers at random positions

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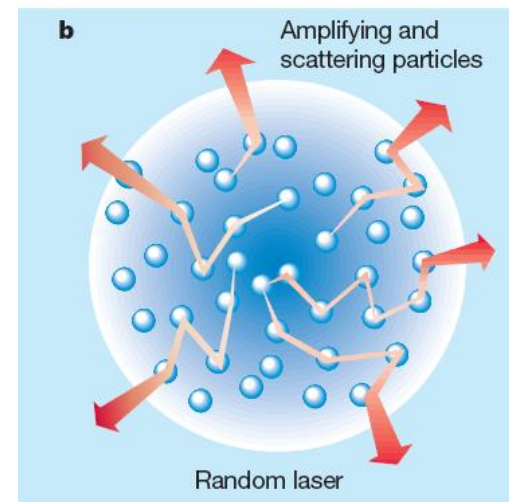
Multiple scattering + gain: “**Random laser**”

→ Emission in all directions

→ Mode and coherence properties: complicated !

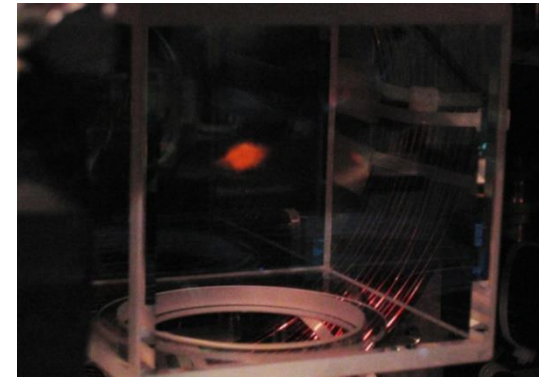
Initial proposal in 1968 !

First realized in 1995, extensively studied since the 2000s



Letokhov, Sov. Phys. JETP **26**, 835 (1968).  
Wiersma, Nature Phys. **4**, 359 (2008).





We use atomic vapors, laser-cooled to  $T \sim 20\text{-}150 \mu\text{K}$ .

→ Almost no Doppler broadening

→ **Very sharp resonance** (width  $6 \text{ MHz} \leftrightarrow 0.000012 \text{ nm} \leftrightarrow 25 \text{ neV} \leftrightarrow 0.0002 \text{ cm}^{-1}$ )

For near-resonant light, a cold-atom vapor is an optical medium with some properties many orders of magnitude different than usual (standard dielectric media):

- Highly diffusive: very opaque without absorption
- Highly dispersive
- Highly nonlinear (a few mW)
- Very sensitive to external fields → highly versatile

$$\sigma = \frac{3\lambda^2}{2\pi} \frac{1}{1 + 4\delta^2/\Gamma^2}$$

- Introduction
  
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- DFB lasing with cold atoms
  
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Rubidium 85

$\lambda = 780 \text{ nm}$

$\Gamma/2\pi = 6 \text{ MHz}$

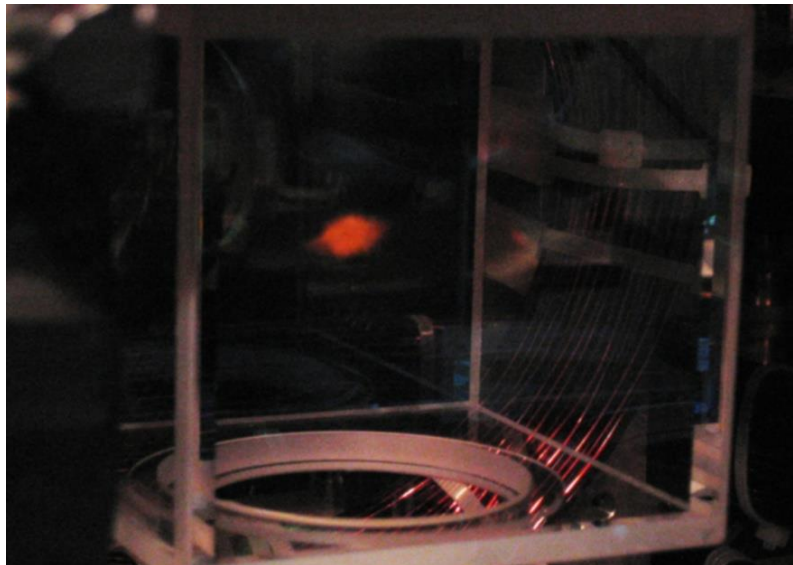
MOT parameters:

$N \sim 10^8\text{-}10^{10}$  atoms

$T \sim 20\text{-}150 \text{ }\mu\text{K}$

$L \sim 1\text{-}2 \text{ mm}$

$\rho \sim 10^{11} \text{ at/cm}^3$

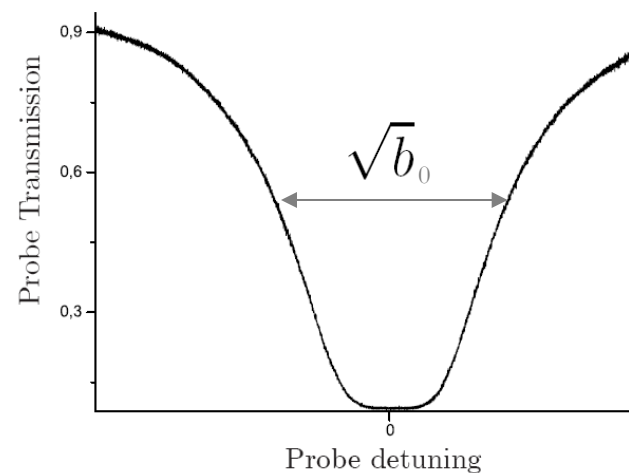
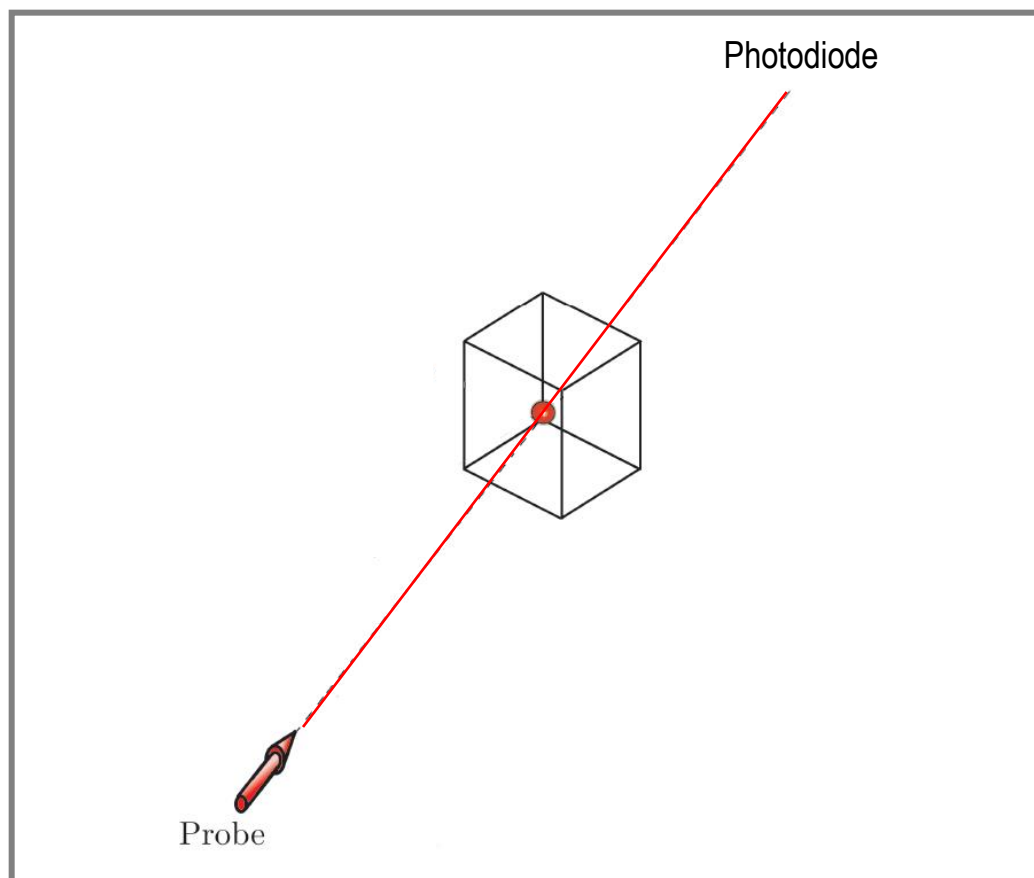


$$T = e^{-b} = e^{-\sigma\rho L} = e^{-L/\ell_{sc}}$$

Typically, on resonance,  $b_0 = 20 - 100$

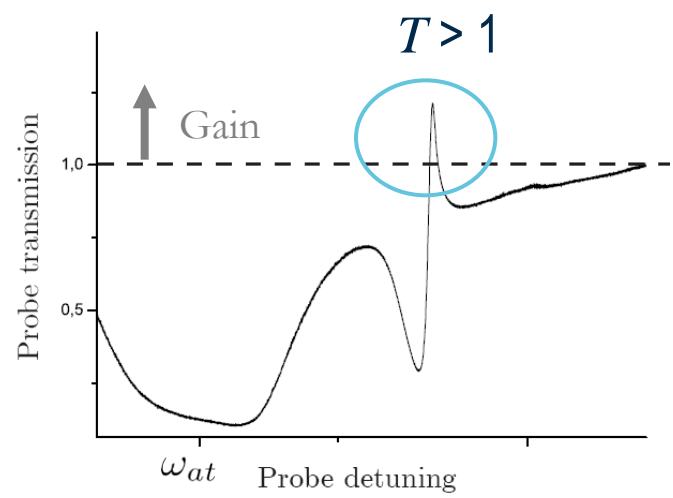
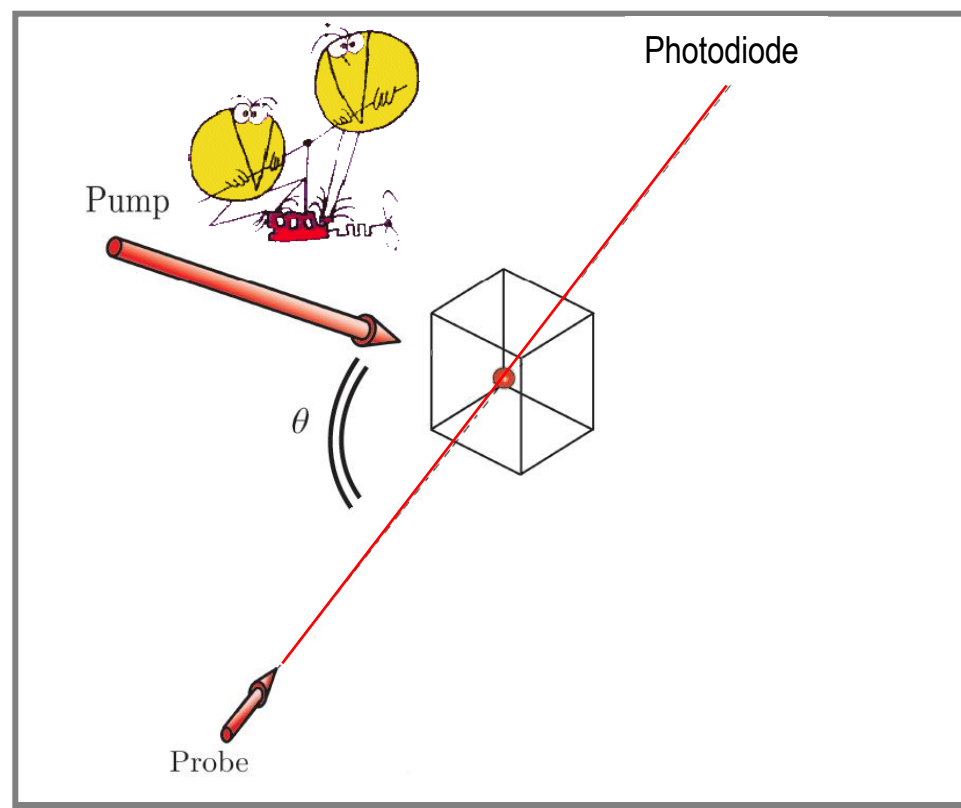
With some efforts: up to  $b_0 \sim 250$

