

# A traffic flow model for the dynamics of traffic jams: a Pressureless Gas Dynamics system under a maximal constraint

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Joint work with

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1. Traffic models: overview on fluid models
2. Modified Aw-Rascle model
3. Limit  $\varepsilon \rightarrow 0$ : Constrained Pressureless Gas Dynamics
4. CPGD: additional laws
5. Existence theorem for CPGD
6. Numerical simulations
7. Conclusion

# 1. Traffic models: overview on fluid models

⇒ Conservation of car density

$$\partial_t n + \partial_x q = 0$$

⇒ What expression for the flux  $q$  ?

- Conservation of car density

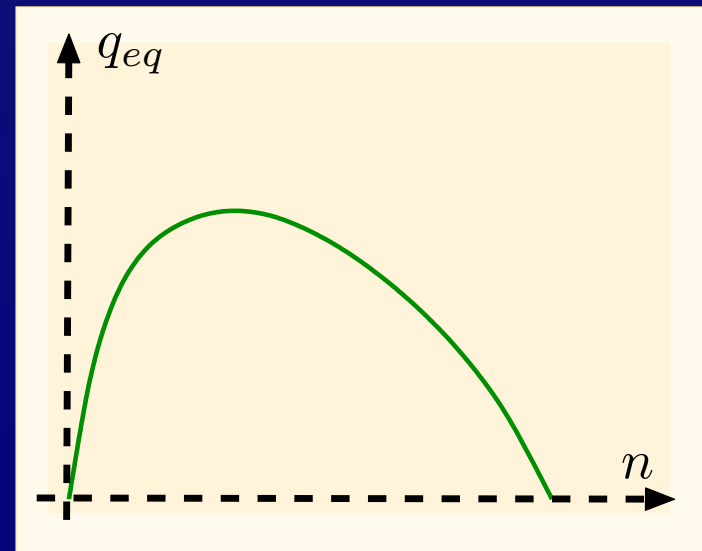
$$\partial_t n + \partial_x q = 0$$

- What expression for the flux  $q$  ?

- First order models:

$$q = q_{eq}(n)$$

[Lighthill, Witham (1955)], ...



- ⇒ Second order models:  $q = nu$  and gas dynamics-like eq. for  $u$ :

$$\partial_t nu + \partial_x(nu^2 + p) = -\frac{nu - q_{eq}(n)}{\tau}$$

- ⇒ [Payne (1971)], ...

- ➡ Second order models:  $q = nu$  and gas dynamics-like eq. for  $u$ :

$$\partial_t nu + \partial_x(nu^2 + p) = -\frac{nu - q_{eq}(n)}{\tau}$$

- ➡ [Payne (1971)], ...
- ➡ [Daganzo (1995)]: Inacceptable properties (e.g. Vehicles going backwards)
  - ➡ Fluid  $\Rightarrow$  sound propagation is isotropic in a comoving frame
  - ➡ Traffic: information propagates backwards

- ➡ Modified 2nd order model (see also [Zhang (2002)])
- ➡ Preferred velocity  $w$  is a Lagrangian quantity:

$$\dot{w} := (\partial_t + u\partial_x)w = 0$$



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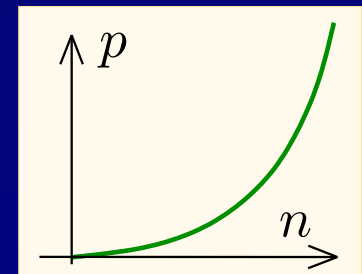
⇒ Preferred velocity  $w$  is a Lagrangian quantity:

$$\dot{w} := (\partial_t + u\partial_x)w = 0$$

⇒ The actual velocity  $u$  offsets the preferred velocity  $w$  by a quantity  $p(n)$  which increases with  $n$

$$w = u + p(n), \quad p \nearrow \text{ as } n \nearrow$$

⇒ Typically  $p(n) = n^\gamma, \gamma > 0$



$$\partial_t n + \partial_x (nu) = 0$$

$$(\partial_t + u\partial_x)(u + p(n)) = 0$$

⇒ Second eq. equivalent to

$$(\partial_t + (u - np'(n))\partial_x)u = 0$$

⇒ Two characteristic velocities:

⇒  $\lambda_1 = u - np'(n)$  (assoc.w.  $u$ , GNL)

⇒  $\lambda_2 = u$  (assoc.w.  $w = u + p(n)$ , LD)

⇒ Invariant regions:  $(u, w)$  - rectangles

⇒ If

$$a < u_0 < b \quad \text{and} \quad c < w_0 < d$$

then for all times

$$a < u(t) < b \quad \text{and} \quad c < w(t) < d$$

⇒ Prevents  $u < 0$  (no vehicle going backwards !)

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⇒ AR model in Lagrangian coordinates = continuous version of Follow-the-Leader model [Aw, Klar, Materne, Rascle (2002)]

# No invariant region for $n$ !

- ➡ Problem: there is no invariant region for  $n$ 
  - ➡  $n > 0$  BUT:
  - ➡  $n$  can exceed the upper limit  $n^*$  (if any) even if initially  $n < n^*$ )

- Problem: there is no invariant region for  $n$ 
  - $n > 0$  BUT:
  - $n$  can exceed the upper limit  $n^*$  (if any) even if initially  $n < n^*$ )
- Goal:
  - Modify AR model s.t. it enforces an upper limit constraint on the density  $n \leq n^*$
  - Investigate the dynamics of the clustered regions where  $n = n^*$

## 2. Modified Aw-Rascle model

⇒ AR model which guarantees the constraint

$$n < n^*$$

at all times



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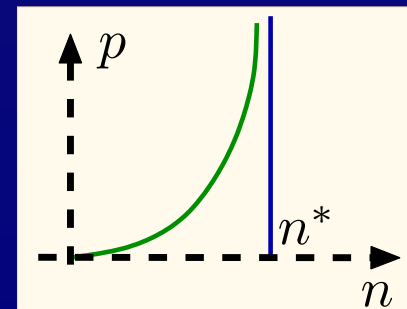
at all times

- ➡ Modify  $p(n)$  s.t.

$$p(n) \longrightarrow \infty \quad \text{as} \quad n \longrightarrow n^*$$

- ➡ For instance

$$p(n) = \frac{1}{\left(\frac{1}{n} - \frac{1}{n^*}\right)^\gamma}$$



- M-AR has the same properties as the standard AR model
  - Hyperbolicity
  - Invariant regions
- Satisfies the density constraint

$$n < n^*$$

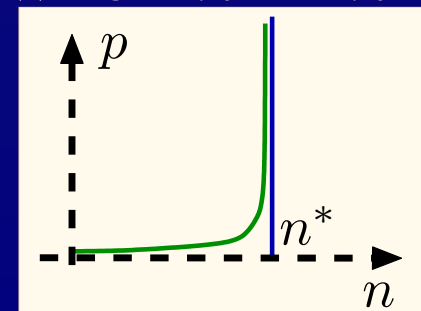
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- Modeled by the rescaling:

$$p(n) = \varepsilon \tilde{p}(n)$$



▶▶▶ Perturbed AR system

$$\partial_t n^\varepsilon + \partial_x (n^\varepsilon u^\varepsilon) = 0$$

$$(\partial_t + u^\varepsilon \partial_x)(u^\varepsilon + \varepsilon p(n^\varepsilon)) = 0$$

▶▶▶ with modified velocity offset:

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▶▶▶ Question: what happens in the limit

$$\varepsilon \longrightarrow 0$$

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### 3. Limit $\varepsilon \rightarrow 0$ : Constrained Pressureless Gas Dynamics



