

ON THE EXISTENCE AND COMPACTNESS OF A TWO-DIMENSIONAL RESONANT SYSTEM OF CONSERVATION LAWS

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ABSTRACT. We prove the existence of a weak solution to a two-dimensional resonant 3×3 system of conservation laws with BV initial data. Due to possible resonance (coinciding eigenvalues), spatial BV estimates are in general not available. Instead, we use an entropy dissipation bound combined with the time translation invariance property of the system to prove existence based on a two-dimensional compensated compactness argument adapted from [36]. Existence is proved under the assumption that the flux functions in the two directions are linearly independent.

1. INTRODUCTION

This paper studies certain two-dimensional resonant 3×3 systems of conservation laws of the form

$$(1.1) \quad \begin{aligned} k_t &= 0, & l_t &= 0, \\ u_t + f(k, u)_x + g(l, u)_y &= 0, \end{aligned}$$

which are augmented with $L^\infty \cap BV$ initial data

$$(1.2) \quad k|_{t=0} = k(x, y), \quad l|_{t=0} = l(x, y), \quad u|_{t=0} = u_0(x, y).$$

The goal is to prove that there exists a weak solution to (1.1)–(1.2).

In recent years the one-dimensional version of the above system,

$$(1.3) \quad \begin{aligned} k_t &= 0, \\ u_t + f(k, u)_x &= 0, \end{aligned}$$

has received a considerable amount of attention. This system may be viewed as an alternative way of writing a scalar conservation law with a discontinuous flux, namely

$$(1.4) \quad u_t + f(k(x), u)_x = 0.$$

Equations like (1.4) occur in a variety of applications, including flow in porous media, sedimentation processes, traffic flow, radar shape-from-shading problems, blood flow, and gas flow in a variable duct.

If $k(x)$ is a smooth function, Kruřkov's theory [21] tells us that there exists a unique entropy solution to the initial value problem for (1.4), for general flux functions f . The scalar Kruřkov theory does not apply when $k(x)$ is discontinuous. Instead it proves useful to rewrite (1.4) as a 2×2 system of equations (1.3), which makes it possible to apply ideas from the theory of systems of conservation laws.

As a starting point, it is necessary to introduce conditions on the flux $f(k, u)$ that guarantee that solutions stay uniformly bounded. For example, one can require $f(k, a) = f(k, b) = 0$ for all k , which in fact implies that the interval $[a, b] \subset \mathbb{R}$ becomes an invariant region. The system (1.4) has two eigenvalues, namely $\lambda_1 = 0$ and $\lambda_2 = f_u(k, u)$. Consequently, if $f_u(k, u)$ vanishes for some value of (k, u) , then (1.4)

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is nonstrictly hyperbolic and experiences so-called nonlinear resonant behavior, which implies that wave interactions are more complicated than in strictly hyperbolic systems. As a matter of fact, one cannot expect to bound the total variation of the conserved quantities directly, but only when measured under a certain singular mapping. A singular mapping that is relevant for (1.3) is

$$\Psi(k, u) = \int^u |f_u(k, \xi)| d\xi.$$

If $\{u^\rho\}_{\rho>0}$ is a sequence of "reasonable" approximate solutions of (1.3), then one proves that the total variation of the transformed quantity $z^\rho := \Psi(k, u^\rho)$ is bounded independently of ρ . Helly's theorem then gives convergence (along a subsequence) of z^ρ as $\rho \downarrow 0$. Since the continuous mapping $u \mapsto \Psi(k, u)$ is one-to-one, u^ρ also converges.

A singular mapping was used first by Temple [39] to establish convergence of the Glimm scheme (and thereby the existence of a weak solution) for a 2×2 resonant system of conservation laws modeling the displacement of oil in a reservoir by water and polymer, which is now known to be equivalent to a conservation law with a discontinuous coefficient (see, e.g., [20]). Since then the singular mapping approach has been used and adapted by great many authors to prove existence of weak solutions to resonant systems of conservation laws/scalar conservation laws with discontinuous flux functions, by establishing convergence of various approximations schemes (Glimm and Godunov schemes, front tracking, upwind and central type schemes, vanishing viscosity/smoothing method, ...), see (the list is far from being complete) [1, 2, 3, 4, 5, 6, 8, 9, 10, 11, 12, 18, 19, 20, 22, 23, 29, 35, 40, 41]. Similar ideas have been used also in the context of degenerate parabolic equations [15]. Regarding uniqueness and entropy conditions for scalar conservation laws with discontinuous coefficients, see [16, 17] and the references therein.

As an alternative to the singular mapping approach, the papers [13, 14, 17] has suggested to use the compensated compactness method and "scalar entropies" for the convergence analysis of approximate solutions. The results obtained with this approach are more general (and to some extent the proofs are easier) than those obtained with the singular mapping approach.

All the papers up to now have addressed the one-dimensional case. The aim of the present paper is to take a first look at the multi-dimensional case, which is completely unexplored. More precisely, we will prove the existence of at least one weak solution to the initial value problem for the two-dimensional system (1.1).

Our existence proof is based on studying the " $(\varepsilon, \delta) \downarrow (0, 0)$ limit" of classical solutions $u^{\varepsilon, \delta}$ of the uniformly parabolic equation

$$u_t^{\varepsilon, \delta} + f(k^\delta, u^{\varepsilon, \delta})_x + g(l^\delta, u^{\varepsilon, \delta})_y = \varepsilon (u_{xx}^{\varepsilon, \delta} + u_{yy}^{\varepsilon, \delta}), \quad \varepsilon > 0, \delta > 0,$$

where k^δ, l^δ converge to k, l in $L^1_{\text{loc}}(\mathbb{R}^2)$, respectively, as $\delta \downarrow 0$.