## Spectroscopy in transmission



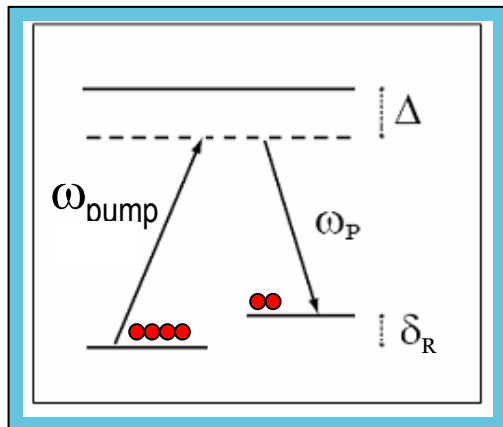
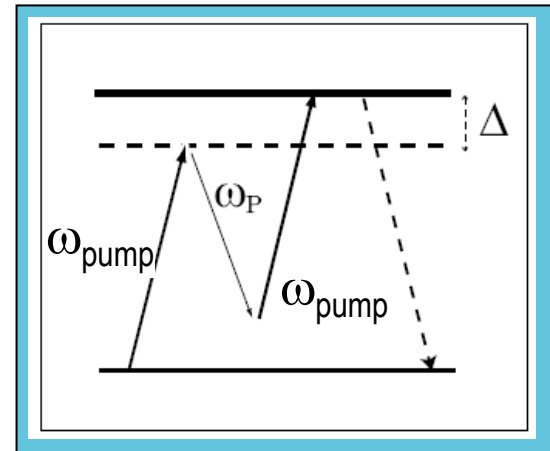
$b_0$  : on-resonance optical thickness

## Pump-probe spectroscopy



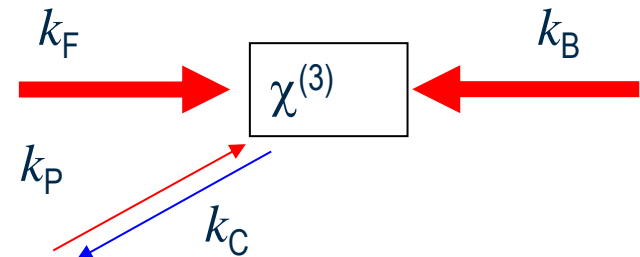
- **Mollow gain.** Two-level atoms + one pump:  
3-photon transition (population inversion in the dressed-state basis)

Mollow, Phys. Rev. A **5**, 2217 (1972).



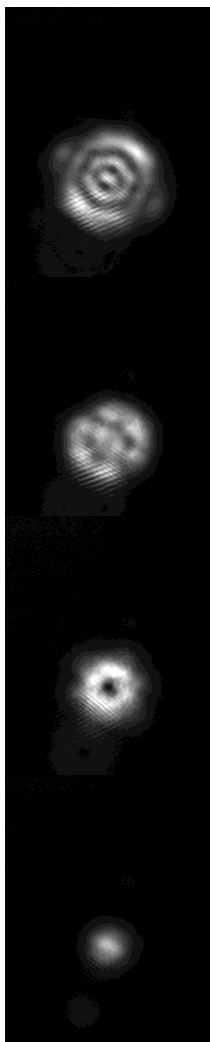
- **Raman gain.** Three-level atoms + one pump:  
2-photon Raman transition (population inversion between the two ground states – hyperfine or Zeeman levels)

- **Degenerate four-wave mixing.** Parametric gain using the nonlinear atomic susceptibility (needs two pumps)





**Laser radiation  $\approx 300 \mu\text{W}$**



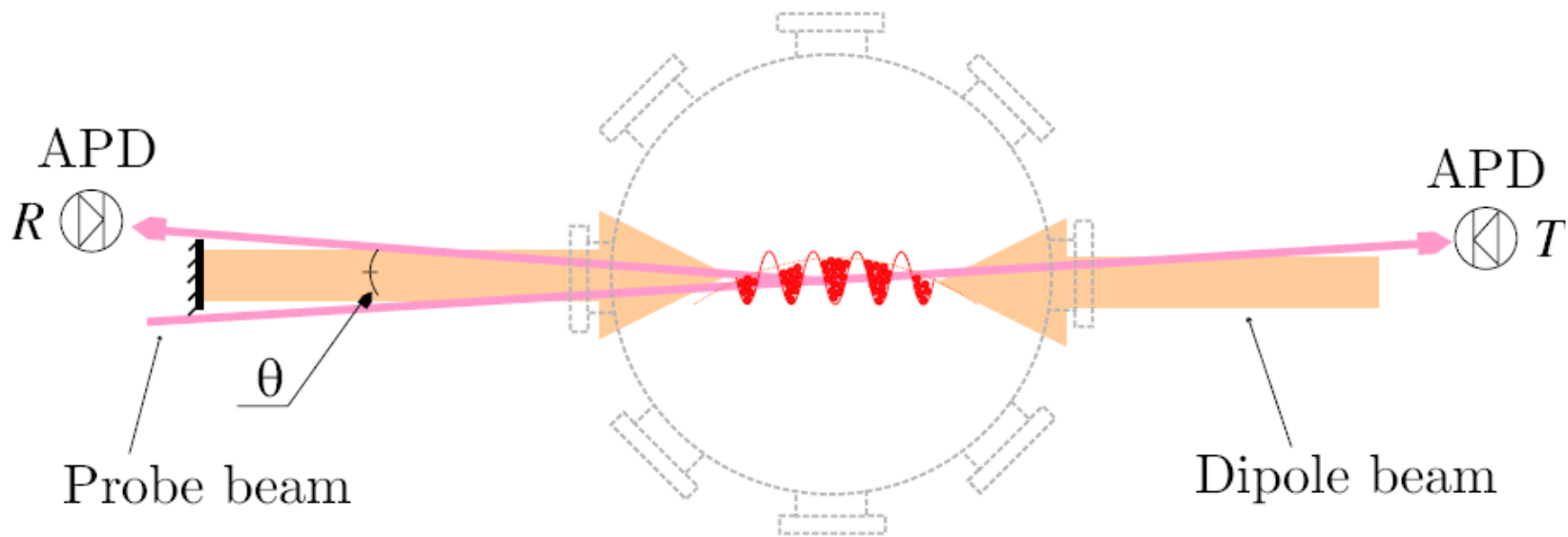
**Cold atoms inside !**

- Mollow laser for small pump detuning.
- (Zeeman) Raman laser for larger pump detuning, single pump.
- DFWM laser for larger pump detuning and two pumps.

W. Guerin, F. Michaud, R. Kaiser, Phys. Rev. Lett. **101**, 093002 (2008).

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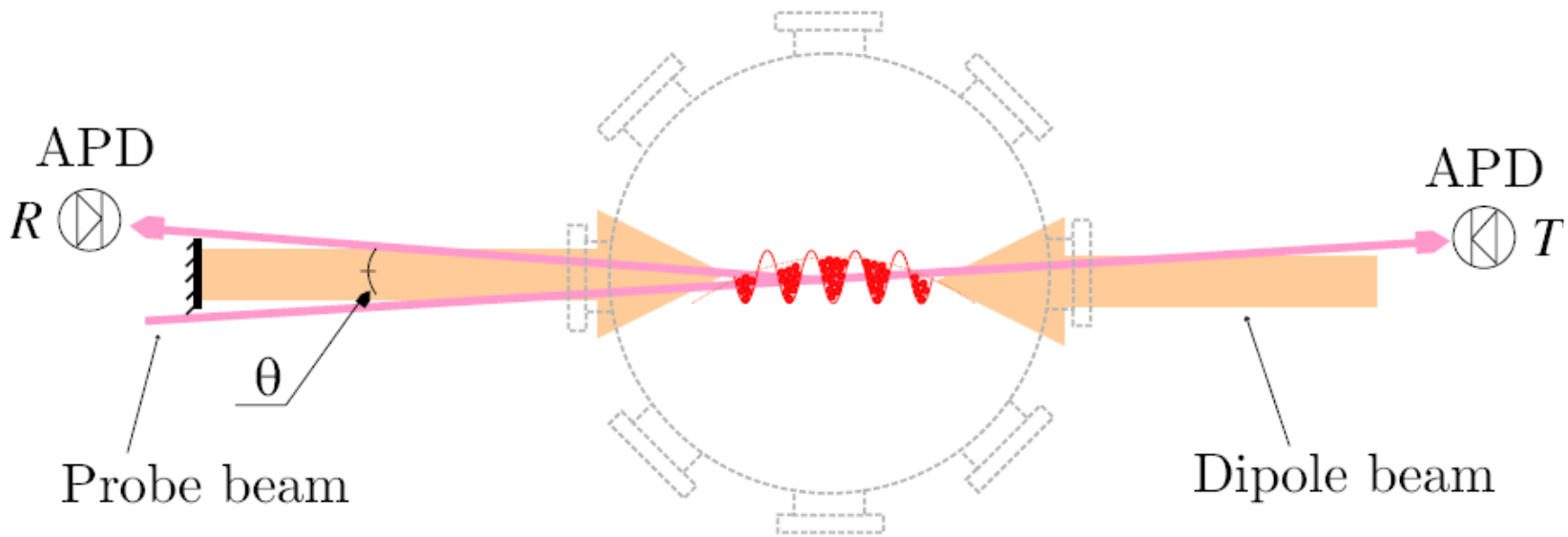