- ➡ Suppose  $n^\varepsilon \rightarrow n < n^*$  (uncongested case)
- ➡ Then  $\varepsilon p(n^\varepsilon) \rightarrow 0$  in (RM-AR) model:

$$\partial_t n^\varepsilon + \partial_x (n^\varepsilon u^\varepsilon) = 0$$

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⇒ Limit system = Pressureless Gas Dynamics

$$\partial_t n + \partial_x (nu) = 0$$

$$(\partial_t + u \partial_x)u = 0$$

⇒ ⇒ Mass conservation

⇒ Burger's eq. for the velocity

⇒ Not strictly hyperbolic

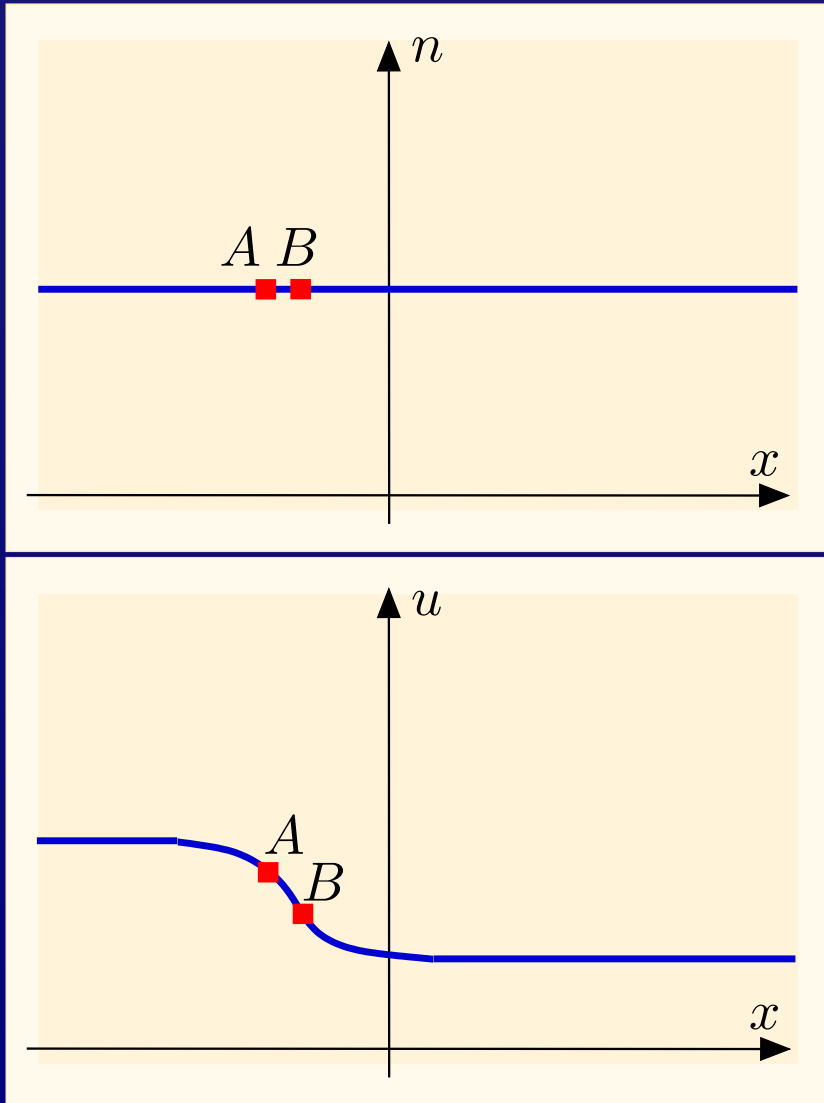
⇒ 2 identical eigenvalues  $u$

⇒ But not diagonalizable: Jacobian = 
$$\begin{pmatrix} u & n \\ 0 & u \end{pmatrix}$$

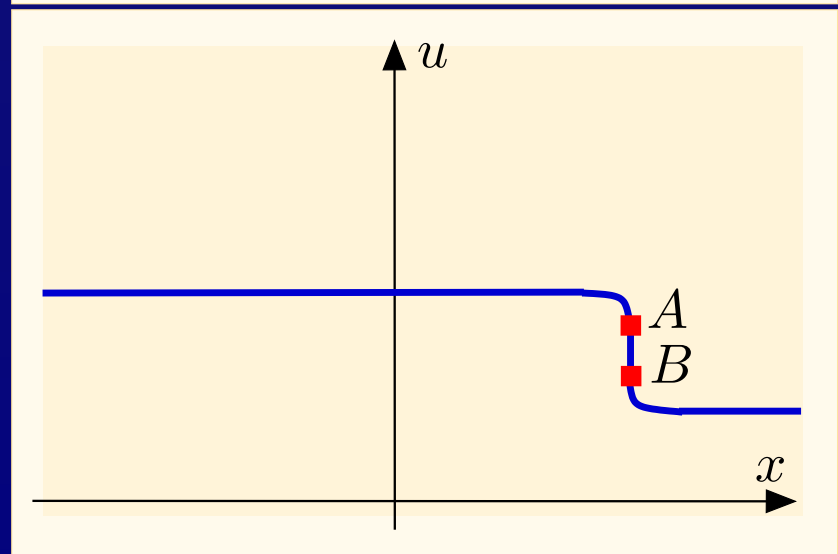
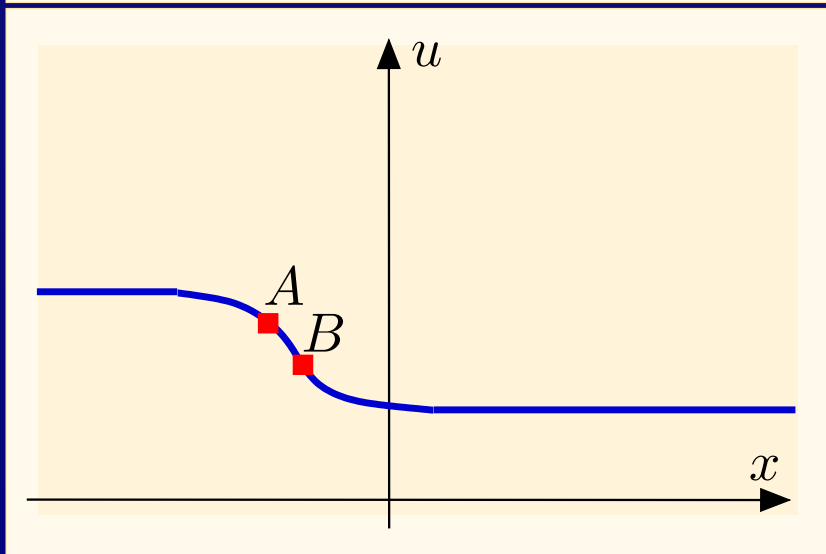
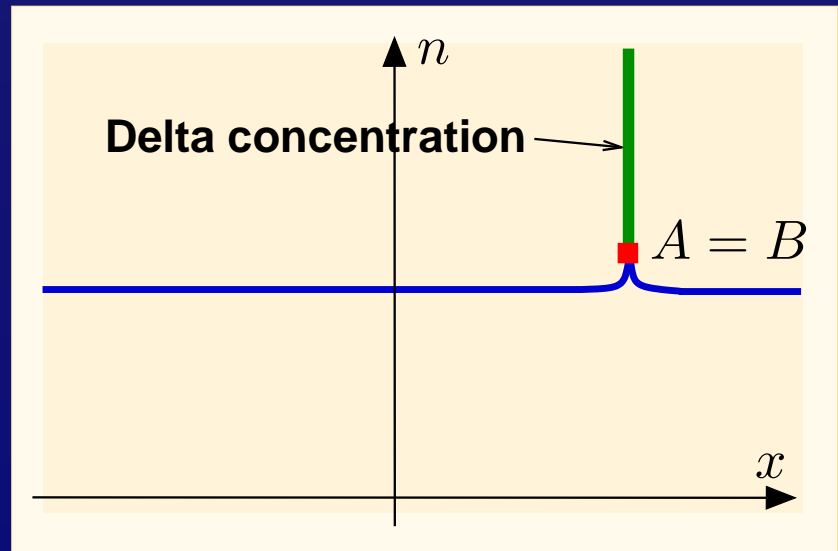
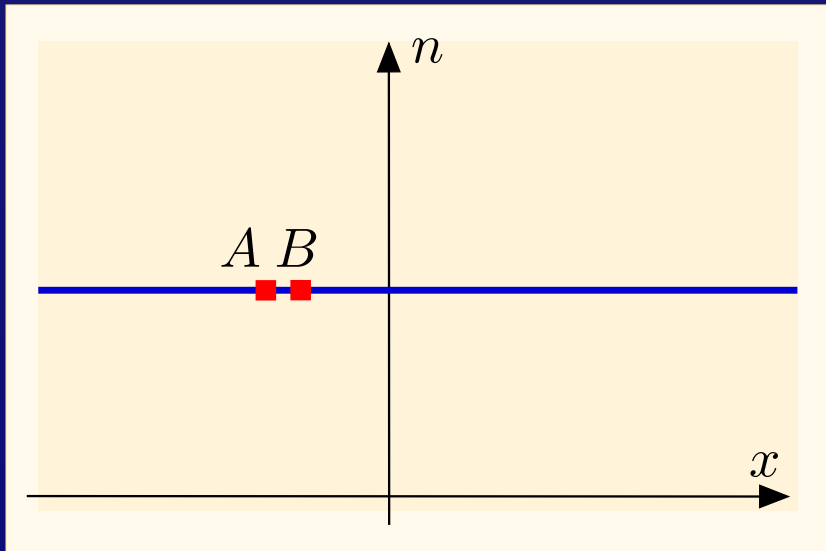
⇒ Weak instability:

⇒ linearized solution increase like  $O(t)$

⇒ Generates mass concentrations

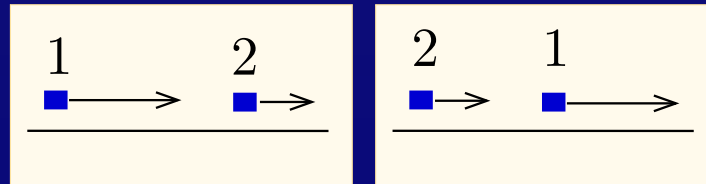


# Pressureless Gas Dyn. concentrations 18

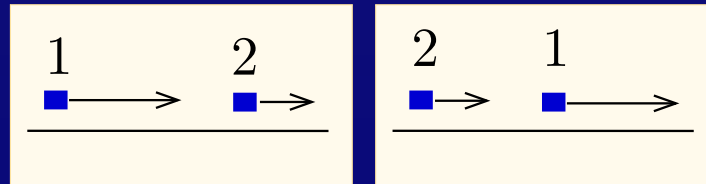


- ➡ Concentrations = 'particles'
- ➡ Beyond concentration: solution not unique  
Depends on particle interaction model

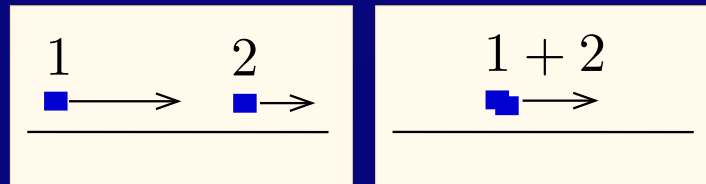
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- Sticky particles (Zeldowitch, E, ...)



- see e.g. [Bouchut (94)], [Grenier (95)], [Rykov, Sinai (96)], [Brenier, Grenier (98)], ...



- Density constraint: no concentration formation
  - No need to define a particle dynamics
- Instead: formation of 'clusters' (traffic jams)
  - Cluster dynamics follows from the asymptotic limit

⇒ Suppose  $n^\varepsilon \rightarrow n^*$  (then  $p(n^\varepsilon) \rightarrow \infty$ )

⇒ Suppose  $\varepsilon p(n^\varepsilon) \rightarrow \bar{p} < \infty$

⇒ Then  $\varepsilon \rightarrow 0$  in (RM-AR) model:

$$\partial_t n^\varepsilon + \partial_x (n^\varepsilon u^\varepsilon) = 0$$

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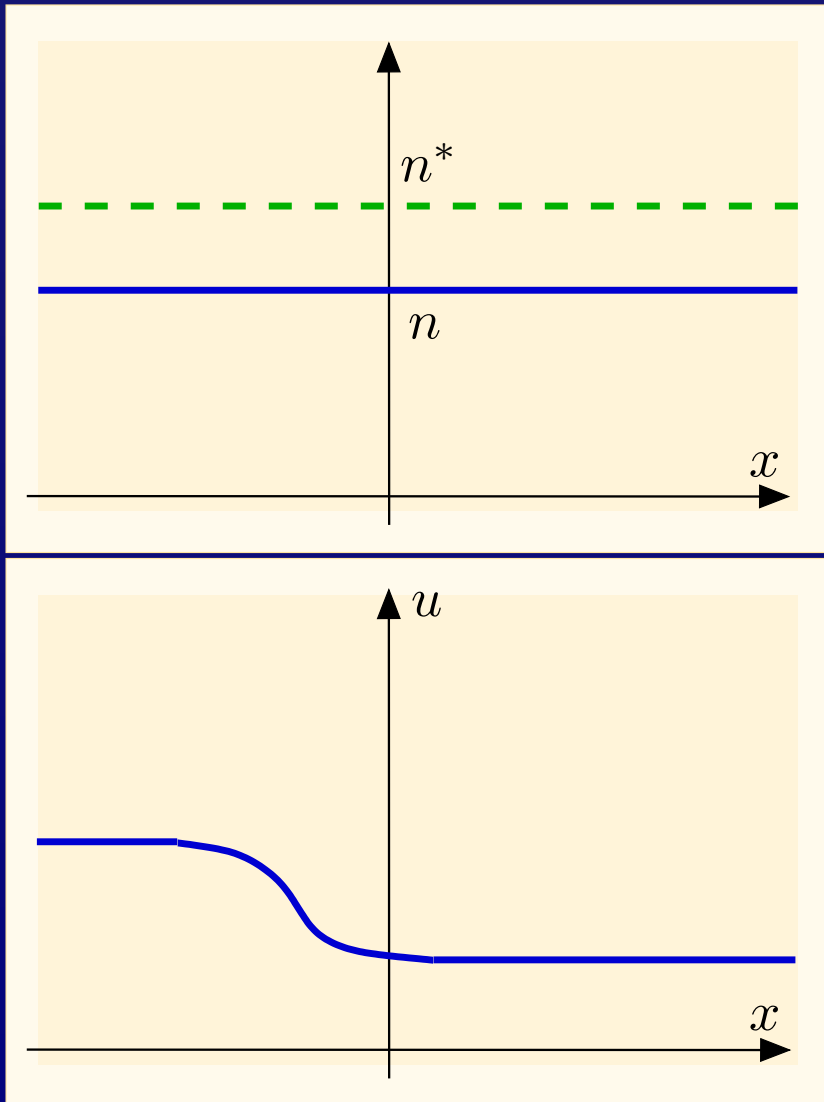
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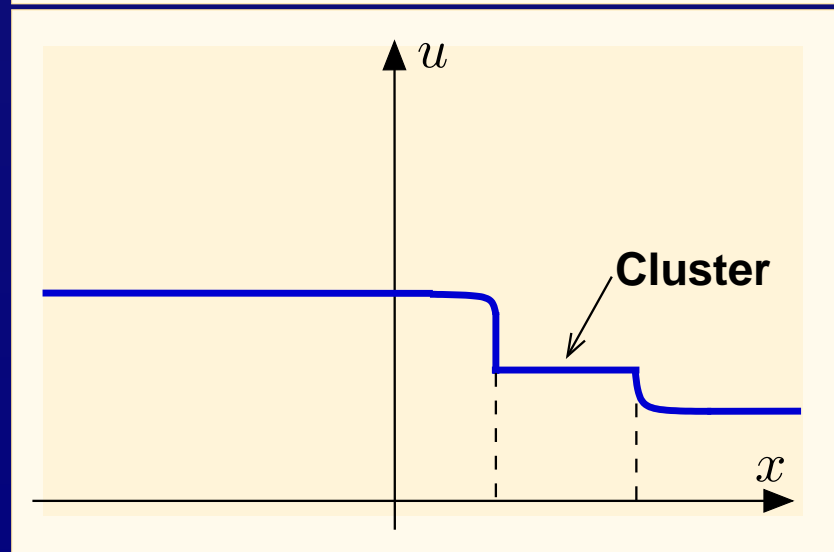
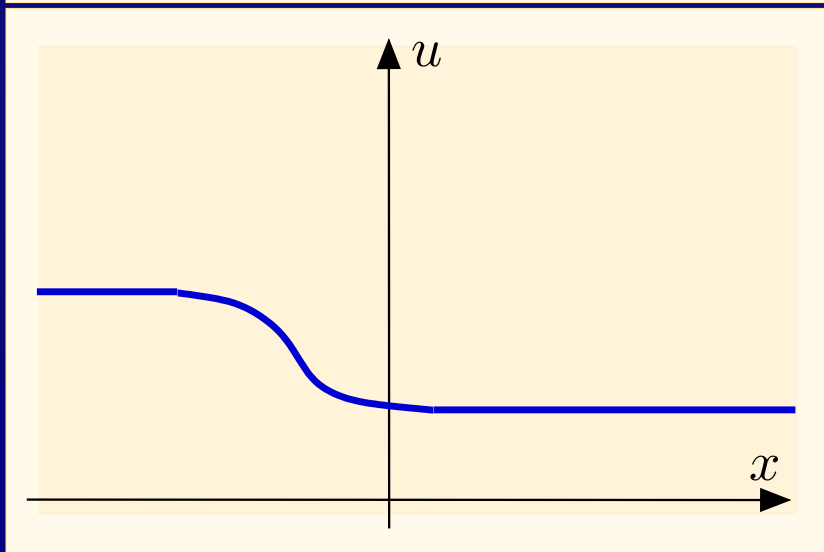
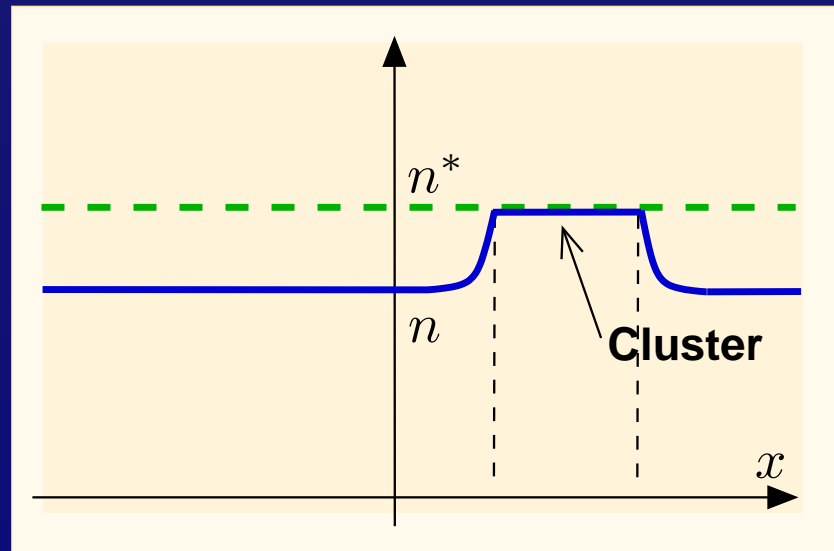
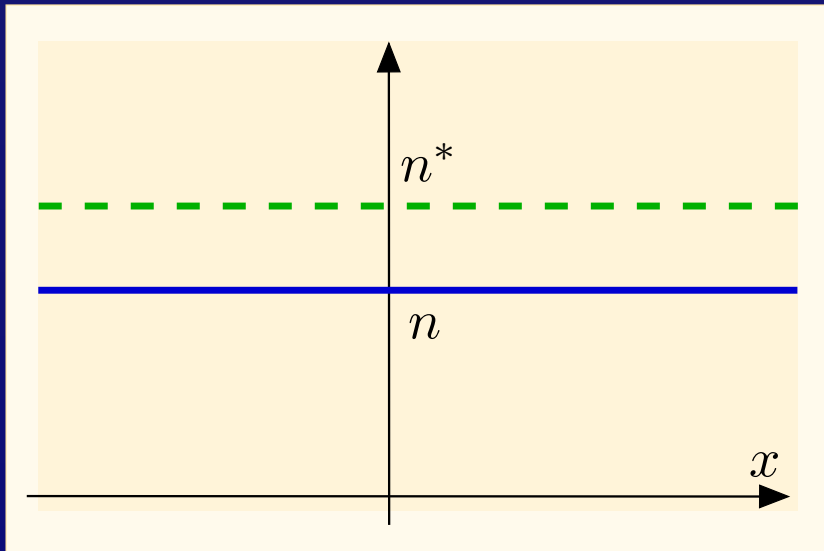
$$\partial_t n + \partial_x (nu) = 0$$

$$(\partial_t + u \partial_x)(u + \bar{p}) = 0$$

$$n = n^*$$

⇒  $\bar{p}$  unknown: Lagrange multiplier





## ➡ Constrained Pressureless Gas Dynamics (CPGD)

$$\partial_t n + \partial_x(nu) = 0$$

$$(\partial_t + u\partial_x)(u + \bar{p}) = 0$$

$$\bar{p}(n^* - n) = 0$$

$$\bar{p} \geq 0, \quad 0 \leq n \leq n^*$$

➡ see e.g. [Brenier, ...], [B. and Bouchut] for gaseous corks in pipes

## 4. Constrained Pressureless Gas Dynamics: additional laws

- ➡ CPGD formulation ill-posed  
lack of information for defining a unique solution



- ⇒ CPGD formulation ill-posed  
lack of information for defining a unique solution
- ⇒ To be defined
  - ⇒ Cluster dynamics
  - ⇒ Value of  $\bar{p}$  inside clusters
  - ⇒ What if clusters meet ?

⇒ If  $n^\varepsilon \rightarrow n^*$  with  $\varepsilon p(n^\varepsilon) \rightarrow \bar{p} < \infty$ , then:

$$\varepsilon n^\varepsilon p'(n^\varepsilon) \rightarrow \infty$$

⇒ Characteristic velocities:

$$\Rightarrow \lambda_1^\varepsilon = u^\varepsilon - n^\varepsilon p'(n^\varepsilon) \rightarrow -\infty$$

$$\Rightarrow \lambda_2^\varepsilon = u^\varepsilon \rightarrow u$$

⇒ Riemann invariant associated with  $\lambda_1$  is  $u$ :

$$\partial_t u^\varepsilon + \lambda_1^\varepsilon \partial_x u^\varepsilon = 0$$

⇒ In the limit  $\varepsilon \rightarrow 0$ ,  $u$  is constant in a cluster

$$\partial_x u = 0$$

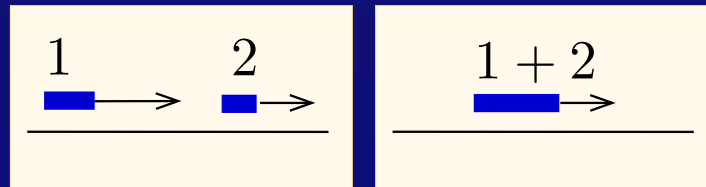
- ➡ In a cluster, all cars have the same velocity (the cluster velocity)
- ➡ Since  $\lambda_1^\varepsilon \rightarrow -\infty$ , all variations of the cluster velocity originate from the leading vehicle of the cluster
- ➡ Any variation of the velocity of the leading vehicle instantaneously propagates to the whole cluster

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  - ⇒ Gives no information about cluster dynamics beyond what has been noticed above

- ▶▶▶ Limit (RM-AR)  $\rightarrow$  (CPGD) is formal
  - ▶▶ Gives no information about cluster dynamics beyond what has been noticed above
  
- ▶▶▶ But Riemann problem solutions of (RM-AR) are explicit
  - ▶▶ Limit  $\varepsilon \rightarrow 0$  in these solutions give information about cluster dynamics