Observe that we are essentially considering a scalar approximation scheme for (1.1), see [5, 6, 15, 13, 14, 17, 40, 41] for other scalar approximation schemes for one-dimensional discontinuous flux problems.

Although spatial BV bounds are out of reach, we still have a time translation invariance property at our disposal, which, together with the assumption of BV initial data, implies that $u_t^{\varepsilon, \delta}$ is uniformly bounded in L^1 . Consider three functions $F(k, u), G(l, u), H(k, l, u)$ defined by

$$F_u = (f_u)^2, \quad G_u = (g_u)^2, \quad H_u = f_u g_u.$$

We prove, at least under the assumption that ε and δ are of comparable size, that the two sequences

$$F(k(x, y), u^{\varepsilon, \delta})_x + H(k(x, y), l(x, y), u^{\varepsilon, \delta})_y$$

and

$$H(k(x, y), l(x, y), u^{\varepsilon, \delta})_x + G(k(x, y), u^{\varepsilon, \delta})_y$$

are compact in $W_{\text{loc}}^{-1,2}(\mathbb{R}^2)$, for each fixed $t > 0$.

The crux of the convergence analysis is then to prove that the above $W_{\text{loc}}^{-1,2}(\mathbb{R}^2)$ compactness is sufficient to establish a "two-dimensional" compensated compactness argument in the spirit of the classical Tartar-Murat results for one-dimensional conservation laws, [26, 27, 28, 37, 38]. Here we follow the recent two-dimensional compensated compactness framework developed in Tadmor et. al. [36] for *non-linear* conservation laws. We extend their results to the case involving additional discontinuous "variable coefficients". Accordingly, we make the *nonlinearity assumption* that for each fixed k, l the functions $u \mapsto f_u(k, u)$ and $u \mapsto g_u(l, u)$ are almost everywhere linearly independent (see (2.4) in the next section for a precise statement). Our main existence result is based on an application of the two-dimensional compensated compactness lemma with "variable coefficients" — lemma 3.2 stated in Section 3 below. Granted the nonlinearity assumption, it then yields that (a subsequence of) $u^{\varepsilon, \delta}(\cdot, \cdot, t)$ converges in $L_{\text{loc}}^1(\mathbb{R}^2)$ to a bounded function $u(\cdot, \cdot, t)$, for a.e. $t > 0$. Since $u^{\varepsilon, \delta}$ is uniformly L^1 Lipschitz continuous in time we obtain, in Section 4 below, our main Theorem 2.1, stating that $u^{\varepsilon, \delta} \rightarrow u$ in $L_{\text{loc}}^1(\mathbb{R}^2 \times \mathbb{R}_+)$ and that the limit function u is a weak solution of (1.1)–(1.2).

Although we have chosen to analyze the vanishing viscosity/smoothing method, the techniques used here for that purpose can also be applied to various numerical schemes, including appropriate two-dimensional versions of the scalar finite difference schemes studied in [13, 17, 40, 41].

2. ASSUMPTIONS AND STATEMENT OF MAIN RESULTS

We start by listing the assumptions on the initial conditions u_0 and the fluxes k, l, f, g that are needed for the existence result.

Regarding the initial function we assume

$$(2.1) \quad u_0 \in L^\infty(\mathbb{R}^2) \cap BV(\mathbb{R}^2), \quad a \leq u_0 \leq b \quad \text{for a.e. in } \mathbb{R}^2.$$

For the discontinuous coefficients $k, l : \mathbb{R}^2 \rightarrow \mathbb{R}$ we assume

$$(2.2) \quad \begin{cases} k, l \in L^\infty(\mathbb{R}^2) \cap BV(\mathbb{R}^2), \\ \alpha \leq k, l \leq \beta \quad \text{a.e. in } \mathbb{R}^2. \end{cases}$$

For the flux functions $f, g : [\alpha, \beta] \times [a, b] \rightarrow \mathbb{R}$ we assume

$$(2.3) \quad \begin{cases} u \mapsto f(k, u), u \mapsto g(l, u) \in C^2[a, b] \text{ for all } k, l \in [\alpha, \beta]; \\ k \mapsto f(k, u), l \mapsto g(l, u) \in C^1[\alpha, \beta] \text{ for all } u \in [a, b]. \end{cases}$$

Moreover, we make the *nonlinearity assumption* which excludes the possibility of $\xi_1 f(k, u) + \xi_2 g(l, u)$ being an affine function (in u) on any nontrivial interval for all $k, l \in [\alpha, \beta]$,

$$(2.4a) \quad \forall |\xi| = 1 \text{ and } k, l \in [\alpha, \beta] : \quad \xi_1 f(k, \cdot) + \xi_2 g(l, \cdot) \not\equiv \text{affine function on any nontrivial interval.}$$

In its slightly stronger version, this assumption requires that $f_u(k, \cdot)$ and $g_u(l, \cdot)$ are a.e. linearly independent so that the symbol $s(\xi, k, l, u) := \xi_1 f_u(k, u) + \xi_2 g_u(l, u)$ satisfies

$$(2.4b) \quad \forall |\xi| = 1 \text{ and } k, l \in [\alpha, \beta] : \quad \text{meas}\{u \mid s(\xi, k, l, u) = 0\} = 0.$$

This is a straightforward generalization of the notion of nonlinearity found in [24], in their study of kinetic formulations for nonlinear conservation laws.