**Atoms:** laser-cooled  $^{87}\text{Rb}$ ,  $\lambda_0 = 780.24 \text{ nm}$ .

**Lattice beam:** tunable Ti-Sa laser, 1W, waist  $200 \mu\text{m}$ , wavelength  $\lambda_{\text{lat}} > \lambda_0$ .

**Detection tools:** probe beam and avalanche photodiodes (APD).

**Measurements:** transmission  $T$  and reflection  $R$  spectra.



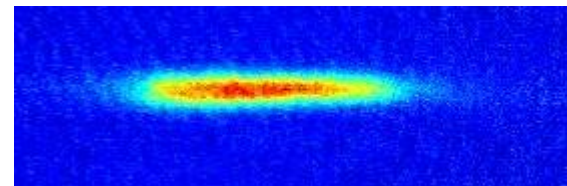
## Atomic sample:

$$N = 5 \times 10^7$$

$$T \sim 100 \mu\text{K}$$

$$\rho \sim 10^{11} - 10^{12} \text{ cm}^{-3}$$

$$\rightarrow n - 1 \sim 10^{-4} - 10^{-3}$$

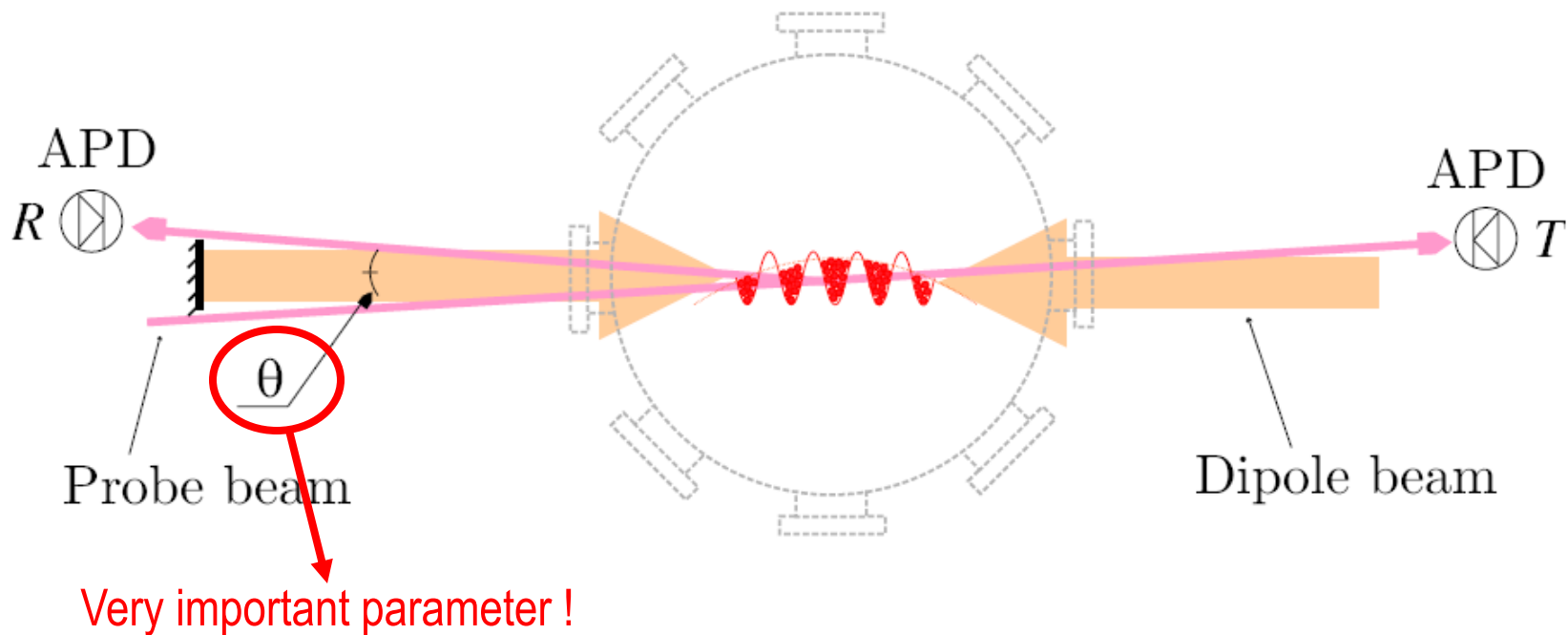


↑ ~ 200 μm



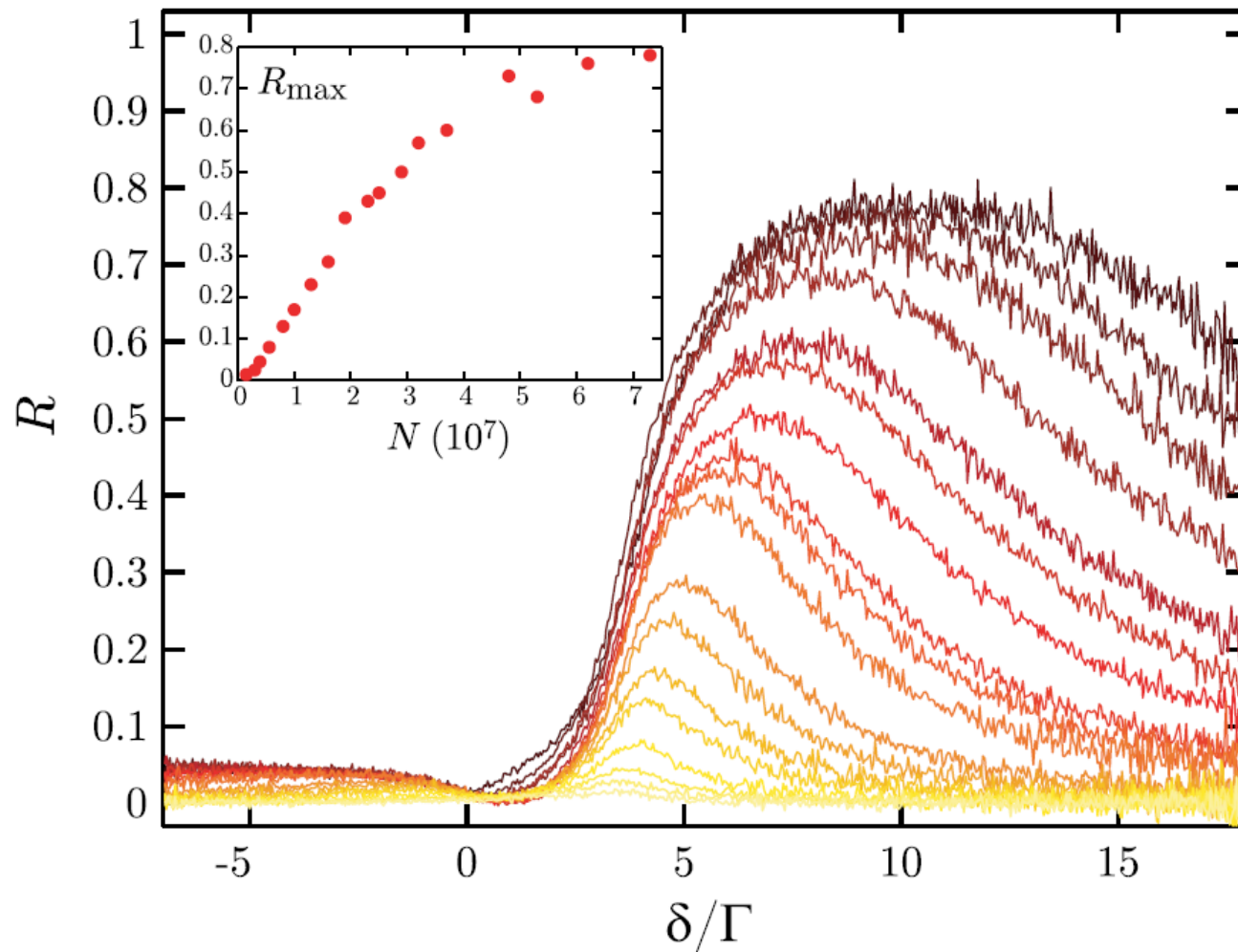
$L \sim 3 \text{ mm}$

→ 7700 atomic layers



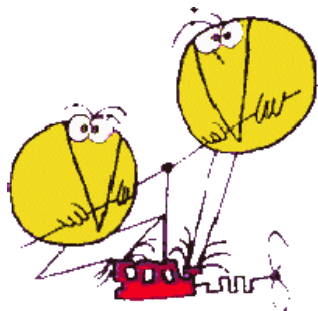
- $\lambda_{\text{lat}} > \lambda_0$  to trap the atoms, and the lattice period is  $\lambda_{\text{lat}}/2$
  - the refractive index  $n$  is nonnegligible only around  $\lambda \sim \lambda_0$
  - The Bragg condition can only be fulfilled with an angle such that  $\lambda_{\text{lat}} \sim \lambda_0/\cos(\theta)$
- If  $\theta$  too large : bad overlap between the probe beam and the atomic cloud.

Bragg reflection spectra for increasing atom number (or density  $\rho$ ), at the optimum  $\lambda_{\text{lat}}$ .

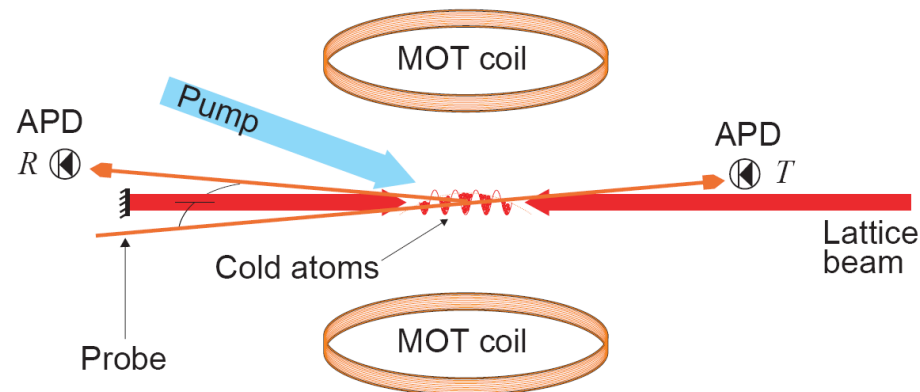


→ 80% reflection

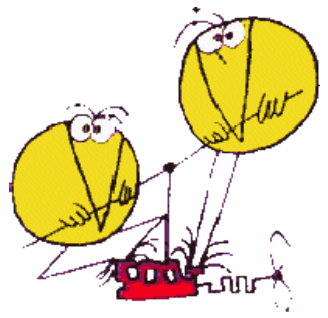
Schilke et al., Phys. Rev. Lett. **106**, 223903 (2011).