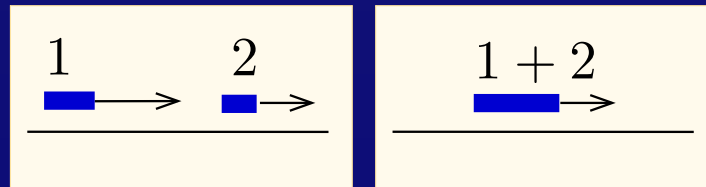
# Cluster dynamics (from Riemann pbm) 29

- When two clusters meet, they merge
  - The resulting cluster takes the velocity of the front cluster (the slowest one)



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- At the right-most end of a cluster, there is
  - either vacuum (and  $\bar{p} = 0$  at the right-most point of the cluster)
  - or the velocity is continuous across the right most point of the cluster

- These rules allow to determine
  - $\bar{p}$  everywhere
  - The velocity of the extreme points of the cluster

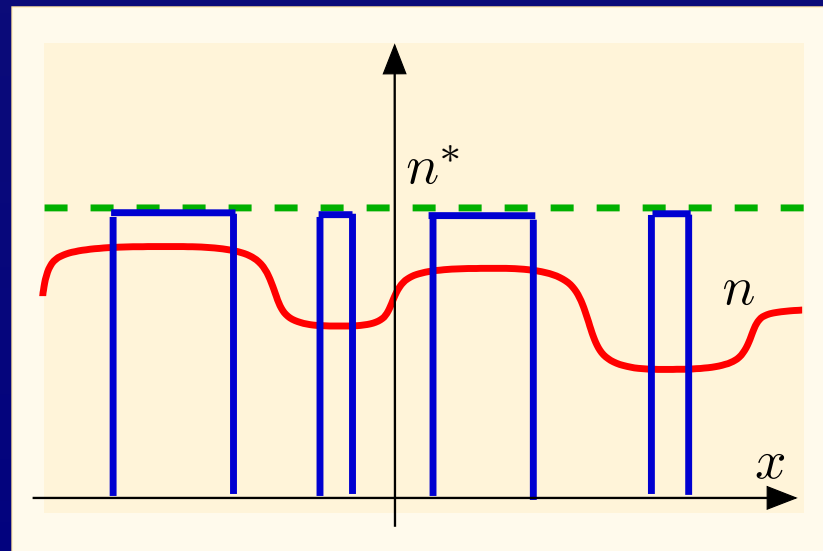


- These rules allow to determine
  - $\bar{p}$  everywhere
  - The velocity of the extreme points of the cluster
- Cluster dynamics  $\neq$  from that of [Brenier, ...] and [B. and Bouchut]
  - no momentum conservation at cluster collapse

## 5. Existence theorem for CPGD

- Idea (follows from [B. and Bouchut (2002, 2003)]),
  - ➔ Approximate (in  $\mathcal{D}'$ ) the solution by clusters

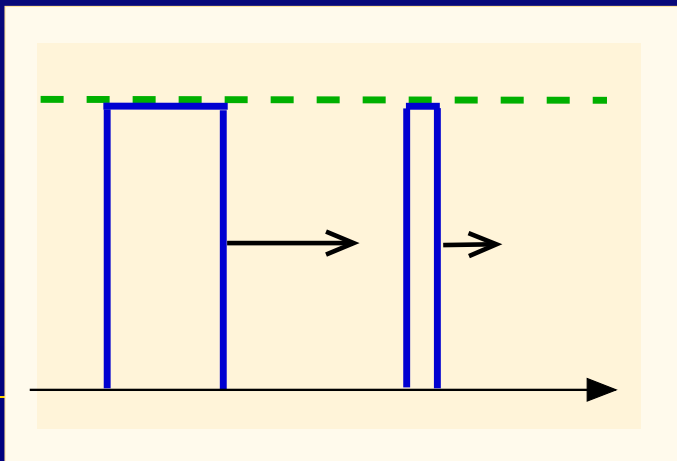
$$\begin{pmatrix} n(x, t) \\ (nu)(x, t) \end{pmatrix} \approx \sum_1^N \begin{pmatrix} n^* \\ n^* u_i(t) \end{pmatrix} \chi_{a_i(t) \leq x \leq b_i(t)}$$



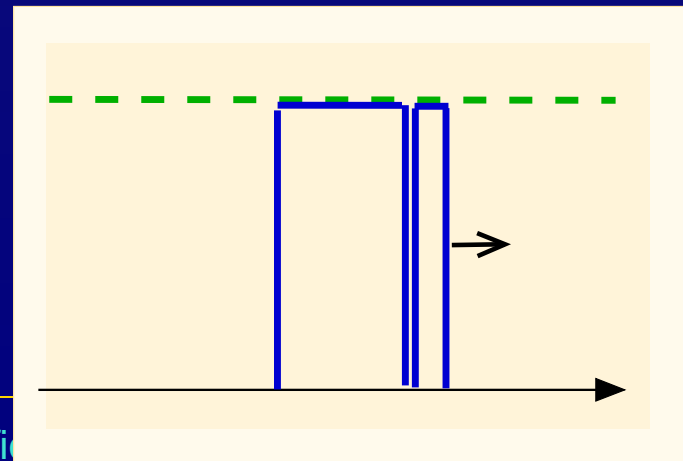
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(Summary)



(Conclusion)

- Is a weak solution of (CPGD)
- Satisfies an 'Oleinik' condition

$$\partial_x u \leq \frac{1}{t}$$

- Satisfies  $L^\infty$  and BV bounds:

$$\inf u_0 \leq u(., t) \leq \sup u_0$$

$$TV_K(\bar{p}) \leq 2TV(u_0), \quad \forall K \text{ compact}$$

$$0 \leq \bar{p} \leq 2 \sup u_0$$

⇒ Satisfies an 'entropy' equality

⇒ For any  $S \in C^1$ ,  $\exists \bar{p}^S$  s.t.

$$\partial_t(nS(u) + n\bar{p}^S) + \partial_x(nuS(u) + nu\bar{p}^S) = 0$$

⇒  $\bar{p}^S$  has a BV bound:

$$TV_K(\bar{p}^S) \leq 2|S'|_\infty TV u_0$$



- ⇒ Step 1: approximate initial condition  $(n_0, u_0)$  by a converging sequence of clusters  $(n_0^k, u_0^k)$  [B. (02)]
- ⇒ Defines a sequence of cluster sol.  $(n^k, u^k)$  satisfying the above a priori bounds

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- ➡ Key point: an extension of the compensated compactness lemma using Oleinik estimate [B. (02)]

⇒ Suppose

$$\Rightarrow n_0 \in L^1, \quad 0 \leq n_0 \leq n^*$$

$$\Rightarrow u_0 \in L^\infty \cap BV$$

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⇒  $u_0 \in L^\infty \cap BV$

⇒  $\exists n \in L_t^\infty(L_x^\infty \cap L_x^1)$  ,  $u, \bar{p} \in L_{x,t}^\infty$

⇒ a solution of CPGD

⇒ satisfying Oleinik estimate

⇒ satisfying entropy equality for any  $S \in C^1$

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⇒ No uniqueness so far

⇒ No convergence proof (RM-AR)  $\rightarrow$  (CPGD)

## 6. Numerical simulations



## Rescaled Modified Follow-the-Leader

$$\dot{x}_i = v_i, \quad \dot{v}_i = \frac{\varepsilon}{\gamma} \frac{v_{i+1} - v_i}{(x_{i+1} - x_i - d)^{\gamma+1}}$$

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## Limit $\varepsilon \rightarrow 0$

$$\dot{x}_i = v_i, \quad \dot{v}_i = \begin{cases} \dot{v}_i = 0 & \text{if } x_{i+1} - x_i > d \\ v_i = v_{i+1} & \text{if } x_{i+1} = x_i + d \end{cases}$$

