

# CONFORMAL HARMONIC FORMS, BRANSON-GOVER OPERATORS AND DIRICHLET PROBLEM AT INFINITY.

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ABSTRACT. For odd dimensional Poincaré-Einstein manifolds  $(X^{n+1}, g)$ , we study the set of harmonic  $k$ -forms (for  $k < \frac{n}{2}$ ) which are  $C^m$  (with  $m \in \mathbb{N}$ ) on the conformal compactification  $\bar{X}$  of  $X$ . This is infinite dimensional for small  $m$  but it becomes finite dimensional if  $m$  is large enough, and in one-to-one correspondence with the direct sum of the relative cohomology  $H^k(\bar{X}, \partial\bar{X})$  and the kernel of the Branson-Gover [3] differential operators  $(L_k, G_k)$  on the conformal infinity  $(\partial\bar{X}, [h_0])$ . In a second time we relate the set of  $C^{n-2k+1}(\Lambda^k(\bar{X}))$  forms in the kernel of  $d + \delta_g$  to the conformal harmonics on the boundary in the sense of [3], providing some sort of long exact sequence adapted to this setting. This study also provides another construction of Branson-Gover differential operators, including a parallel construction of the generalization of  $Q$  curvature for forms.

## 1. INTRODUCTION

Let  $(M, [h_0])$  be an  $n$ -dimensional compact manifold equipped with a conformal class  $[h_0]$ . The  $k$ -th cohomology group  $H^k(M)$  can be identified with  $\ker(d + \delta_h)$  for any  $h \in [h_0]$  by usual Hodge-De Rham Theory. However, the choice of harmonic representatives in  $H^k(M)$  is not conformally invariant with respect to  $[h_0]$ , except when  $n$  is even and  $k = \frac{n}{2}$ . Recently, Branson and Gover [3] defined new complexes, new conformally invariant spaces of forms and new operators to somehow generalize this  $k = \frac{n}{2}$  case. More precisely, they introduce conformally covariant differential operators  $L_k^{\text{BG}, \ell}$  of order  $2\ell$  on the bundle  $\Lambda^k(M)$  of  $k$ -forms, for  $\ell \in \mathbb{N}$  (resp.  $\ell \in \{1, \dots, \frac{n}{2}\}$ ) if  $n$  is odd (resp.  $n$  is even). A particularly interesting case is the critical one in even dimension, this is

$$(1.1) \quad L_k^{\text{BG}} := L_k^{\text{BG}, \frac{n}{2} - k}.$$

The main features of this operator are that it factorizes under the form  $L_k^{\text{BG}} = G_{k+1}^{\text{BG}} d$  for some operator

$$(1.2) \quad G_{k+1}^{\text{BG}} : C^\infty(M, \Lambda^{k+1}(M)) \rightarrow C^\infty(M, \Lambda^k(M))$$

and that  $G_k^{\text{BG}}$  factorizes under the form  $G_k^{\text{BG}} = \delta_{h_0} Q_k^{\text{BG}}$  for some differential operator

$$(1.3) \quad Q_k^{\text{BG}} : C^\infty(M, \Lambda^k(M)) \cap \ker d \rightarrow C^\infty(M, \Lambda^k(M))$$

where  $\delta_{h_0}$  is the adjoint of  $d$  with respect to  $h_0$ . This gives rise to an elliptic complex

$$\dots \xrightarrow{d} \Lambda^{k-1}(M) \xrightarrow{d} \Lambda^k(M) \xrightarrow{L_k^{\text{BG}}} \Lambda^k(M) \xrightarrow{\delta_{h_0}} \Lambda^{k-1}(M) \xrightarrow{\delta_{h_0}} \dots$$

named the *detour complex*, whose cohomology is conformally invariant. Moreover, the pairs  $(L_k^{\text{BG}}, G_k^{\text{BG}})$  and  $(d, G_k^{\text{BG}})$  on  $\Lambda^k(M) \oplus \Lambda^k(M)$  are graded injectively elliptic in the sense that  $\delta_{h_0} d + dG_k^{\text{BG}}$  and  $L_k^{\text{BG}} + dG_k^{\text{BG}}$  are elliptic. Their finite dimensional kernel

$$(1.4) \quad \mathcal{H}_L^k(M) := \ker(L_k^{\text{BG}}, G_k^{\text{BG}}), \quad \mathcal{H}^k(M) := \ker(d, G_k^{\text{BG}})$$

are conformally invariant, the elements of  $\mathcal{H}^k(M)$  are named *conformal harmonics*, providing a type of Hodge theory for conformal structure. The operator  $Q_k^{\text{BG}}$  above generalizes Branson  $Q$ -curvature in the sense that it satisfies, as operators on closed  $k$ -forms,

$$\hat{Q}_k^{\text{BG}} = e^{\mu(2k-n)}(Q_k^{\text{BG}} + L_k^{\text{BG}} \mu)$$

if  $\hat{h}_0 = e^{2\mu}h_0$  is another conformal representative.