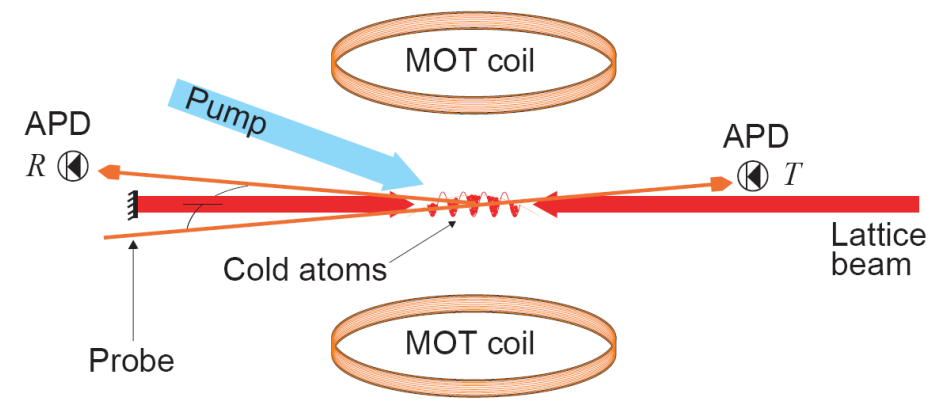
We have to **pump** !



Several gain mechanisms are possible with cold atoms (see previous part!).



We have to **pump** !

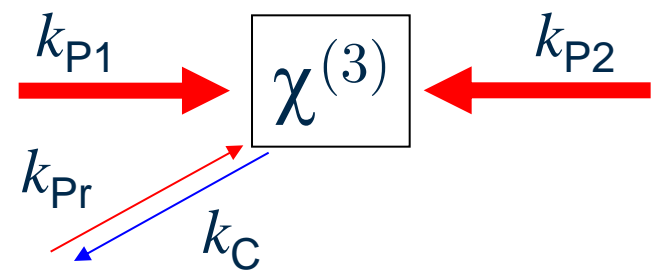


Several gain mechanisms are possible with cold atoms (see previous part!).

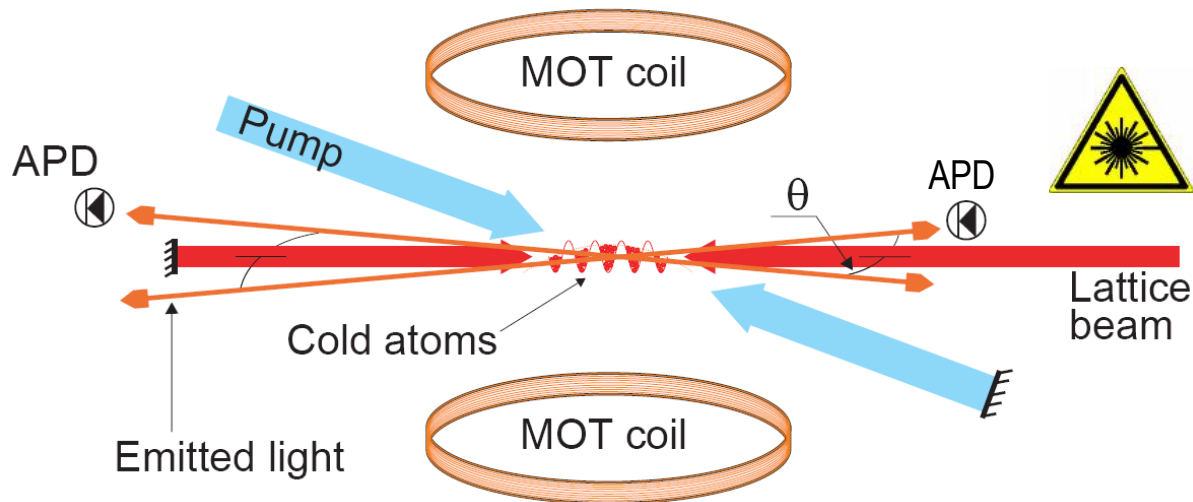
One possibility: **four-wave mixing**

**Phase-conjugation mechanism**

→ “backward gain”



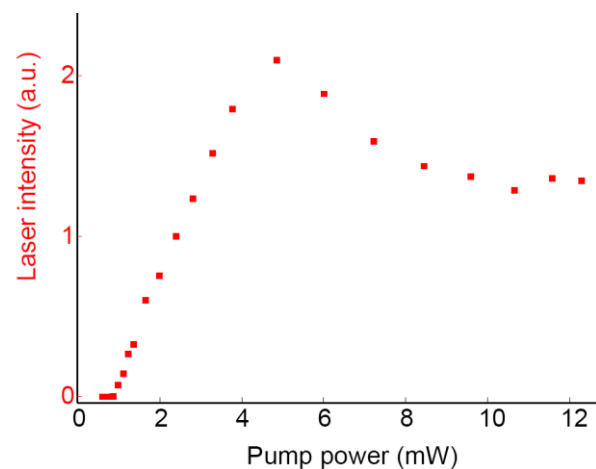
**Degenerate FWM:**  
 $\omega_{P1} = \omega_{P2} = \omega_{Pr} = \omega_C$

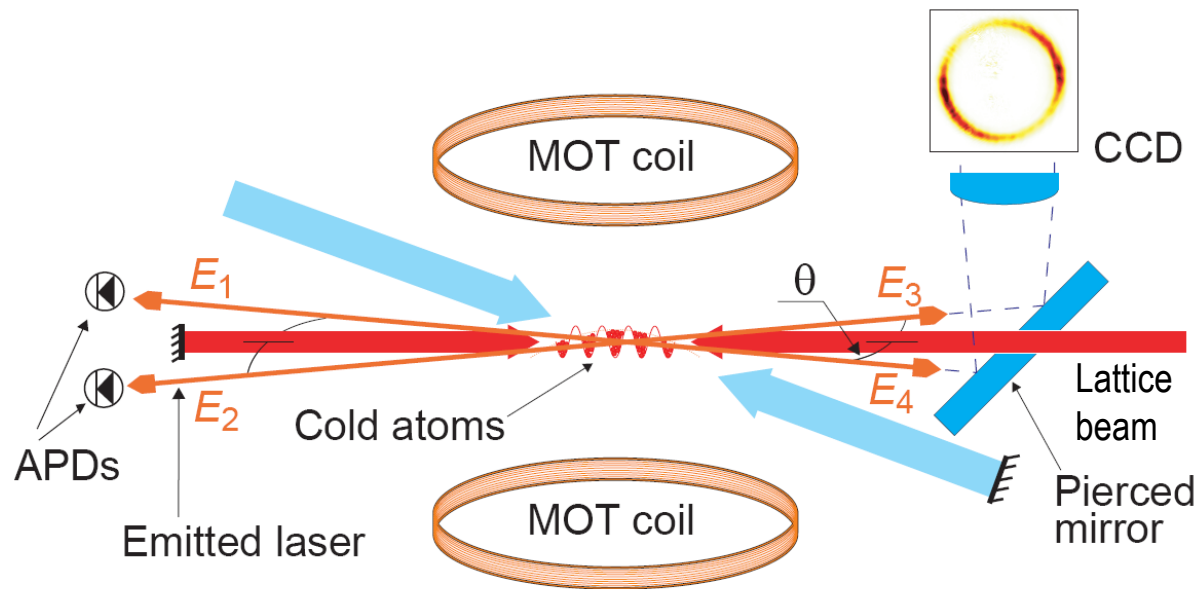


→ Huge signals on our  $R$  and  $T$  photodiodes **even without probe beam !**

→ Threshold with the pump power

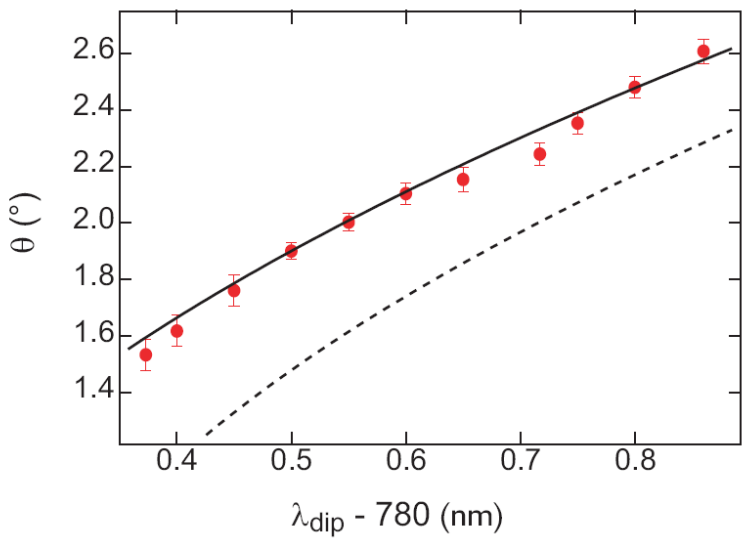
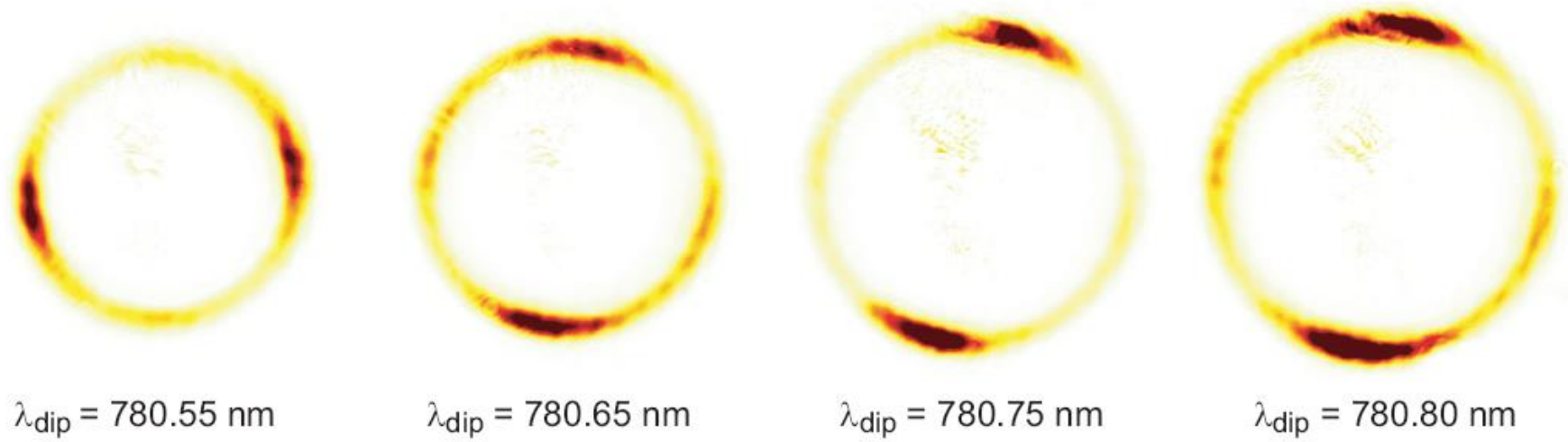
→ **Laser**





→ Cone-shaped emission





Well explained by the **Bragg condition**:

$$\lambda_{\text{dip}} \cos(\theta) = \lambda_0 / \bar{n}$$

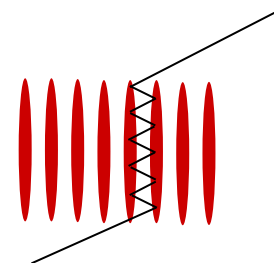
$$\bar{n} - 1 = (2.2 \pm 0.5) \times 10^{-4}$$

Schilke *et al.*, Nature Photon. **6**, 101 (2012).