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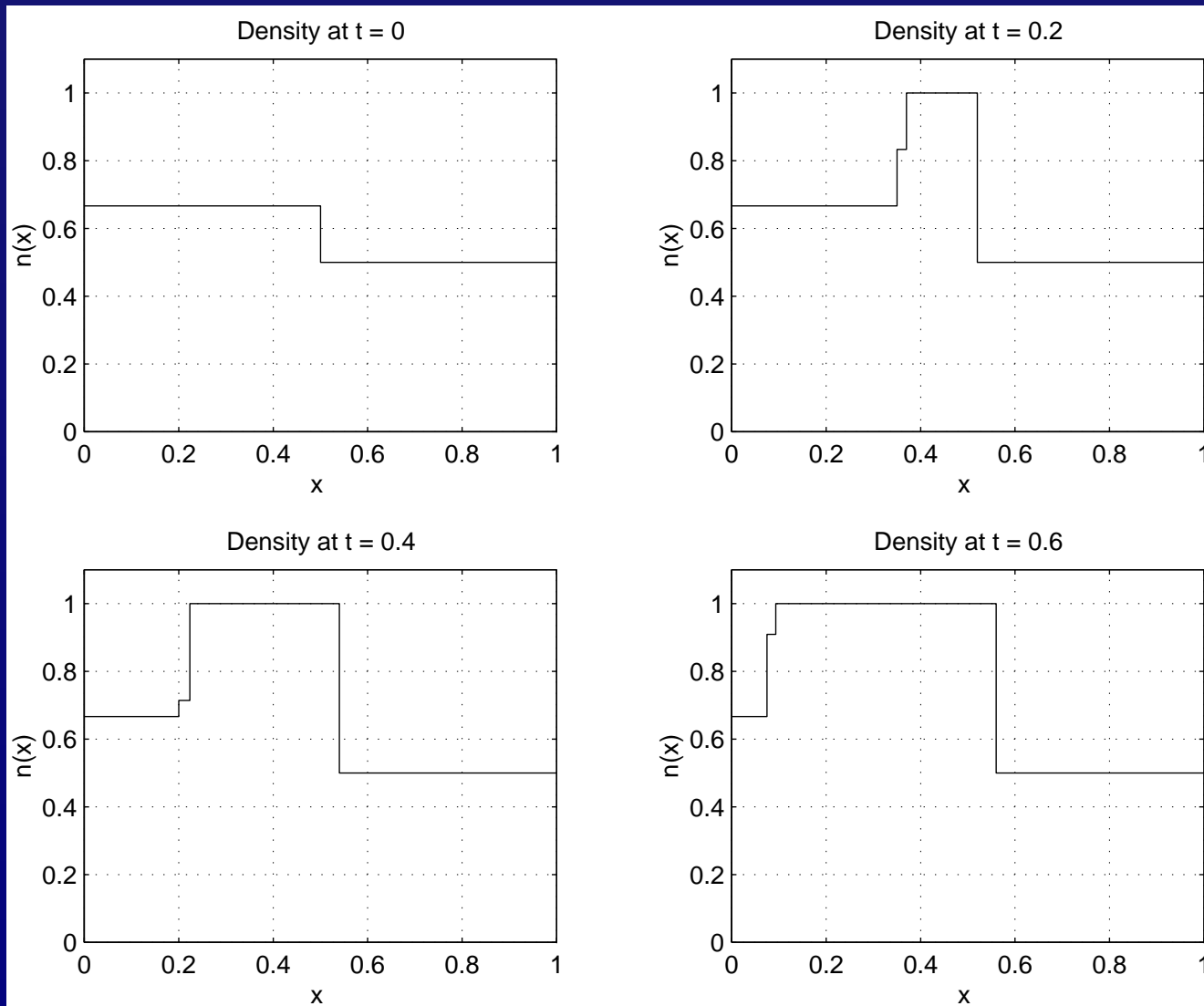
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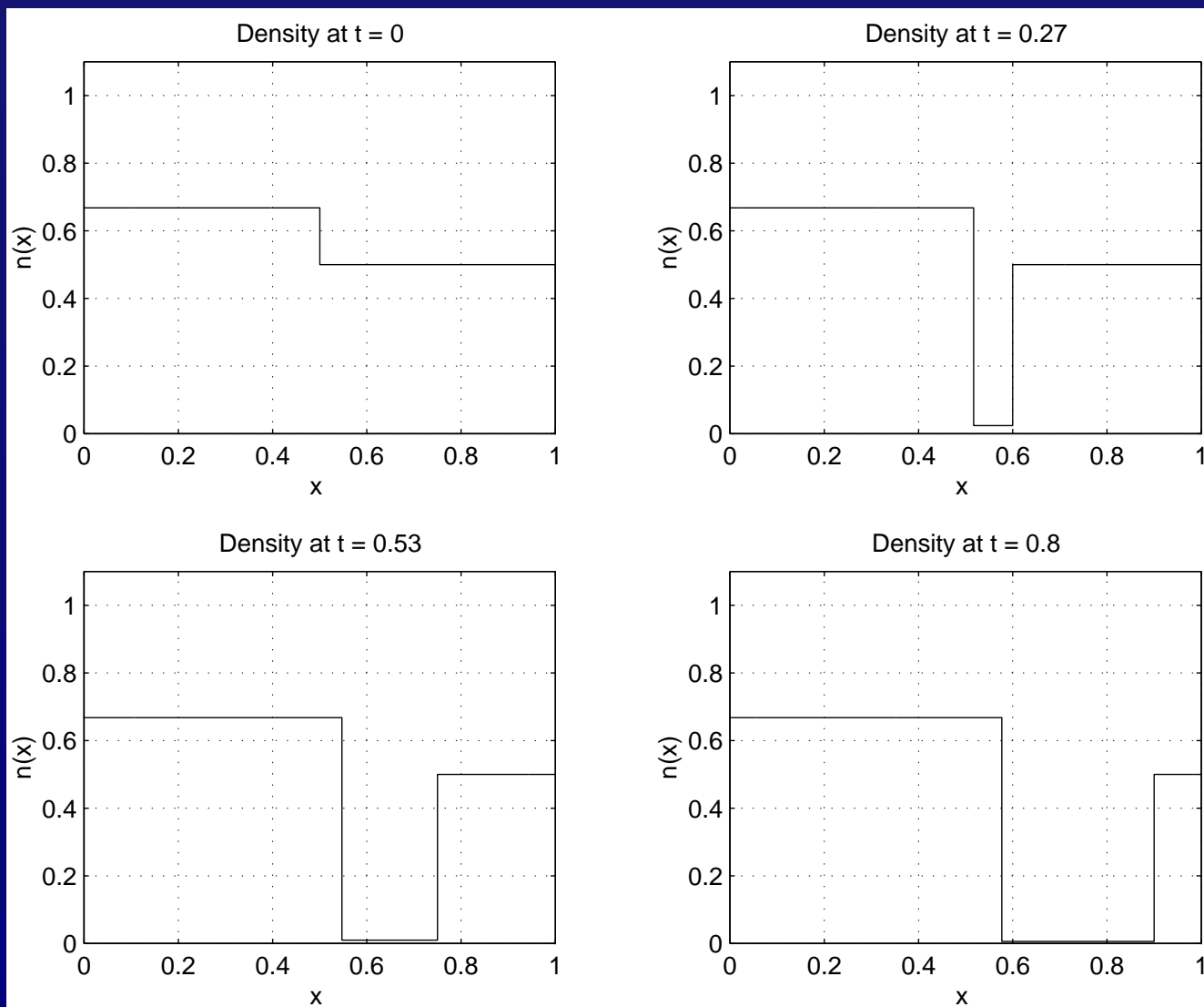
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## Cluster = sequence of vehicles separated by $d$