Finally, we need to know that our approximate solutions stay uniformly bounded. For example, this is ensured by the assumption

$$(2.5) \quad f(k, a), f(k, b), g(l, a), g(l, b) = 0 \text{ for all } k, l \in [\alpha, \beta],$$

which implies that the interval $[a, b]$ becomes an invariant region.

We are now ready to state our main result.

Theorem 2.1. *Suppose (2.1), (2.2), (2.3), (2.4), and (2.5) hold. Then, there exists a weak solution of the initial value problem (1.1)–(1.2), $u \in L^\infty(\mathbb{R}^2 \times \mathbb{R}_+) \cap \text{Lip}(\mathbb{R}_+; L^1(\mathbb{R}^2))$, satisfying*

$$(2.6) \quad \int_{\mathbb{R}_+} \int_{\mathbb{R}} \left(u \phi_t + f(k(x, y), u) \phi_x + g(l(x, y), u) \phi_y \right) dx dt + \int_{\mathbb{R}} u_0(x) \phi(x, 0) dx = 0, \quad \forall \phi \in \mathcal{D}(\mathbb{R}^2 \times [0, \infty)).$$

This weak solution can be constructed as a strong $L^1_{\text{loc}}(\mathbb{R}^2 \times \mathbb{R}_+)$ -limit of classical solutions, $u^{\varepsilon, \delta}$, of the uniformly parabolic problems,

$$u_t^{\varepsilon, \delta} + f(k^\delta, u^{\varepsilon, \delta})_x + g(l^\delta, u^{\varepsilon, \delta})_y = \varepsilon \Delta u^{\varepsilon, \delta},$$

with the smoothly mollified coefficients, $k^\delta := \omega_\delta \star k$ and $l^\delta := \omega_\delta \star l$ (outlined in section 4 below). Moreover, if v is another weak solution constructed as a vanishing viscosity limit of $v^{\varepsilon, \delta}$, then for any $t > 0$,

$$(2.7) \quad \|u(\cdot, t) - v(\cdot, t)\|_{L^1(\mathbb{R}^2)} \leq \|u_0 - v_0\|_{L^1(\mathbb{R}^2)}.$$

Thus, if we let S_t denote the solution operator associated with weak solutions constructed as such vanishing viscosity limits, then S_t is L^1 -contractive.

The proof of this theorem is given in the following two sections. We close this section with the following summary.

Remark 2.2. Stated differently, Theorem 2.1 shows that there exists a weak solution to the following two-dimensional scalar conservation law with discontinuous coefficients $k, l \in L^\infty(\mathbb{R}^2) \cap BV(\mathbb{R}^2)$:

$$\begin{aligned} u_t + f(k(x, y), u)_x + g(l(x, y), u)_y &= 0, \\ u|_{t=0} &= u_0 \in L^\infty(\mathbb{R}^2) \cap BV(\mathbb{R}^2). \end{aligned}$$

We note in passing that the solution operator in this case of discontinuous "variable coefficients" is not translation invariant in space and hence, the L^1 -contraction statement in (2.7) does not imply spatial BV compactness.

Moreover, if we let $S_t : L^\infty(\mathbb{R}) \rightarrow L^\infty(\mathbb{R})$ denote the solution operator of (1.1)–(1.2), so that $u(t, \cdot) = S_t u_0(\cdot)$ is the (vanishing viscosity) weak solution constructed in theorem 2.1, then by adapting standard arguments, we can prove that $\{S_t\}_{t>0}$ is compact with respect to the L^1_{loc} norm.

3. A COMPENSATED COMPACTNESS LEMMA

In this section we prove a "two-dimensional" compensated compactness lemma. We refer [7, 25, 26, 27, 28, 37, 38] for background information on the compensated compactness theory. We start by recalling the celebrated div-curl lemma.

Lemma 3.1 (div-curl). *Let $\Omega \subset \mathbb{R}^2$ be an open domain and let $\rho > 0$ denote a parameter taking its values in a sequence which tends to zero. Suppose $D^\rho \rightharpoonup \overline{D}$, $E^\rho \rightharpoonup \overline{E}$ in $(L^2(\Omega))^2$ and $\{\text{div } D^\rho\}_{\rho>0}$, $\{\text{curl } E^\rho\}_{\rho>0}$ belong to a compact subset of $W_{\text{loc}}^{-1,2}(\Omega)$. Then, after extracting a subsequence if necessary, we have $D^\rho \cdot E^\rho \rightarrow \overline{D} \cdot \overline{E}$ in $\mathcal{D}'(\Omega)$ as $\rho \downarrow 0$.*

The compensated compactness lemma below is tailored for two-dimensional equations, whose spatial part involve discontinuous coefficients:

$$f(k(x, y), v(x, y))_x + g(l(x, y), v(x, y))_y.$$

If $g(l, u) = g(u)$ and $f(k, u) = f(u)$, then the lemma below coincides with the two-dimensional result of [36, Theorem 3.1]. If we set $g = 0$ then the result coincides with Tartar's compensated compactness lemma for the one-dimensional scalar conservation with genuinely nonlinear flux f .