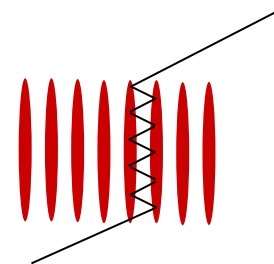
$\theta \neq 0 \rightarrow$  the Bragg feedback alone is **unstable** (walk-off)

Why is it working ?

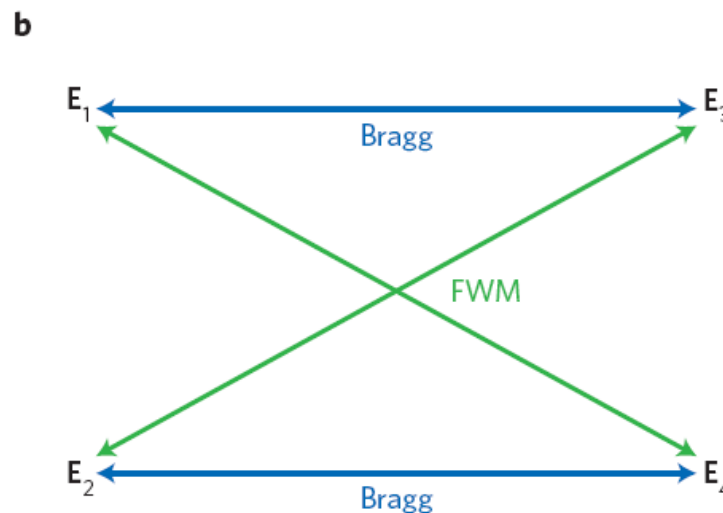
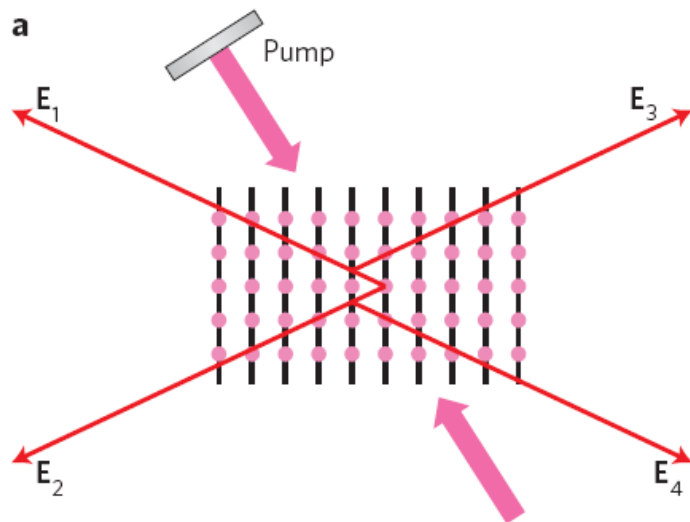




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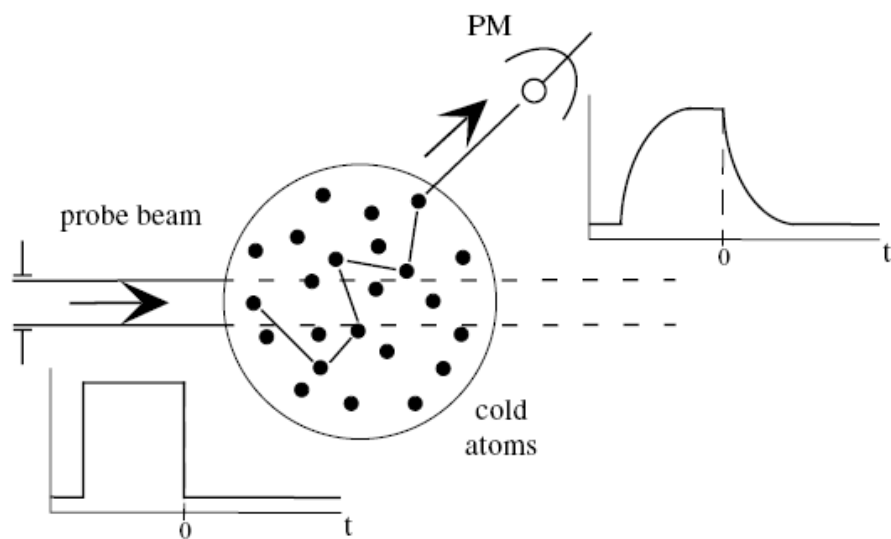


FWM is a **phase-conjugation** process (backward gain)

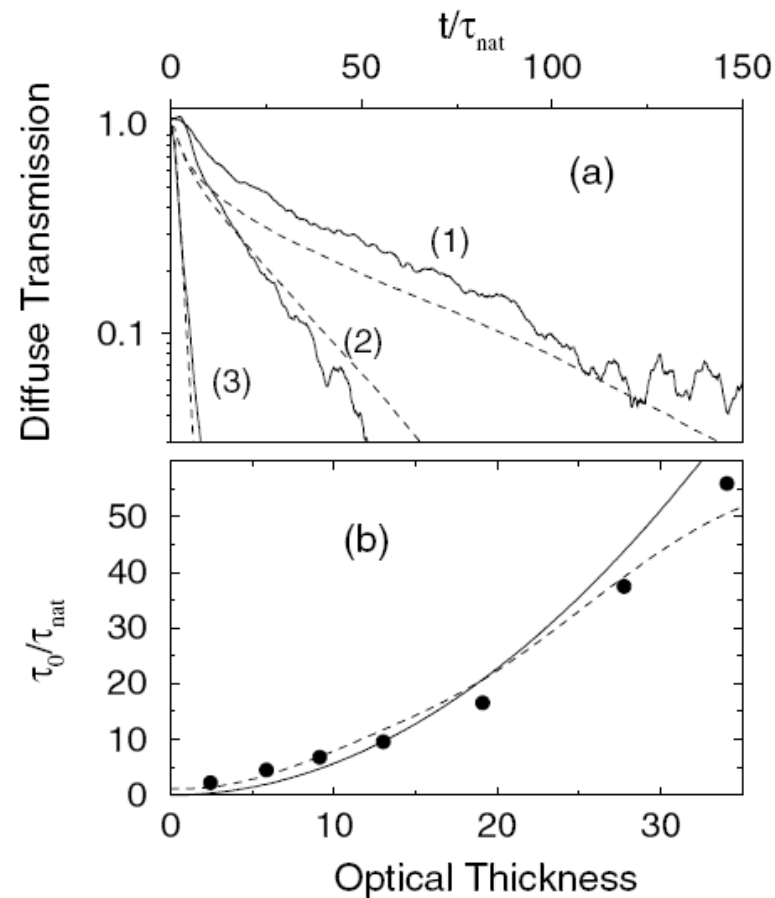


$\rightarrow$  creates a feedback loop without walk-off  
(No observed DFB laser with Raman gain !)

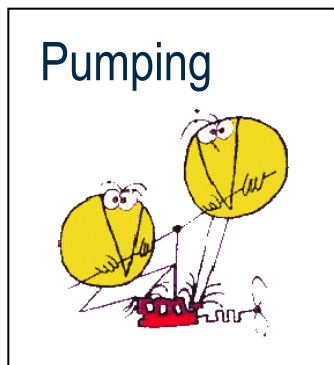
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$$v_E \approx 3.1 \times 10^{-5} c$$



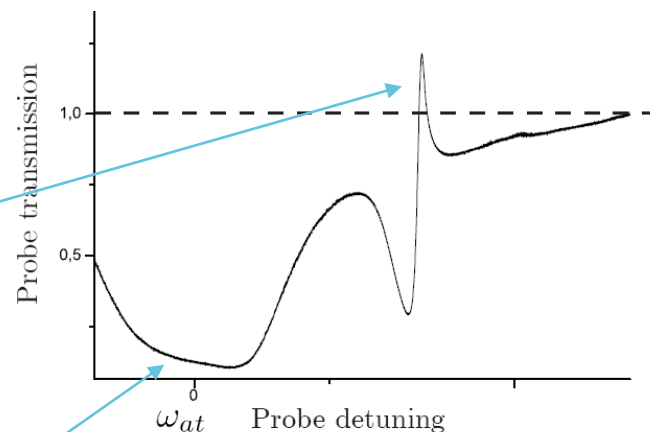
The scatterers and the amplifiers are **the same atoms** !