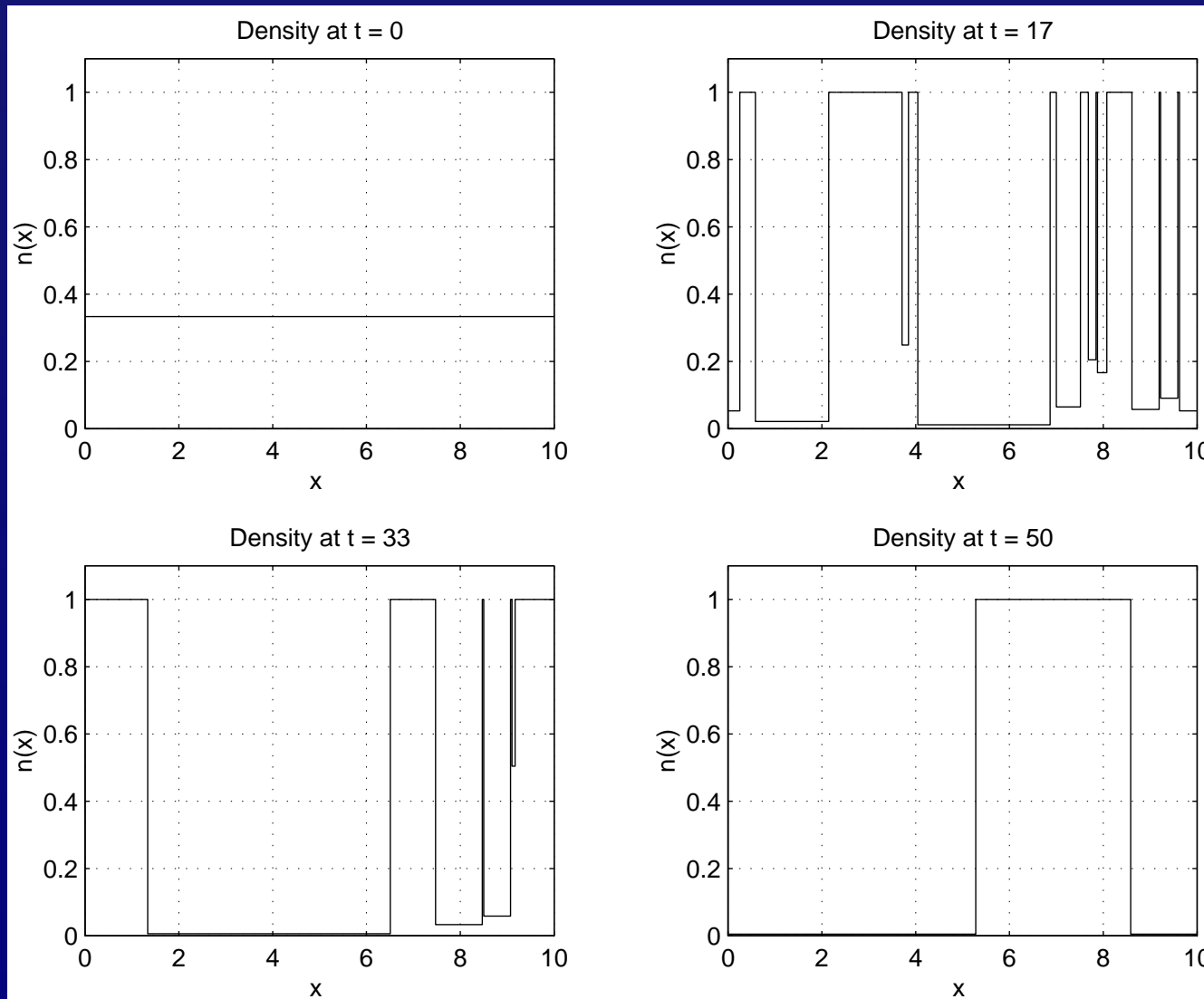
# Riemann problem: Cluster formation 40



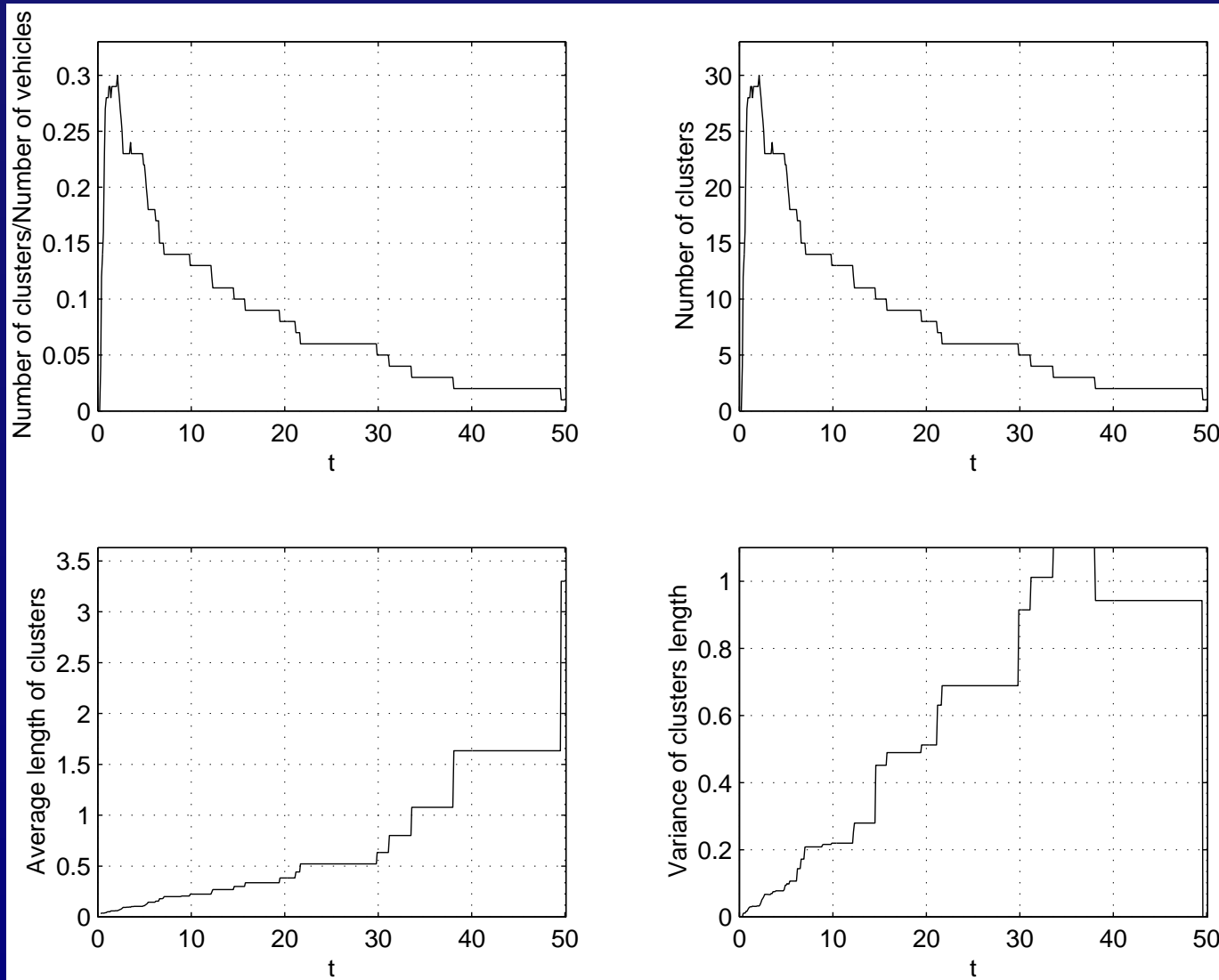
# Riemann problem: vacuum formation 41



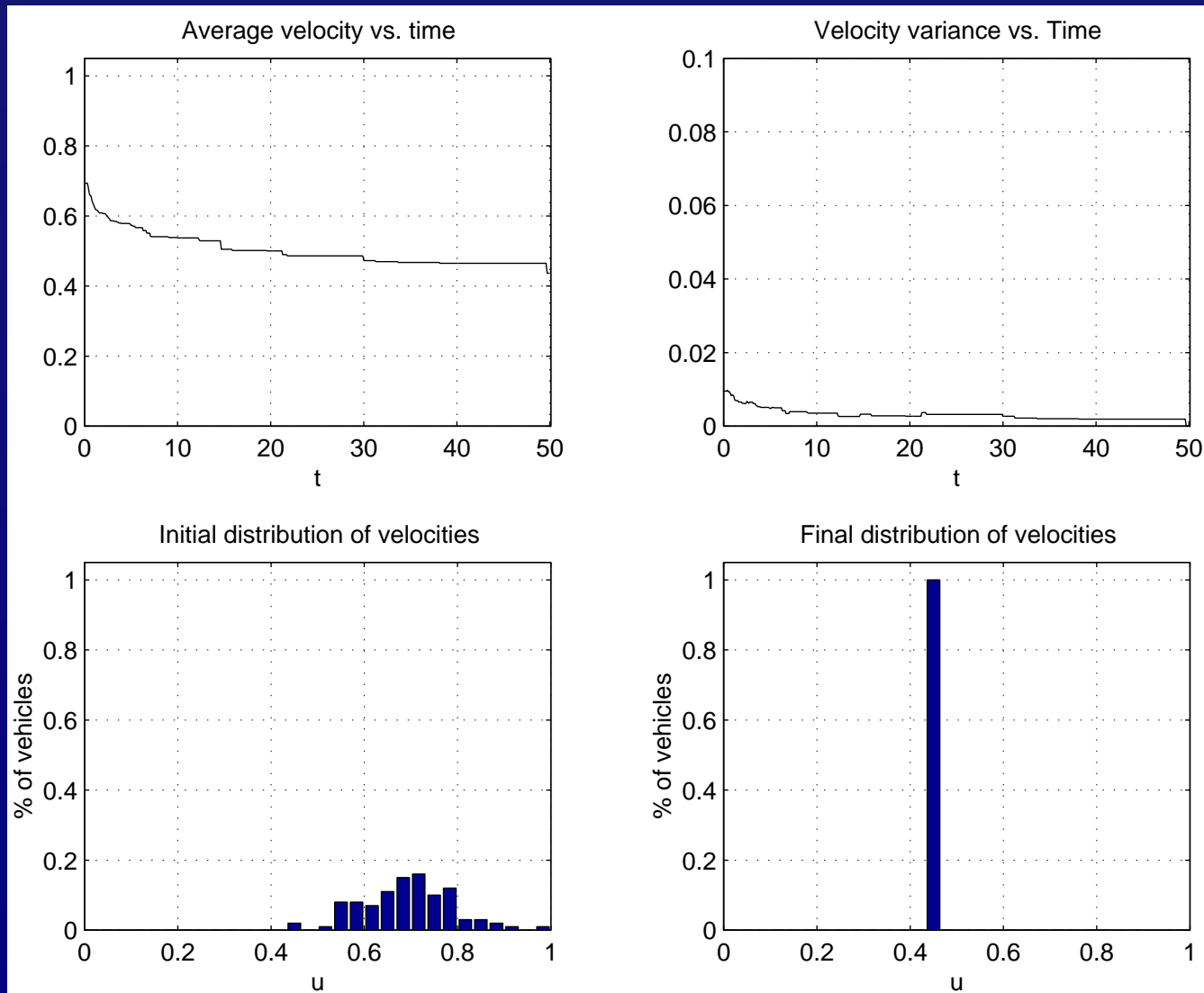
# Random initial velocity: cluster formation<sub>42</sub>



# Random initial velocity: cluster statistics 43

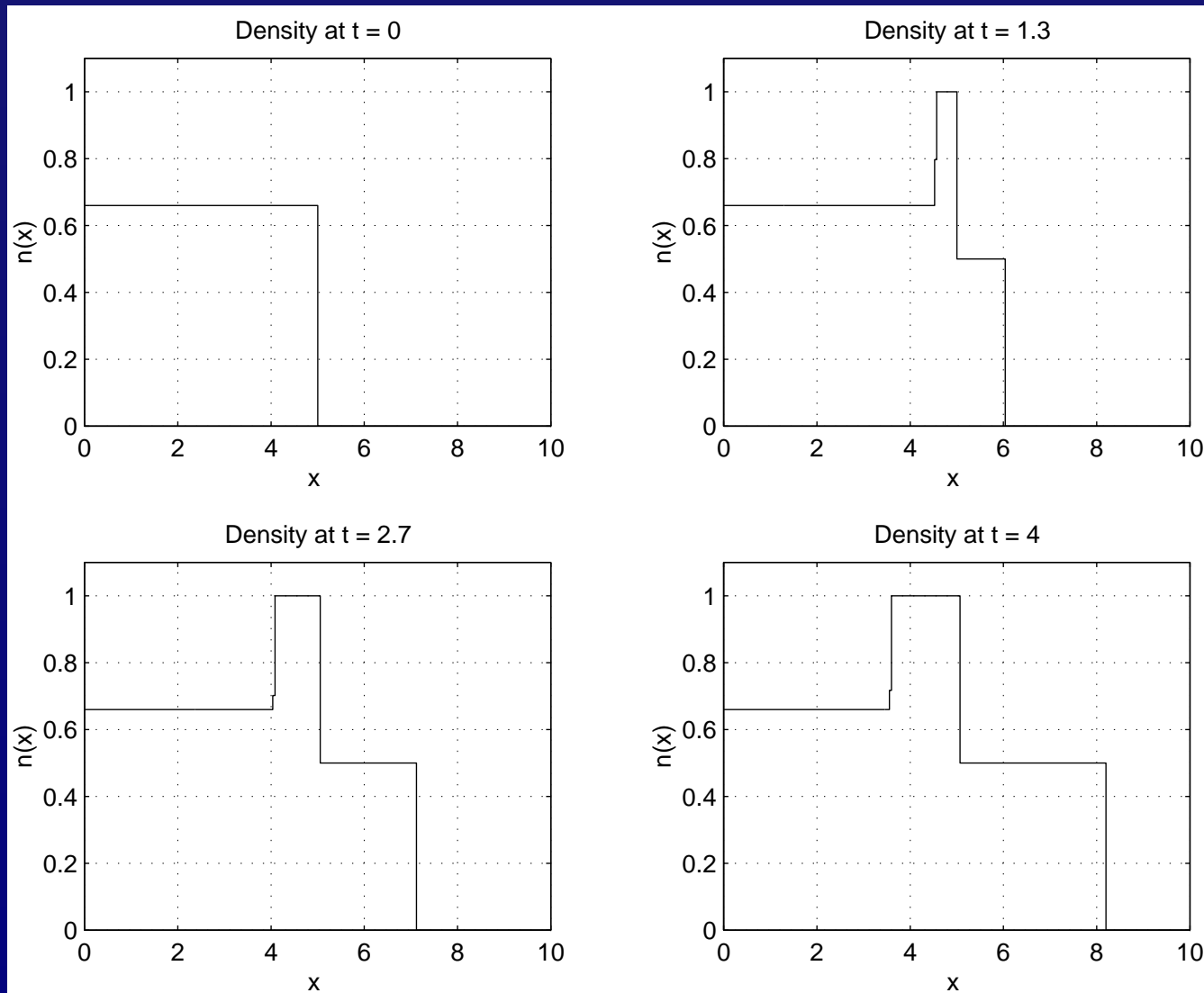


# Random initial velocity: velocity distribution





# Bottleneck: jam formation



## 7. Conclusion

- Modified Aw-Rascle model
  - Density constraint
  - Rescaled for small difference between preferred velocity and actual velocity in uncongested situations

- Modified Aw-Rascle model
  - Density constraint
  - Rescaled for small difference between preferred velocity and actual velocity in uncongested situations
  
- Limit model
  - Constrained Pressureless Gas Dynamics
  - Describes well cluster formation and dynamics
  - Existence theorem
  - Numerical simulations

## ⇒ CPGD:

⇒ Convergence proof (RM-AR)  $\rightarrow$  (CPGD)

⇒ Lagrangian formulation and scheme

- ⇒ CPGD:
  - ⇒ Convergence proof (RM-AR)  $\rightarrow$  (CPGD)
  - ⇒ Lagrangian formulation and scheme
- ⇒ More elaborate model
  - ⇒ Density constraint depends on velocity (work in progress with P. degond, V. Leblanc, M. Rascle and J. Royer)
  - ⇒ Multi-lane
  - ⇒ Multi-class
  - ⇒ etc.