Lemma 3.2 (Compensated compactness). *Let $\Omega \subset \mathbb{R}^2$ be an open domain. Let k, l, f, g be functions satisfying (2.2), (2.3), (2.4), and (2.5). Suppose $\{v^\rho(x, y)\}_{\rho>0}$ is a sequence of measurable functions on Ω that satisfies the following two conditions:*

1. *There exist two finite constants $a < b$ independent of ρ such that*

$$a \leq v^\rho(x, y) \leq b \quad \text{for a.e. } (x, y) \in \Omega.$$

2. *Let the functions F, G, H be defined by*

$$F_u(k, u) = (f_u(k, u))^2, \quad G_u(l, u) = (g_u(l, u))^2, \quad H_u(k, l, u) = f_u(k, u)g_u(l, u).$$

We assume that the two sequences

$$(3.1) \quad \begin{aligned} & \left\{ F(k(x, y), v^\rho)_x + H(k(x, y), l(x, y), v^\rho)_y \right\}_{\rho>0}, \\ & \left\{ H(k(x, y), l(x, y), v^\rho)_x + G(l(x, y), v^\rho)_y \right\}_{\rho>0} \end{aligned}$$

belong to a compact subset of $W_{\text{loc}}^{-1,2}(\mathbb{R} \times \mathbb{R}_+)$.

Then, there exists a subsequence of $\{v^\rho(x, y)\}_{\rho>0}$ that converges a.e. to a function $v \in L^\infty(\mathbb{R}^2)$, and $a \leq v(x, y) \leq b$ for a.e. $(x, y) \in \Omega$.

Proof. To simplify the notation let

$$F^\rho := F(k(x, y), v^\rho), \quad G^\rho := G(l(x, y), v^\rho), \quad H^\rho := H(k(x, y), l(x, y), v^\rho),$$

and denote their $L^\infty(\mathbb{R} \times \mathbb{R}_+)$ weak- \star limits by $\overline{F}, \overline{G}, \overline{H}$, respectively. Introduce the vector fields

$$D^\rho = (F^\rho, H^\rho) \quad E^\rho = (-G^\rho, H^\rho),$$

and denote their respective $L^\infty(\mathbb{R} \times \mathbb{R}_+)$ weak- \star limits by $\overline{D}, \overline{E}$.

Thanks to (3.1), we can apply the div-curl lemma to the sequences $\{D^\rho\}_{\rho>0}, \{E^\rho\}_{\rho>0}$ to produce

$$\overline{D^\rho \cdot E^\rho} = \overline{D} \cdot \overline{E} \quad \text{a.e. in } \Omega,$$

that is,

$$\overline{(H^\rho)^2 - F^\rho G^\rho} = (\overline{H})^2 - \overline{F} \overline{G},$$

which implies

$$(3.2) \quad \overline{(H^\rho - \overline{H})^2 - (F^\rho - \overline{F})(G^\rho - \overline{G})} = 0.$$

Fix $c = c(x, y) \in L^\infty(\Omega)$. Following [36], we now consider the function $I : [a, b] \rightarrow \mathbb{R}$ defined by

$$\begin{aligned} I(v) &= \left(H(k(x, y), l(x, y), v) - H(k(x, y), l(x, y), c) \right)^2 \\ &\quad - \left(F(k(x, y), v) - F(k(x, y), c) \right) \cdot \left(G(l(x, y), v) - G(l(x, y), c) \right). \end{aligned}$$

Note that

$$\begin{aligned} I(v^\rho) &= \left([H^\rho - \overline{H}] + [\overline{H} - H(k(x, y), l(x, y), c)] \right)^2 \\ &\quad - \left([F^\rho - \overline{F}] + [\overline{F} - F(k(x, y), c)] \right) \cdot \left([G^\rho - \overline{G}] + [\overline{G} - G(l(x, y), c)] \right). \end{aligned}$$

Using this and (3.2), we compute

$$(3.3) \quad \overline{I(v^\rho)} = \left(\overline{H} - H(k(x, y), l(x, y), c) \right)^2 - \left(\overline{F} - F(k(x, y), c) \right) \cdot \left(\overline{G} - G(l(x, y), c) \right).$$

By the Cauchy-Schwarz inequality we have for any $u \in [a, b]$

$$(3.4) \quad \begin{aligned} & \left(H(k(x, y), l(x, y), u) - H(k(x, y), l(x, y), c) \right)^2 \\ &= \left(\int_c^u f_u(k(x, y), \xi) g_u(l(x, y), \xi) d\xi \right)^2 \\ &\leq \left(F(k(x, y), u) - F(k(x, y), c) \right) \left(G(l(x, y), u) - G(l(x, y), c) \right), \end{aligned}$$

and hence $I(\cdot) \leq 0$ with $I(c) = 0$. Thanks to (2.4), the Cauchy-Schwarz inequality in (3.4) is in fact a *strict* inequality. This shows that the function $I(v)$ has a strict global maximum at $v = c$.

Since $u \mapsto F(k(x, y), u)$ is strictly increasing, we can choose c as

$$c(x, y) := F^{-1}(k(x, y), \overline{F}(x, y)),$$

so that (3.3) becomes

$$(3.5) \quad \overline{I(v^\rho)} = \left(\overline{H} - H(k(x, y), l(x, y), c) \right)^2.$$

Since $I(\cdot) \leq 0$, we conclude that $\overline{H} = H(k(x, y), l(x, y), c)$, and thus $\overline{I(v^\rho)} = 0$. In fact, we have $I(v^\rho) \rightarrow 0$ a.e. in Ω .