Gain 😊

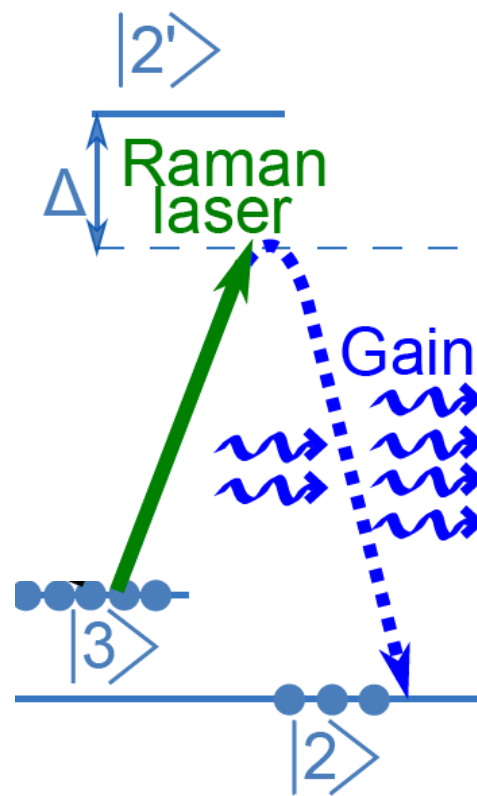
Saturation ☹️

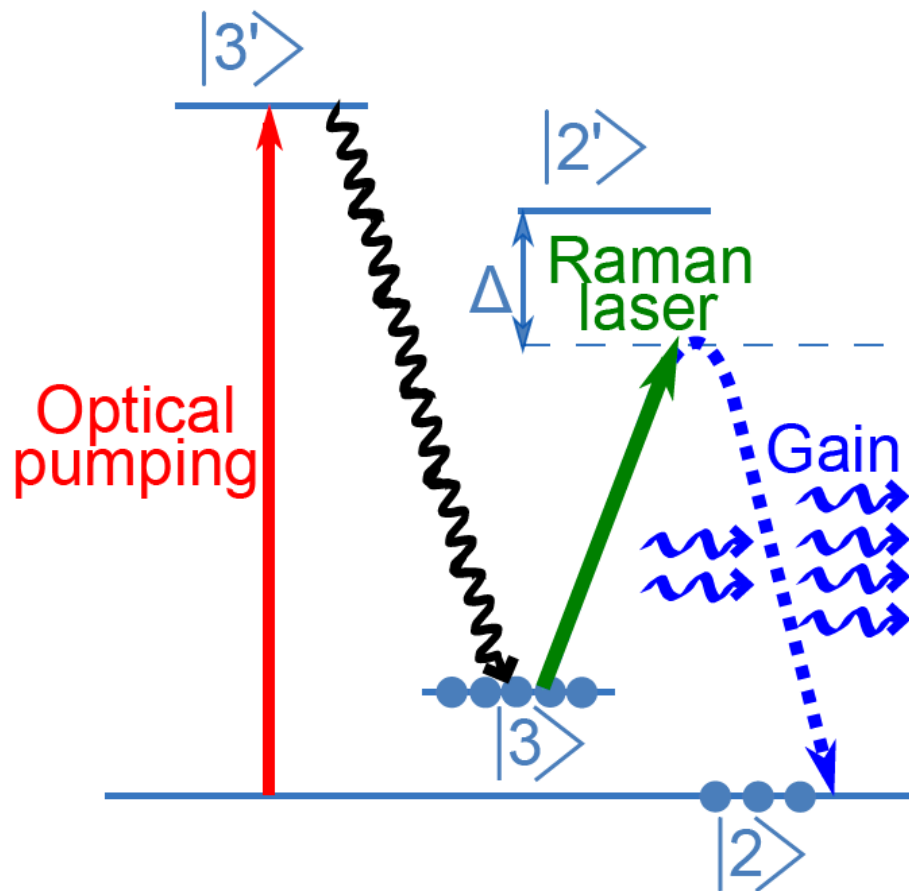
- elastic scattering ↘
- inelastic scattering ↗



Gain and scattering do not occur at the same frequency !!! ☹️ ☹️ ☹️

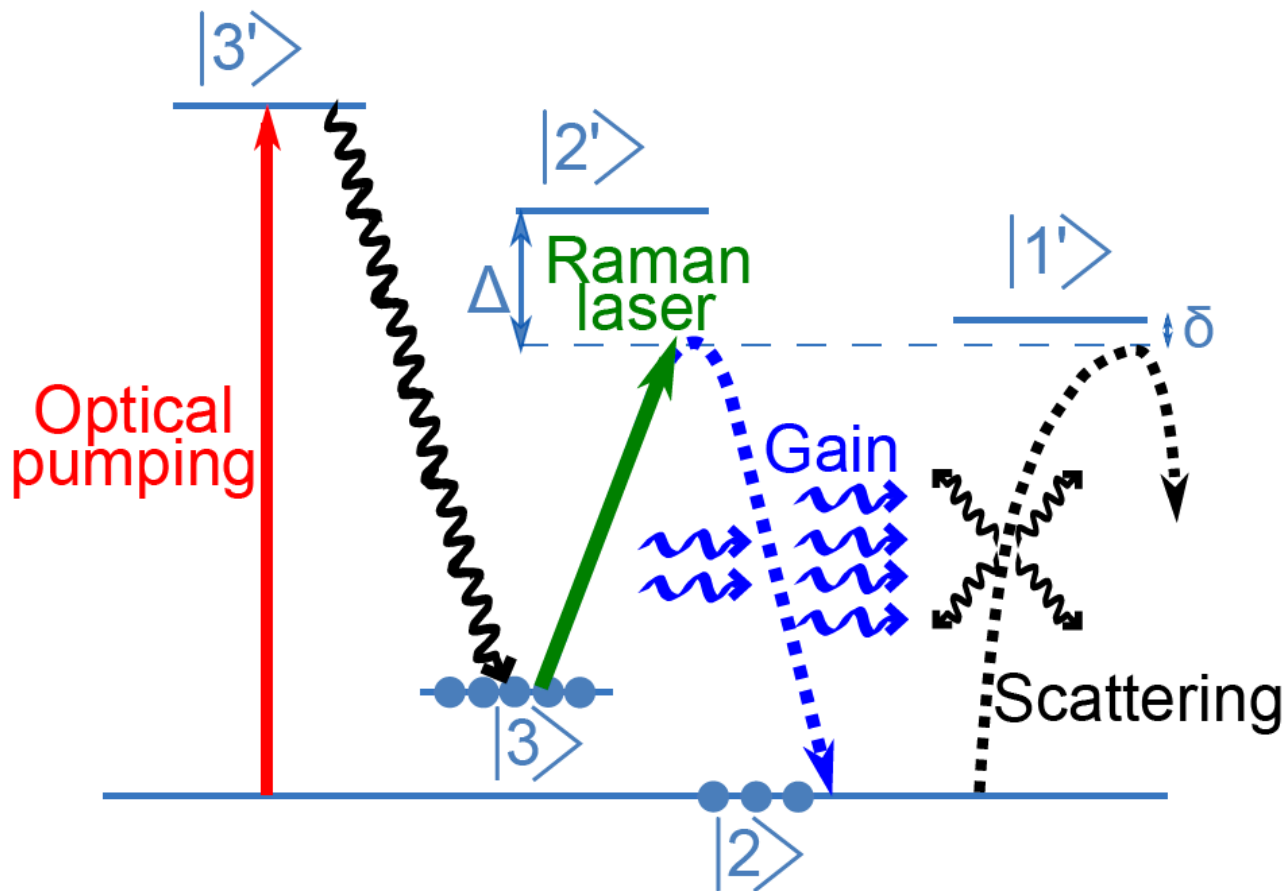
Is it possible to get enough scattering and gain *simultaneously* ?