The general approach of Fefferman-Graham [4] for dealing with conformal invariants is related to Poincaré-Einstein manifolds, roughly speaking it provides a correspondence between Riemannian invariants in the bulk  $(X, g)$  and conformal invariants on the conformal infinity  $(\partial\bar{X}, [h_0])$  of  $(X, g)$ , inspired by the identification of the conformal group of the sphere  $\mathbb{S}^n$  with the isometry group of the hyperbolic space  $\mathbb{H}^{n+1}$ . A smooth Riemannian manifold  $(X, g)$  is said to be a *Poincaré-Einstein manifold* with conformal infinity  $(M, [h_0])$  if the space  $X$  compactifies smoothly to  $\bar{X}$  with boundary  $\partial\bar{X} = M$ , and if there is a boundary defining function of  $\bar{X}$  and some collar neighbourhood  $(0, \epsilon)_x \times \partial\bar{X}$  of the boundary such that

$$(1.5) \quad g = \frac{dx^2 + h_x}{x^2}$$

$$(1.6) \quad \text{Ric}(g) = -ng + O(x^\infty)$$

where  $h_x$  is a one-parameter family of smooth metrics on  $\partial\bar{X}$  such that there exist some family of smooth tensors  $h_x^j$  ( $j \in \mathbb{N}_0$ ) on  $\partial\bar{X}$ , depending smoothly on  $x \in [0, \epsilon)$  with

$$(1.7) \quad \begin{cases} h_x \sim \sum_{j=0}^{\infty} h_x^j (x^n \log x)^j \text{ as } x \rightarrow 0 \text{ if } n+1 \text{ is odd} \\ h_x \text{ is smooth in } x \in [0, \epsilon) \text{ if } n+1 \text{ is even} \end{cases}$$

$$(1.8) \quad h_x|_{x=0} \in [h_0].$$

The tensor  $h_0^1$  is called *obstruction tensor* of  $h_0$ , it is defined in [4] and studied further in [9]. We shall say that  $(X, g)$  is a smooth Poincaré-Einstein manifold if  $x^2g$  extends smoothly on  $\bar{X}$ , i.e. either if  $n+1$  is even or  $n+1$  is odd and  $h_x^j = 0$  for all  $j > 0$ . It is proved in [6] that  $h_0^1 = 0$  implies that  $(X, g)$  is a smooth Poincaré-Einstein manifold.

The boundary  $\partial\bar{X} = \{x = 0\}$  inherits naturally from  $g$  the conformal class  $[h_0]$  of  $h_x|_{x=0}$  since the boundary defining function  $x$  satisfying such conditions are not unique. A fundamental result of Fefferman-Graham [4], which we do not state in full generality, is that for any  $(M, [h_0])$  compact that can be realized as the boundary of smooth compact manifold with boundary  $\bar{X}$ , there is a Poincaré-Einstein manifold  $(X, g)$  for  $(M, [h_0])$ , and  $h_x$  in (1.7) is uniquely determined by  $h_0$  up to order  $O(x^n)$  and up to diffeomorphism which restricts to the Identity on  $M$ . The most basic example is the hyperbolic space  $\mathbb{H}^{n+1}$  which is a smooth Poincaré-Einstein manifold for the canonical conformal structure of the sphere  $\mathbb{S}^n$ , as well as quotients of  $\mathbb{H}^{n+1}$  by convex co-compact groups of isometries.