Using the fact that c is a strict maximizer of $I(v)$, we have

$$I(v) \leq -C_\alpha, \quad \text{whenever } |v - c| > \alpha,$$

for some constant $C_\alpha > 0$ that depends on α . Consequently,

$$(3.6) \quad \text{meas} \{ |v^\rho - \overline{v}| > \alpha \} \leq \frac{-1}{C_\alpha} \iint_{\Omega \cap \{|v^\rho - c| > \alpha\}} I(v^\rho(x, y)) dx dy \rightarrow 0 \quad \text{as } \rho \downarrow 0.$$

Since $\alpha > 0$ was arbitrary, $v^\rho \rightarrow c$ in measure, which in turn implies that a subsequence of $\{v^\rho\}_{\delta > 0}$ converges to c a.e. in Ω . \square

We remark that the idea of using the Cauchy-Schwarz inequality along the lines of (3.4) for proving strong compactness can be traced back to [33, 34]

4. PROOF OF THEOREM 2.1

Let $k^\delta, l^\delta, u_0^\delta$ be smooth functions converging strongly to k, l, u_0 respectively. More precisely, let $\omega_\delta \in C_0^\infty(\mathbb{R})$ be a nonnegative function satisfying

$$\omega(x) \equiv 0 \quad \text{for } |x| \geq 1, \quad \iint_{\mathbb{R}^2} \omega(x) dx = 1.$$

For $\delta > 0$, let $\omega_\delta(x) = \frac{1}{\delta^2} \omega\left(\frac{x}{\delta}\right)$ and introduce the mollified functions

$$k^\delta = \omega_\delta \star k, \quad l^\delta = \omega_\delta \star l.$$

We approximate the initial data u_0 by cut-off and mollification as follows:

$$u_0^\delta = \omega_\delta \star (u_0 \chi_\mu),$$

where $\chi_\delta(x) = 1$ for $|x| \leq 1/\delta$ and 0 otherwise. In particular, we have the estimate

$$(4.1) \quad \left\| (u_0^\delta)_{xx} + (u_0^\delta)_{yy} \right\|_{L^1(\mathbb{R}^2)} \leq \frac{1}{\delta} \iint_{\mathbb{R}^2} \left(|(u_0^\delta)_x| + |(u_0^\delta)_y| \right) dx dy \leq \frac{1}{\delta} |u_0|_{BV(\mathbb{R}^2)}.$$

Observe that for $h^\delta = k^\delta, l^\delta, u_0^\delta$ and $h = k, l, u_0$, we have $h^\delta \in C^\infty(\mathbb{R}^2)$ and

$$h^\delta \rightarrow h \quad \text{a.e. in } \mathbb{R}^2 \text{ and in } L^p(\mathbb{R}^2) \text{ for any } p \in [1, \infty) \text{ as } \delta \downarrow 0.$$

Additionally, u_0^δ is compactly supported.

Now let $u^{\varepsilon, \delta}$ be the unique classical solution of the uniformly parabolic equation

$$(4.2) \quad u_t^{\varepsilon, \delta} + f(k^\delta, u^{\varepsilon, \delta})_x + g(l^\delta, u^{\varepsilon, \delta})_y = \varepsilon (u_{xx}^{\varepsilon, \delta} + u_{yy}^{\varepsilon, \delta}),$$

with initial data $u^{\varepsilon, \delta}|_{t=0} = u_0^\delta$. The proof proceeds through a series of lemmas, which in the end show that for each $t \in [0, T]$ a subsequence of $u^{\varepsilon, \delta}(\cdot, \cdot, t)$ converges a.e. as $\varepsilon, \delta \downarrow 0$.

Our first lemma confirms the uniform bound.

Lemma 4.1 (L^∞ bound). *There exists a constant $C > 0$, independent of ε, δ , such that*

$$\|u^{\varepsilon, \delta}(\cdot, \cdot, t)\|_{L^\infty(\mathbb{R}^2)} \leq C, \quad \text{for all } t \in (0, T).$$

Proof. The proof is standard and exploits assumption (2.5) to conclude that $a \leq u^{\varepsilon, \delta}(x, y, t) \leq b$ for a.e. $(x, y, t) \in \mathbb{R}^2 \times \mathbb{R}_+$. \square

Using that (4.2) is translation invariant in time, we can prove that $u_t^{\varepsilon, \delta}$ is uniformly bounded in $L^\infty(\mathbb{R}_+; L^1(\mathbb{R}^2))$.

Lemma 4.2 (Lipschitz regularity in time). *Suppose the two smoothing parameters ε and δ are kept in balance in the sense that*

$$(4.3) \quad \delta = C\varepsilon, \quad \text{for some constant } C > 0.$$

There is a constant C_0 (which is possibly dependent on u_0 but otherwise is independent of ε, δ), such that for any $t > 0$

$$\iint_{\mathbb{R}^2} |\partial_t u^{\varepsilon, \delta}(\cdot, \cdot, t)| \, dx \, dy \leq C_0.$$

Proof. To prove this, set $w^{\varepsilon, \delta} = u_t^{\varepsilon, \delta}$. Then w satisfies

$$w_t^{\varepsilon, \delta} + (f_u(k^\delta, u^{\varepsilon, \delta})w^{\varepsilon, \delta})_x + (g_u(l^\delta, u^{\varepsilon, \delta})w^{\varepsilon, \delta})_y = \varepsilon (w_{xx}^{\varepsilon, \delta} + w_{yy}^{\varepsilon, \delta}).$$