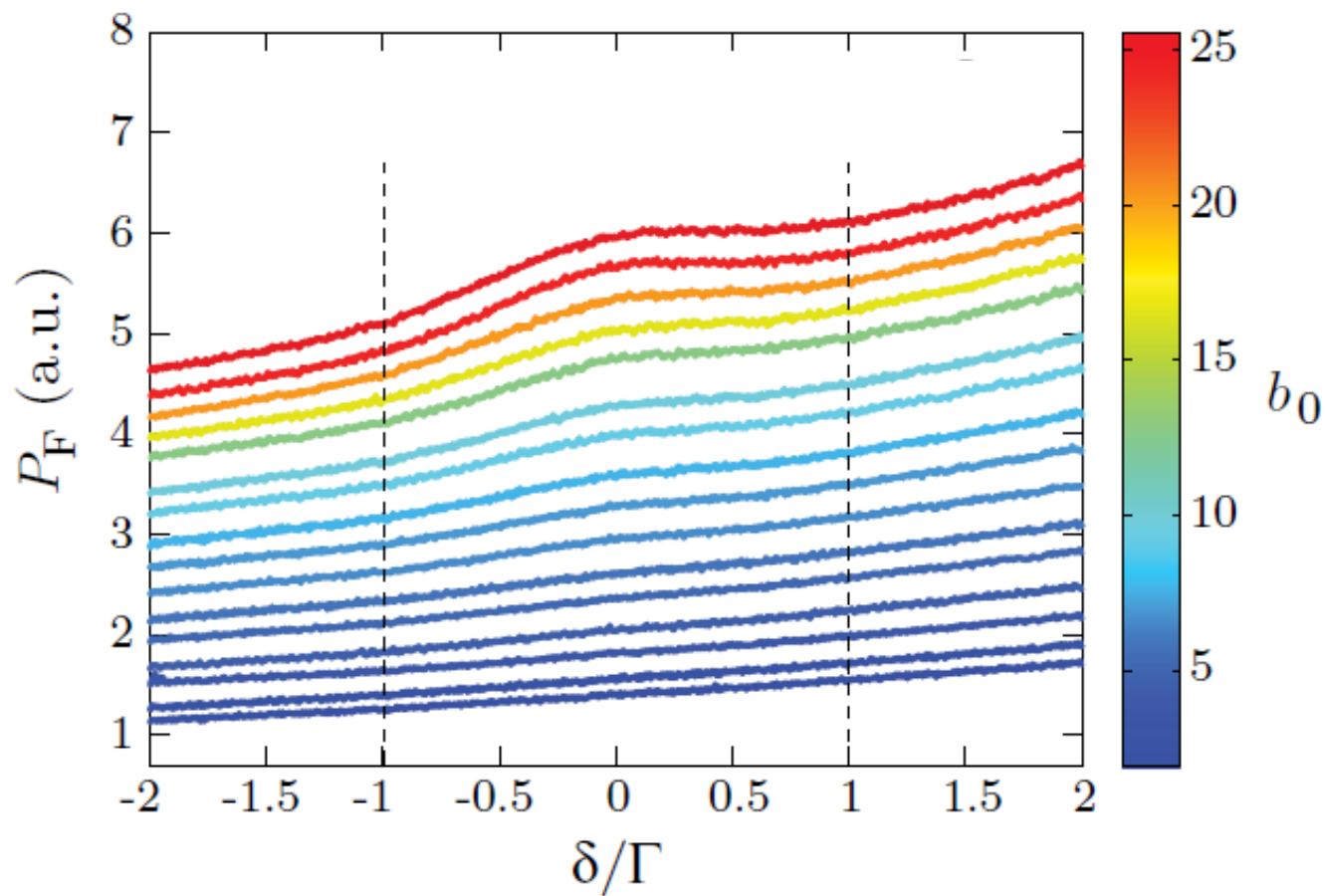




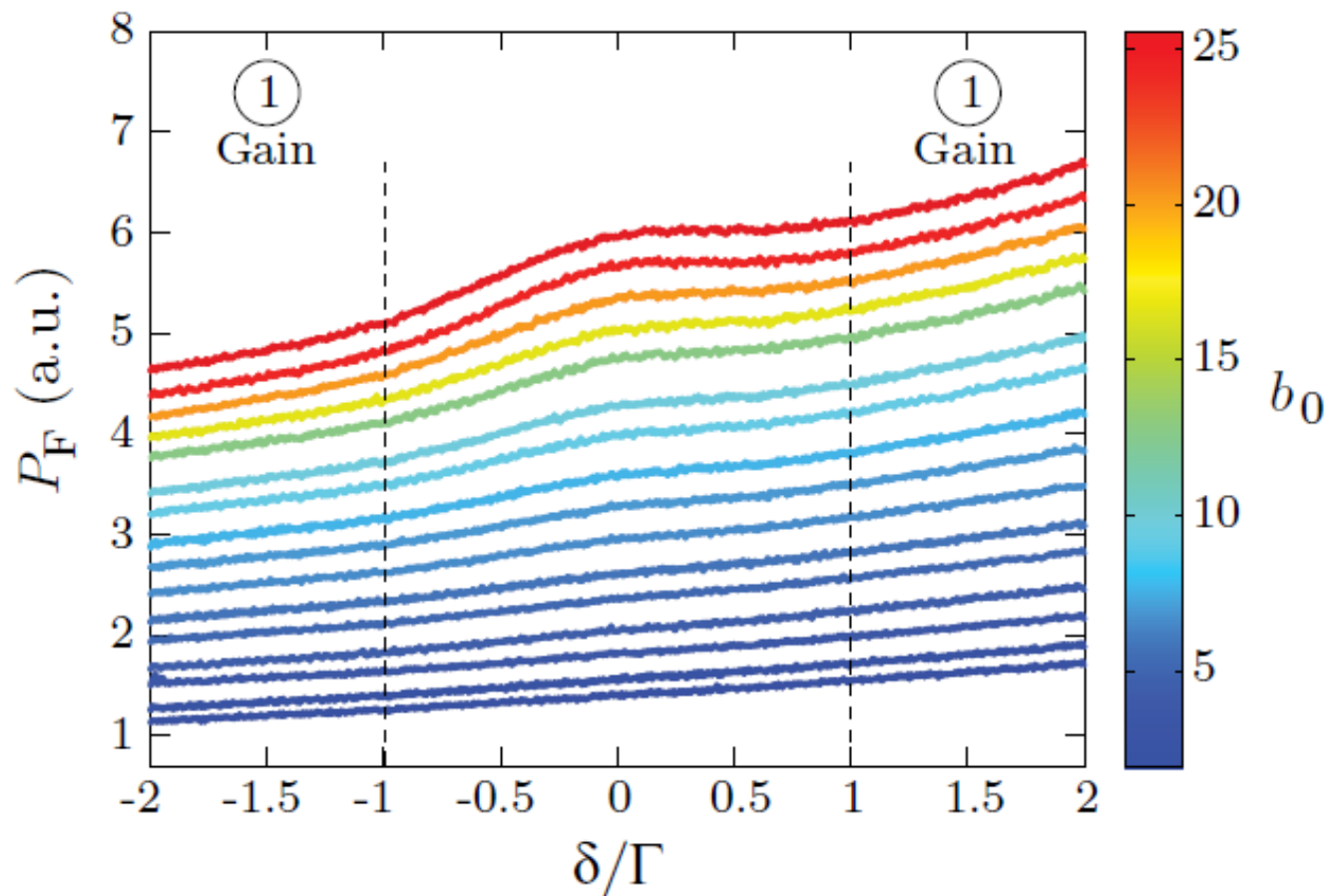
## with additional scattering



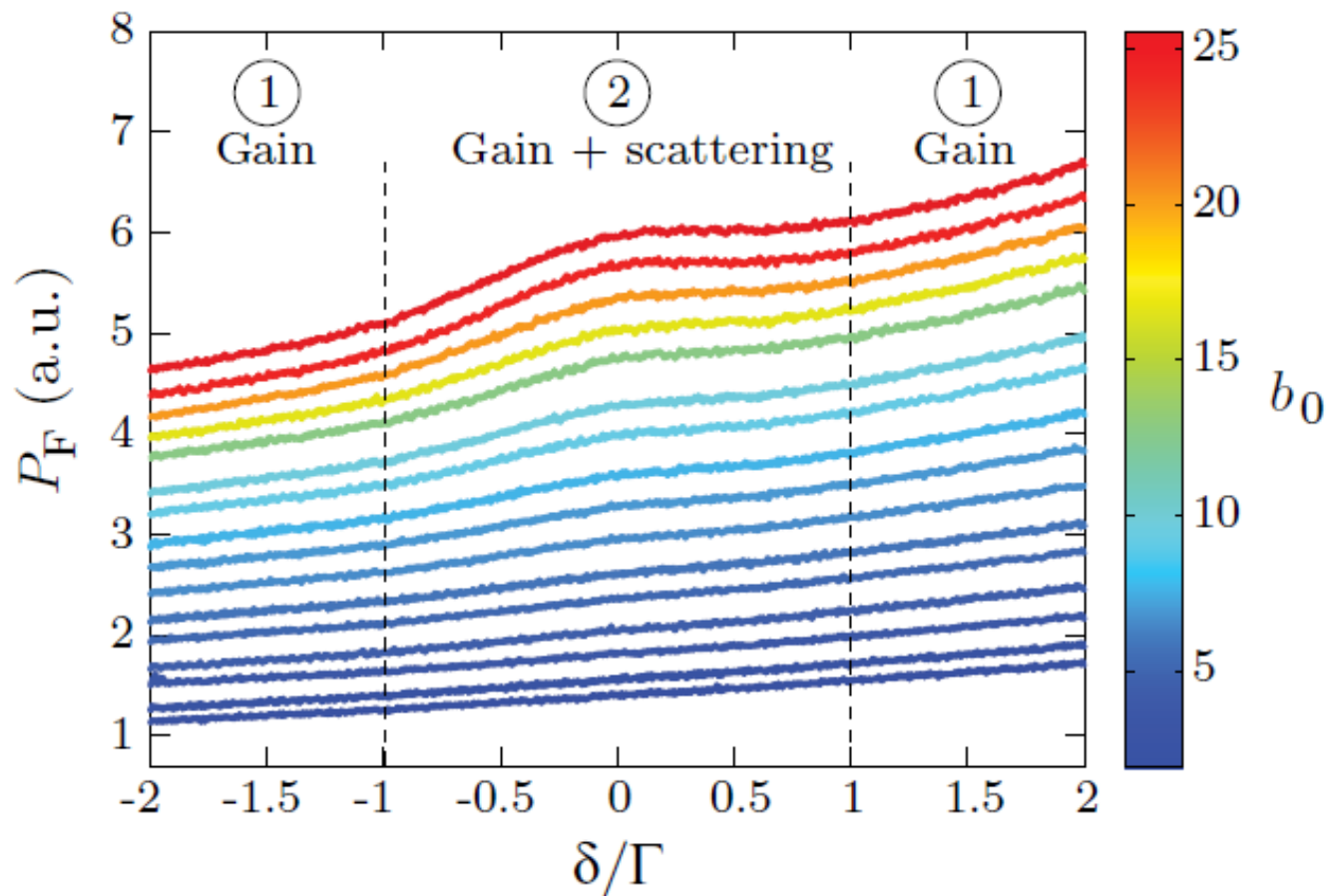
- The random laser emission:
  - is not spatially separated from elastic scattering from the external lasers
  - is very hard to spectrally separate
 → **We look at the total fluorescence** (= pump depletion)
  
- We change  $b_0$  (defines the threshold) with a **constant atom number**.
  - changes are only due to collective effects
  
- We sweep slowly (steady-state) the Raman laser (no probe) around the frequency where Raman gain is on resonance with the  $|2\rangle \rightarrow |1'\rangle$  transition.



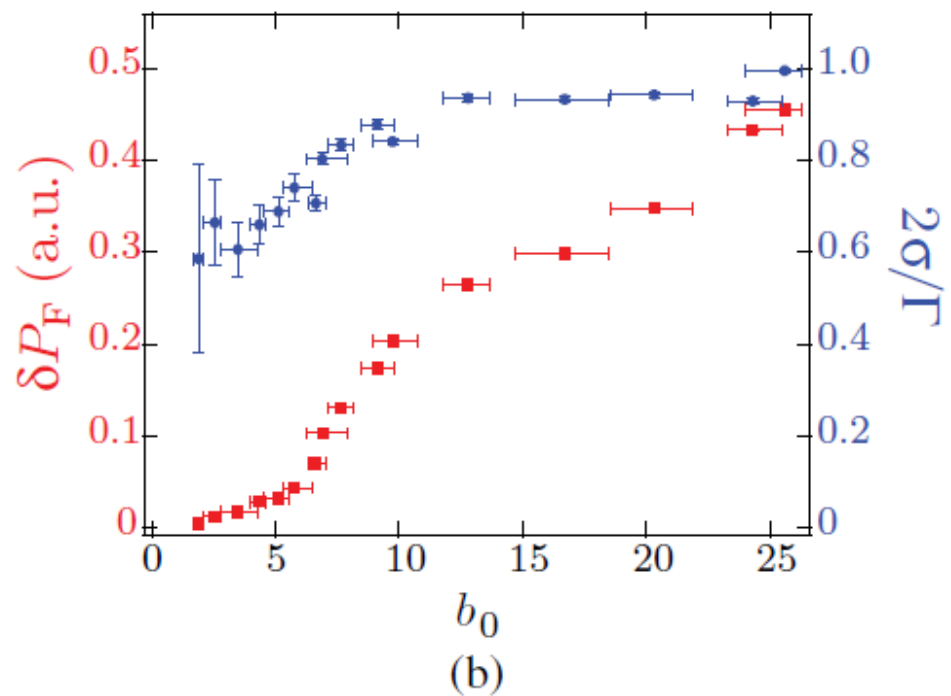
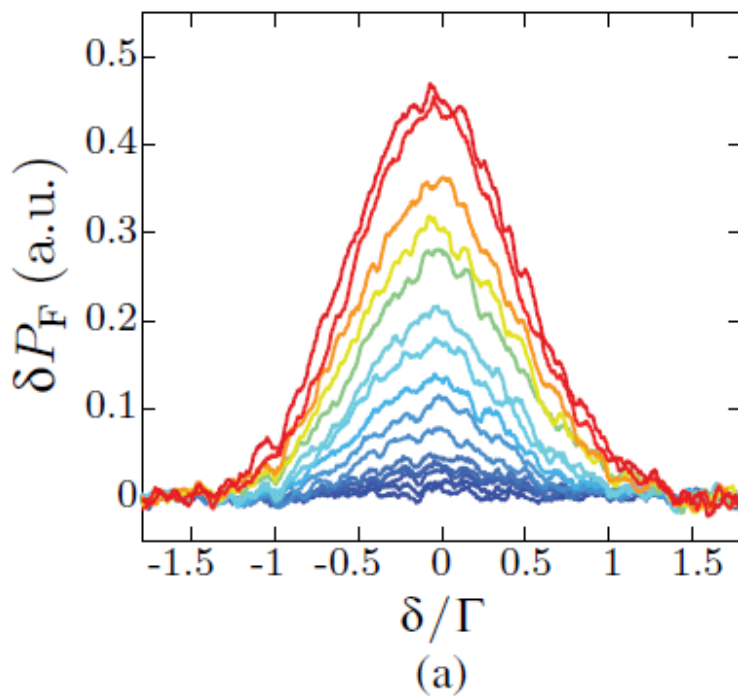
1- Overall increase of fluorescence  $\rightarrow$  Amplified spontaneous emission