It has been proved by Mazzeo [16] that<sup>1</sup> for a Poincaré-Einstein manifold  $(X, g)$ , the relative cohomology  $H^k(\bar{X}, \partial\bar{X})$  is canonically isomorphic to the  $L^2$  kernel  $\ker_{L^2}(\Delta_k)$  of the Laplacian  $\Delta_k = (d + \delta_g)^2$  with respect to the metric  $g$ , acting on the bundle  $\Lambda^k(\bar{X})$  of  $k$ -forms if  $k < \frac{n}{2}$ . In other terms the relative cohomology has a basis of  $L^2$  harmonic representatives. In this work, we give an interpretation of the spaces  $\mathcal{H}^k, \mathcal{H}_L^k$  in terms of harmonic forms on the bulk  $X$  with a certain regularity on the compactification  $\bar{X}$ .

**Theorem 1.1.** *Let  $(X^{n+1}, g)$  be an odd dimensional Poincaré-Einstein manifold with conformal infinity  $(M, [h_0])$  and let  $\Delta_k = (d + \delta_g)^2$  be the induced Laplacian on  $k$ -forms on  $X$  where  $0 \leq k < \frac{n}{2} - 1$ . For  $m \in \mathbb{N}$  and  $0 < k < \frac{n}{2} - 1$ , define*

$$K_m^k(\bar{X}) := \{\omega \in C^m(\bar{X}; \Lambda^k(\bar{X})); \Delta_k \omega = 0\},$$

<sup>1</sup>The class of manifold considered by Mazzeo is actually larger and does not require the asymptotic Einstein condition (1.6)

then  $K_m^k(\bar{X})$  is infinite dimensional for  $m < n - 2k + 1$  while it is finite dimensional for  $m \in [n - 2k + 1, n - 1]$  and there is a canonical short exact sequence

$$(1.9) \quad 0 \longrightarrow H^k(\bar{X}, \partial\bar{X}) \xrightarrow{i} K_m^k(\bar{X}) \xrightarrow{r} \mathcal{H}_L^k(M) \longrightarrow 0$$

where  $\mathcal{H}_L^k$  is defined in (1.4) and  $H^k(\bar{X}, \partial\bar{X})$  is the relative cohomology space of degree  $k$  of  $\bar{X}$ ,  $i$  denotes inclusion and  $r$  denotes pull back by the natural inclusion  $\partial\bar{X} \rightarrow \bar{X}$ . If in addition the Fefferman-Graham obstruction tensor of  $(M, [h_0])$  vanishes, i.e. if  $(X, g)$  is a smooth Poincaré-Einstein manifold, then  $K_{n-2k+1}^k(\bar{X}) = K_\infty^k(\bar{X})$ .

When  $k = \frac{n}{2} - 1$ , the same results hold by replacing  $K_{n-2k+1}^k(\bar{X})$  by the set of harmonic forms in  $C^{n-2k+1, \alpha}(\bar{X}, \Lambda^k(\bar{X}))$  for some  $\alpha \in (0, 1)$ .

When  $k = 0$ ,  $K_m^0(\bar{X})$  is infinite dimensional for  $m < n$  while  $K_n^0(\bar{X})$  is finite dimensional and the exact sequence (1.9) holds.

In establishing this Theorem, we show that we can recover the Branson-Gover operators  $L_k^{\text{BG}}, G_k^{\text{BG}}, Q_k^{\text{BG}}$  from harmonic forms on a Poincaré-Einstein manifold with conformal infinity  $(M, [h_0])$ . Let us recall quickly and informally how the GJMS and Branson-Gover operators are defined in [11, 3]. The ambient metric is a Lorentzian metric on  $\tilde{\mathcal{Q}} := M \times (0, \infty)_t \times (-1, 1)_\rho$ , homogeneous of degree 2 in the  $t$  variable, which extends the tautological tensor  $t^2 h_0$  at the cone  $\mathcal{Q} = \{\rho = 0\}$  and with Ricci curvature vanishing to  $n/2 - 1$  order (resp. infinite order) at  $\mathcal{Q}$  when  $n$  is even (resp.  $n$  odd). The GJMS operators  $P_k$  are defined in two ways in [11]: for  $f$  a  $(k - n/2)$ -homogeneous function on  $\mathcal{Q}$ , take a homogeneous extension  $\tilde{f}$  on  $\tilde{\mathcal{Q}}$  and define  $P_k f := [\tilde{\Delta}^k \tilde{f}]|_{\mathcal{Q}}$  where  $\tilde{\Delta}$  is the Laplacian for the ambient metric; the second equivalent way is to consider an extension  $\tilde{f}$  of  $f$  to  $\tilde{\mathcal{Q}}$  which satisfies  $\tilde{\Delta} \tilde{f} = O(\rho^{k-1})$  and, up to multiplicative constant,  $P_k f = [(\rho t^2)^{1-k} \tilde{\Delta} \tilde{f}]|_{\mathcal{Q}}$ . The second definition gives  $P_k$  as an obstruction to extend smoothly from the cone a harmonic homogeneous function  $f$ . The Branson-Gover operators defined in [3] are constructed following the first method in [11] but with many complications due to the fact that one works with bundle valued objects. Our approach is more closely related to the harmonic extension approach of GJMS [11]. We say that a  $k$ -form  $\omega$  is *polyhomogeneous* on  $\bar{X}$  if it is smooth on  $X$  and with an expansion at the boundary  $M = \{x = 0\}$