Multiplying by $\text{sign}(w^{\varepsilon, \delta})$ gives, in the sense of distributions,

$$\begin{aligned} & |w^{\varepsilon, \delta}|_t + (f_u(k^\delta, u^{\varepsilon, \delta})|w^{\varepsilon, \delta}|)_x + (g_u(l^\delta, u^{\varepsilon, \delta})|w^{\varepsilon, \delta}|)_y \\ &= \varepsilon \left(|w^{\varepsilon, \delta}|_{xx} + |w^{\varepsilon, \delta}|_{yy} \right) - \varepsilon \text{sign}'(w^{\varepsilon, \delta}) \left((w_x^{\varepsilon, \delta})^2 + (w_y^{\varepsilon, \delta})^2 \right), \end{aligned}$$

since $f_u(k^\delta, u^{\varepsilon, \delta})w^{\varepsilon, \delta} \text{sign}'(w^{\varepsilon, \delta})w_x^{\varepsilon, \delta} = g_u(l^\delta, u^{\varepsilon, \delta})w^{\varepsilon, \delta} \text{sign}'(w^{\varepsilon, \delta})w_y^{\varepsilon, \delta} = 0$. Hence

$$\frac{d}{dt} \iint_{\mathbb{R}^2} |w^{\varepsilon, \delta}|(x, y, \cdot) \, dx \, dy \leq 0,$$

which, due to (4.1) and (4.3), concludes the proof. \square

Thanks to the previous lemma, we also have uniform $L^\infty(\mathbb{R}_+; L^2(\mathbb{R}^2))$ control over $\sqrt{\varepsilon}u_x^{\varepsilon, \delta}$ and $\sqrt{\varepsilon}u_y^{\varepsilon, \delta}$.

Lemma 4.3 (Entropy dissipation bound). *There is a constant C , independent of ε, δ , such that*

$$\varepsilon \iint_{\mathbb{R}^2} \left((u_x^{\varepsilon, \delta}(\cdot, \cdot, t))^2 + (u_y^{\varepsilon, \delta}(\cdot, \cdot, t))^2 \right) \, dx \, dy \leq C, \quad \text{for any } t > 0.$$

Proof. Multiplying (4.2) by $u^{\varepsilon, \delta}$ and then integrating yield

$$\begin{aligned} & \iint_{\mathbb{R}^2} \varepsilon \left((u_x^{\varepsilon, \delta}(\cdot, \cdot, t))^2 + (u_y^{\varepsilon, \delta}(\cdot, \cdot, t))^2 \right) dx dy \\ &= - \iint_{\mathbb{R}^2} uu_t^{\varepsilon, \delta} dx dy \\ &+ \iint_{\mathbb{R}^2} \left[\left(\int_0^{u^{\varepsilon, \delta}} f_u(k^\delta, \xi) d\xi \right)_x - \left(\int_0^{u^{\varepsilon, \delta}} f_k(k^\delta, \xi) d\xi \right) k_x^\delta(x, y) \right] dx dy \\ &+ \iint_{\mathbb{R}^2} \left[\left(\int_0^{u^{\varepsilon, \delta}} g_u(l^\delta, \xi) d\xi \right)_y - \left(\int_0^{u^{\varepsilon, \delta}} g_l(l^\delta, \xi) d\xi \right) l_y^\delta \right] dx dy. \end{aligned}$$

In view of Lemmas 4.1 and 4.2 and the BV regularity of the coefficients, we derive easily the uniform bound

$$\begin{aligned} & \varepsilon \iint_{\mathbb{R}^2} \left((u_x^{\varepsilon, \delta}(\cdot, \cdot, t))^2 + (u_y^{\varepsilon, \delta}(\cdot, \cdot, t))^2 \right) dx dy \\ & \leq C \left(\|u_t^{\varepsilon, \delta}\|_{L^\infty(\mathbb{R}_+; L^1(\mathbb{R}^2))} + |k|_{BV(\mathbb{R}^2)} + |l|_{BV(\mathbb{R}^2)} \right), \quad t > 0, \end{aligned}$$

for some constant C that is dependent on $\|u^{\varepsilon, \delta}\|_{L^\infty(\mathbb{R}^2)}$ but otherwise is independent of ε, δ . \square

Lemma 4.4 (Pre-compactness at each time instant). *Suppose the two parameters ε and δ are kept in balance in the sense that (4.3) holds. With F, G , and H defined in Lemma 3.2, the two sequences*

$$\begin{aligned} & \left\{ F(k(x, y), u^{\varepsilon, \delta})_x + H(k(x, y), l(x, y), u^{\varepsilon, \delta})_y \right\}_{\varepsilon, \delta > 0}, \\ & \left\{ H(k(x, y), l(x, y), u^{\varepsilon, \delta})_x + G(l(x, y), u^{\varepsilon, \delta})_y \right\}_{\varepsilon, \delta > 0} \end{aligned}$$

then belong to a compact subset of $W_{\text{loc}}^{-1,2}(\mathbb{R}^2)$, for each fixed $t > 0$.

Proof. Let $\phi = \phi(x, y) \in \mathcal{D}(\mathbb{R}^2)$, and, for each fixed $t > 0$, introduce the distribution

$$\langle \mathcal{L}^{\varepsilon, \delta}, \phi \rangle = \iint_{\mathbb{R}^2} \left(F(k(x, y), u^{\varepsilon, \delta}) \phi_x + H(k(x, y), l(x, y), u^{\varepsilon, \delta}) \phi_y \right) dx dy.$$

Let us first write $\mathcal{L}^{\varepsilon, \delta} = \mathcal{L}_1^{\varepsilon, \delta} + \mathcal{L}_2^{\varepsilon, \delta}$, where

$$\begin{aligned} \langle \mathcal{L}_1^{\varepsilon, \delta}, \phi \rangle &= \iint_{\mathbb{R}^2} \left(F(k(x, y), u^{\varepsilon, \delta}) - F(k^\delta(x, y), u^{\varepsilon, \delta}) \right) \phi_x dx dy \\ &+ \iint_{\mathbb{R}^2} \left(H(k(x, y), l(x, y), u^{\varepsilon, \delta}) - H(k^\delta(x, y), l^\delta(x, y), u^{\varepsilon, \delta}) \right) \phi_y dx dy, \\ \langle \mathcal{L}_2^{\varepsilon, \delta}, \phi \rangle &= \iint_{\mathbb{R}^2} \left(F(k^\delta(x, y), u^{\varepsilon, \delta}) \phi_x + H(k^\delta(x, y), l^\delta(x, y), u^{\varepsilon, \delta}) \phi_y \right) dx dy. \end{aligned}$$