- 1- Overall increase of fluorescence  $\rightarrow$  Amplified spontaneous emission
- 2- Increase of fluorescence around  $\delta = 0 \rightarrow$  combined effect of gain and multiple scattering

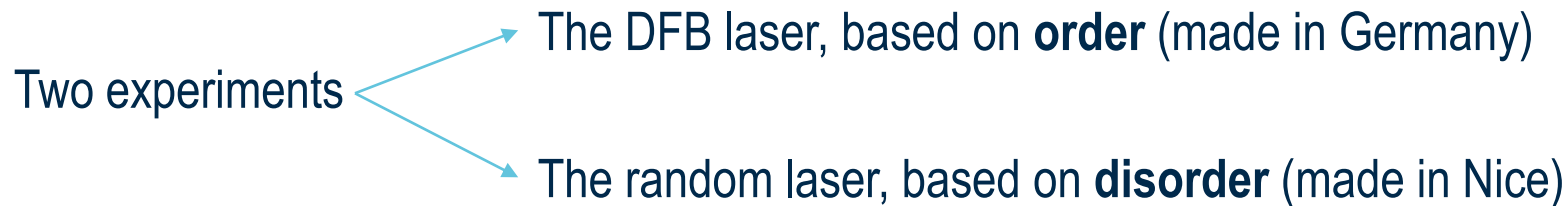


- Fit of the wings  $\rightarrow$  we can subtract the “ASE” background
- $\rightarrow$  More visible bump (Gaussian shape)
- $\rightarrow$  The amplitude has a **threshold with  $b_0$**



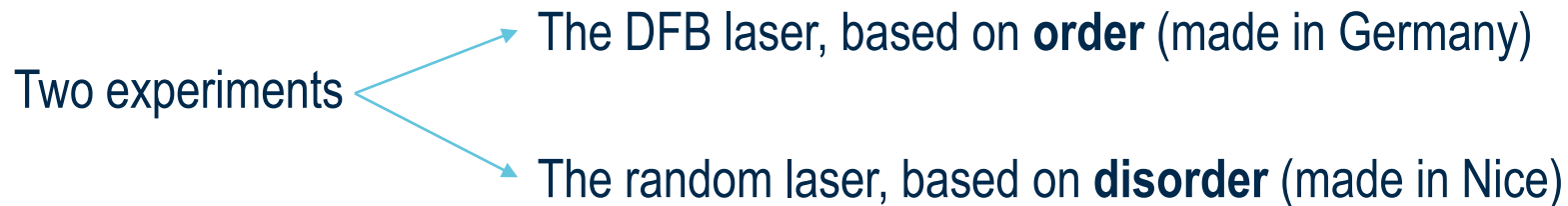
Baudouin *et al.*, Nature Phys. **9**, 357 (2013).

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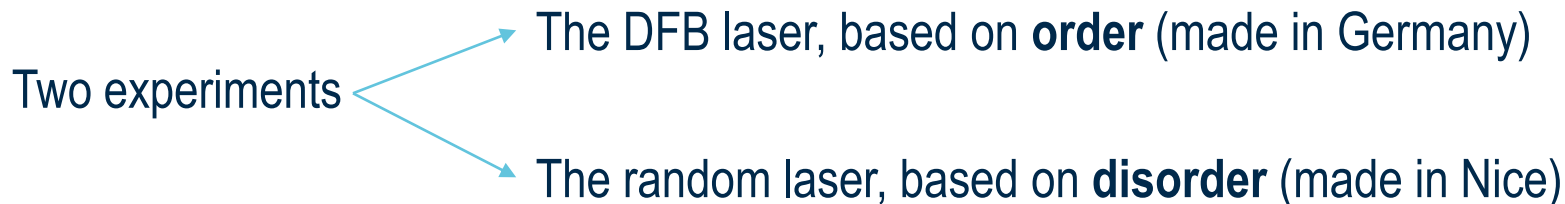
Which one was the *simplest* ?





Which one was the *simplest* ?    The DFB laser !    It took 6 months, the RL 4 years !





Which one was the *simplest* ?    The DFB laser !    It took 6 months, the RL 4 years !

Because **it's only 1D...** → easy to have many layers  
 → directional emission easy to detect

We might investigate the 3D case in the future...

The random laser experiment is far from being finished, we want:

- more data;
- different (more spectacular) signatures. Some spectral or coherence properties ?

→ PhD thesis of Samir Vartabi Kashani, INLN, on-going.

We made the first **mirrorless lasers based on cold atoms**

The whole laser is only made of a few millions atoms in a very dilute gas phase.

→ The *lightest* laser ever !  $M \sim 10$  fg.

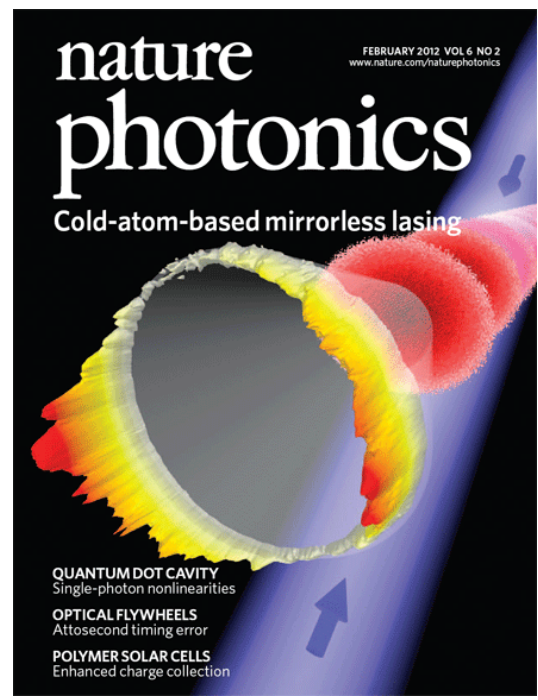
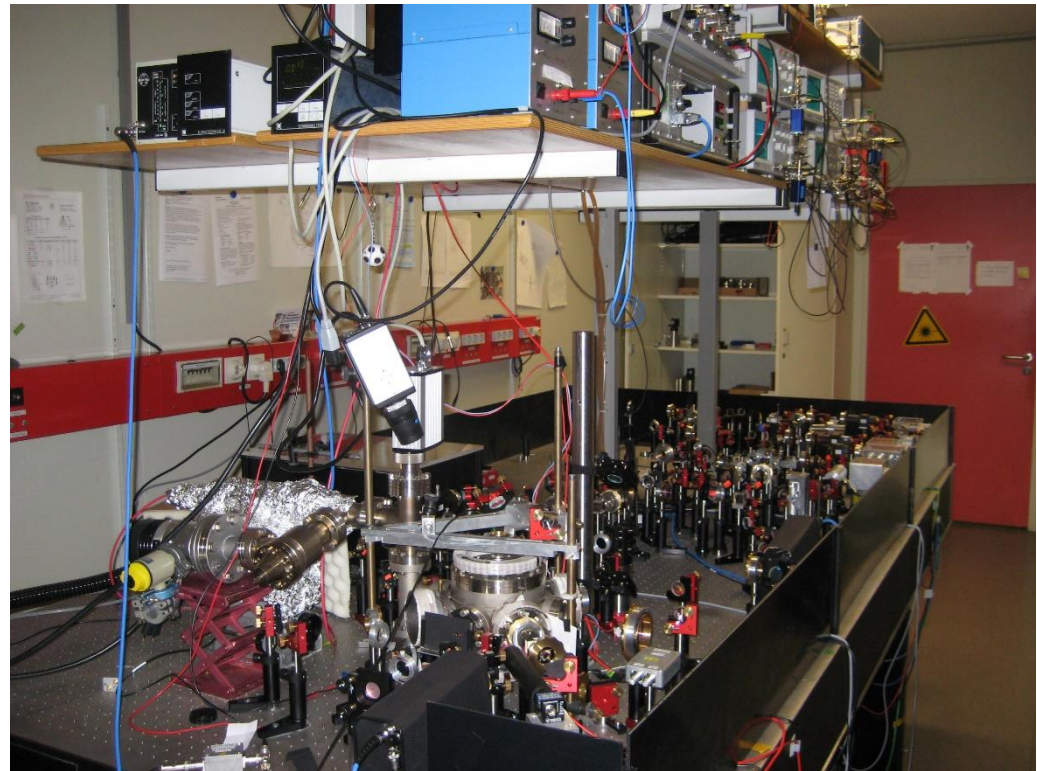


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But...



We made the first **mirrorless lasers with cold atoms**

The whole laser is only made of a few millions atoms in a very dilute gas phase.

→ The *lightest* laser ever !  $M \sim 10$  fg.

But...

There is a big, complex, and expensive machinery behind it...

And: no new  $\lambda$ , low power

→ limited practical interest ☹️



So, what is it interesting for ?

DFB lasers are well known and their physics understood.

But this one has a **cone-shaped emission**. This is **new** !

Why ?

- Because of the **sharp resonance**: in standard DFB laser, the emission wavelength adapts itself to the lattice periodicity, because the gain bandwidth is large.
- Because of the **high nonlinearity and versatility** of cold atoms: just retroreflecting the pumping beam makes a new gain mechanism appear (FWM), which makes the feedback with angle stable.



→ Very good illustration that applying known physics in a new system allows discoveries

Our random laser is not very convenient: hard to produce, hard to characterize...

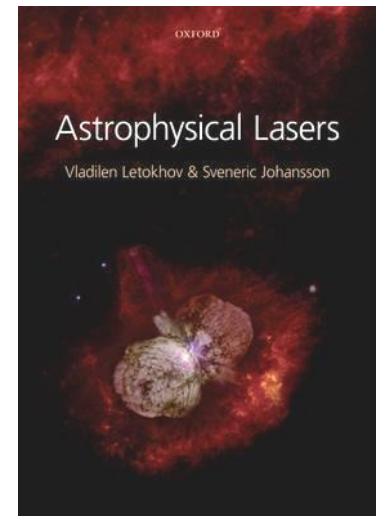
But it has unique features:

- It's truly 3D (homogeneous pumping) thanks to the sharp resonance.
  - The scatterers are all identical (monodisperse sample), and perfectly known, without absorption.
  - The average over the position configuration is done.
- Possible to develop *ab initio* models without any free parameters.

→ Perfect test-bed for theoreticians  
(on-going collaborations)

Also: the first RL based on atomic vapors. Extension to hot atoms ?

→ Would be closer to astrophysical systems (natural RL in space ?)





# Thank you for your attention