In what follows, we let Ω denote an arbitrary but fixed bounded open subset of \mathbb{R}^2 . With $\phi \in W_0^{1,2}(\Omega)$, we have by Hölder's inequality

$$\left| \langle \mathcal{L}_1^{\varepsilon, \delta}, \phi \rangle \right| \leq C \left(\|k - k^\delta\|_{L^2(\Omega)} + \|l - l^\delta\|_{L^2(\Omega)} \right) \|\phi\|_{W_0^{1,2}(\Omega)} \rightarrow 0,$$

as $\delta \downarrow 0$. Thus, $\left\{ \mathcal{L}_1^{\varepsilon, \delta} \right\}_{\varepsilon, \delta > 0}$ is compact in $W^{-1,2}(\Omega)$, for each fixed t .

Multiplying (4.2) by $f_u(k^\delta(x, y), u^{\varepsilon, \delta})$ yields

$$f(k^\delta, u^{\varepsilon, \delta})_t + F(k^\delta, u^{\varepsilon, \delta})_x + H(k^\delta, l^\delta, u^{\varepsilon, \delta})_y = I_1^{\varepsilon, \delta} + I_2^{\varepsilon, \delta} + I_3^{\varepsilon, \delta} + I_4^{\varepsilon, \delta} + I_5^{\varepsilon, \delta},$$

where

$$\begin{aligned} I_1^{\varepsilon, \delta} &= (\varepsilon u_x^{\varepsilon, \delta} f_u(k^\delta, u^{\varepsilon, \delta}))_x + (\varepsilon u_y^{\varepsilon, \delta} f_u(k^\delta, u^{\varepsilon, \delta}))_y, \\ I_2^{\varepsilon, \delta} &= -\varepsilon (u_x^{\varepsilon, \delta})^2 f_{uu}(k^\delta, u^{\varepsilon, \delta}) - \varepsilon (u_y^{\varepsilon, \delta})^2 f_{uu}(k^\delta, u^{\varepsilon, \delta}), \\ I_3^{\varepsilon, \delta} &= -\varepsilon u_x^{\varepsilon, \delta} f_{uk}(k^\delta, u^{\varepsilon, \delta}) k_x^\delta - \varepsilon u_y^{\varepsilon, \delta} f_{uk}(k^\delta, u^{\varepsilon, \delta}) k_y^\delta, \\ I_4^{\varepsilon, \delta} &= F_k(k^\delta, u^{\varepsilon, \delta}) k_x^\delta + H_k(k^\delta, l^\delta, u^{\varepsilon, \delta}) k_y^\delta + H_l(k^\delta, l^\delta, u^{\varepsilon, \delta}) l_y^\delta \\ I_5^{\varepsilon, \delta} &= -f_u(k^\delta, u^{\varepsilon, \delta}) f_k(k^\delta, u^{\varepsilon, \delta}) k_x^\delta - f_u(k^\delta, u^{\varepsilon, \delta}) g_l(l^\delta, u^{\varepsilon, \delta}) l_y^\delta. \end{aligned}$$

Hence, there is a natural decomposition of $\mathcal{L}_{2,0}^{\varepsilon, \delta}$ into six parts. We name the six parts $\mathcal{L}_{2,0}^{\varepsilon, \delta}$, $\mathcal{L}_{2,1}^{\varepsilon, \delta}$, $\mathcal{L}_{2,2}^{\varepsilon, \delta}$, $\mathcal{L}_{2,3}^{\varepsilon, \delta}$, $\mathcal{L}_{2,4}^{\varepsilon, \delta}$, and $\mathcal{L}_{2,5}^{\varepsilon, \delta}$.

Regarding $\mathcal{L}_{2,0}^{\varepsilon, \delta}$,

$$\left| \langle \mathcal{L}_{2,0}^{\varepsilon, \delta}, \phi \rangle \right| \leq \tilde{C} \left\| u_t^{\varepsilon, \delta} \right\|_{L^\infty(\mathbb{R}_+; L^1(\mathbb{R}^2))} \|\phi\|_{L^\infty(\Omega)}, \quad \phi \in C_0(\Omega),$$

which yields $\left\| \mathcal{L}_{2,0}^{\varepsilon, \delta} \right\|_{\mathcal{M}(\Omega)} \leq C$, where $\mathcal{M}(\Omega) = (C_c(\Omega))^*$ denotes the space of bounded measures on Ω .

In view of Lemma 4.3 and the uniform boundedness of the solutions, it follows that $\left\{ \mathcal{L}_{2,1}^{\varepsilon, \delta} \right\}_{\varepsilon, \delta > 0}$ is compact (and in fact converges to zero) in $W^{-1,2}(\Omega)$ and $\left\| \mathcal{L}_{2,2}^{\varepsilon, \delta} \right\|_{\mathcal{M}(\Omega)} \leq C$, for each fixed $t > 0$.

Next, for any $\phi \in C_c(\Omega)$, observe that

$$\begin{aligned} & \left| \iint_{\mathbb{R}^2} I_3^{\varepsilon, \delta} \phi \, dx \, dy \right| \\ & \leq C \left\{ \iint_{\mathbb{R}^2} |\varepsilon k_x^\delta| |k_x^\delta| \right\}^{\frac{1}{2}} \left\{ \iint_{\mathbb{R}^2} \varepsilon (u_x^{\varepsilon, \delta})^2 \, dx \, dy \right\}^{\frac{1}{2}} \\ & \quad + C \left\{ \iint_{\mathbb{R}^2} |\varepsilon k_y^\delta| |k_y^\delta| \right\}^{\frac{1}{2}} \left\{ \iint_{\mathbb{R}^2} \varepsilon (u_y^{\varepsilon, \delta})^2 \, dx \, dy \right\}^{\frac{1}{2}} \end{aligned}$$

The point here is to have ε and δ in balance, so that we can ensure $|\varepsilon k_x^\delta|, |\varepsilon k_y^\delta| \leq C$. More precisely, we have $|\varepsilon k_x^\delta|, |\varepsilon k_y^\delta| \leq C \frac{\varepsilon}{\delta}$, and by choosing ε, δ according to (4.3) we achieve this balance. Consequently,

$$\left| \iint_{\mathbb{R}^2} I_3^{\varepsilon, \delta} \phi \, dx \, dy \right| \leq C \|\phi\|_{L^\infty(\Omega)},$$

and thus $\left\| \mathcal{L}_{2,3}^{\varepsilon, \delta} \right\|_{\mathcal{M}(\Omega)} \leq C$, for each fixed $t > 0$.

Finally, using the *BV* regularity of the coefficients and the boundedness of the solutions,

$$\left| \iint_{\mathbb{R}^2} I_4^{\varepsilon, \delta} \phi \, dx \, dy \right| \leq C \|\phi\|_{L^\infty(\Omega)},$$

and thus $\left\| \mathcal{L}_{2,4}^{\varepsilon, \delta} \right\|_{\mathcal{M}(\Omega)} \leq C$, for each fixed $t > 0$.

Similarly, $\left| \iint_{\mathbb{R}^2} I_5^{\varepsilon, \delta} \phi \, dx \, dy \right| \leq C \|\phi\|_{L^\infty(\Omega)}$, and thus $\left\| \mathcal{L}_{2,5}^{\varepsilon, \delta} \right\|_{\mathcal{M}(\Omega)} \leq C$, for each fixed $t > 0$.

Summarizing, we have shown that the sequence of distributions $\left\{ \mathcal{L}^{\varepsilon, \delta} \right\}_{\varepsilon, \delta > 0}$ satisfies the following two properties: {i} each distribution is the sum of two terms — one is compact in $W^{-1,2}(\Omega)$ and the other

one is bounded in $\mathcal{M}(\Omega)$. In addition, Lemma 4.1 implies that $\{\text{ii}\} \{\mathcal{L}^{\varepsilon,\delta}\}_{\varepsilon,\delta>0}$ belongs to a bounded subset of $W^{-1,\infty}(\Omega)$. We now appeal to Murat lemma [28], which guarantees that $\{\mathcal{L}^{\varepsilon,\delta}\}_{\varepsilon,\delta>0}$ belongs to a compact subset of $W^{-1,2}(\Omega)$. This concludes the proof of the first part of the lemma, since Ω was an arbitrary bounded open subset of \mathbb{R}^2 . The second part of the lemma can be proved in a similar way. \square

Concluding the proof of Theorem 2.1. By combining Lemmas 3.2 and 4.4, we conclude that $u^{\varepsilon,\delta}(\cdot, \cdot, t)$ is pre-compact a.e. for each $t \in [0, T]$. Together with a diagonal argument, we can prove that $u^{\varepsilon,\delta}(\cdot, \cdot, t)$ converges along a subsequence a.e. in $\mathbb{R}^2 \times \mathbb{R}_+$ and in $L^1_{\text{loc}}(\mathbb{R}^2)$, for each fixed $t > 0$. Lemma 4.2 implies that

$$\|u^{\varepsilon,\delta}(\cdot, \cdot, t + \tau) - u^{\varepsilon,\delta}(\cdot, \cdot, t)\|_{L^1(\mathbb{R}^2)} \leq C\tau, \quad \forall \tau \in (0, T - \tau),$$

and using this L^1 time continuity estimate it takes a standard density argument to show that there exists a subsequence of $\{u^{\varepsilon,\delta}\}_{\varepsilon,\delta>0}$ that converges to a limit function u a.e. in $\mathbb{R}^2 \times \mathbb{R}_+$ and in $L^1_{\text{loc}}(\mathbb{R}^2 \times \mathbb{R}_+)$. Moreover, the limit u belongs to $L^\infty(\mathbb{R}^2 \times \mathbb{R}_+) \cap \text{Lip}(\mathbb{R}_+; L^1(\mathbb{R}))$.

Equipped with the strong convergence it is easy to prove that the limit u is a weak solution. Moreover, since each $u^{\varepsilon,\delta}$ satisfies an L^1 contraction principle, it follows that the same is true for limit. In particular, the possibly different limits of different subsequences of $u^{\varepsilon,\delta}$ coincide, and we conclude that the *whole* family of vanishing viscosity solutions converges $L^1_{\text{loc}}(\mathbb{R}^2 \times \mathbb{R}_+)$ -strong to u . This concludes the proof of Theorem 2.1. \square

We close the paper by a couple of remarks. First, we note that the apriori bounds in Lemmas 4.1, 4.2 and 4.3, being uniform in time, enabled us to deduce pre-compactness at each fix $t > 0$, thus circumventing the temporal argument required in [36, Appendix A]. Second, we have herein exclusively dealt with problems that are spatially two-dimensional. A possible strategy for going beyond two dimensions might be to adapt the compactness framework of Panov [30, 32, 31].

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