$$\omega \sim \sum_{j=0}^{\infty} \sum_{\ell=0}^{\ell(j)} x^j \log(x)^\ell (\omega_{j,\ell}^{(t)} + \omega_{j,\ell}^{(n)} \wedge dx)$$

for some forms  $\omega_{j,\ell}^{(t)} \in C^\infty(M, \Lambda^k(M))$  and  $\omega_{j,\ell}^{(n)} \in C^\infty(M, \Lambda^{k-1}(M))$  and some sequence  $j \in \mathbb{N}_0 \rightarrow \ell(j) \in \mathbb{N}_0$ . We show that the Branson-Gover operators appear naturally in the resolution of the absolute or relative Dirichlet type problems for the Laplacian on forms on  $\bar{X}$ .

**Theorem 1.2.** *Let  $(X^{n+1}, g)$  be an odd-dimensional Poincaré-Einstein manifold with conformal infinity  $(M, [h_0])$ , let  $k < \frac{n}{2}$  and  $\alpha \in (0, 1)$ .*

(i) *For any  $\omega_0 \in C^\infty(M, \Lambda^k(M))$ , harmonic forms  $\omega \in C^{\frac{n}{2}-k, \alpha}(\bar{X}, \Lambda^k(\bar{X}))$  with boundary value  $\omega|_M = \omega_0$  exist, are unique modulo  $\ker_{L^2}(\Delta_k)$  and are actually polyhomogeneous with an expansion at  $M$  at order  $O(x^{n-2k+1})$  given by*

$$\begin{aligned} \omega = \omega_0 + \sum_{j=1}^{\frac{n}{2}-k} x^{2j} (\omega_j^{(t)} + \omega_j^{(n)} \wedge \frac{dx}{x}) + x^{n-2k} \log(x) L_k \omega_0 \\ + x^{n-2k+1} \log(x) (G_k \omega_0) \wedge dx + O(x^{n-2k+1}) \end{aligned}$$

where  $L_k, G_k$  are, up to a normalization constant, the Branson-Gover operators in (1.1), (1.2) and  $\omega_j^{(\cdot)}$  are forms on  $M$ .

(ii) *For any closed form  $\omega_0 \in C^\infty(M, \Lambda^{k-1}(M))$ , harmonic forms  $\omega$  such that  $x\omega \in$*

$C^{\frac{n}{2}-k+1,\alpha}(\bar{X}, \Lambda^k(\bar{X}))$  and  $\omega = x^{-1}(\omega_0 \wedge dx) + O(x)$  exist, are unique modulo  $\ker_{L^2}(\Delta_k)$  and  $x\omega$  is polyhomogeneous with expansion at  $M$  given by

$$\omega = \omega_0 \wedge \frac{dx}{x} + \sum_{j=1}^{\frac{n}{2}-k} x^{2j} (\omega_j^{(t)} + \omega_j^{(n)} \wedge \frac{dx}{x}) + x^{n-2k+1} \log(x) (Q_{k-1} \omega_0) \wedge dx + O(x^{n-2k+1})$$

where  $Q_{k-1}$  is, up to a normalization constant, the operator (1.3) of Branson-Gover and  $\omega_j^{(\cdot)}$  are smooth forms on  $M$ .

The Dirichlet problem for functions in this geometric setting is studied by Graham-Zworski [12] and Joshi-Sa Barreto [15]. In a more general setting (but again for functions), it was analyzed by Anderson [1] and Sullivan [20].

We also prove in Subsection 4.6 that, with  $Q_0$  defined by the Theorem above,

$$Q_0 1 = \frac{n(-1)^{\frac{n}{2}+1}}{2^{n-1} \frac{n!}{2} (\frac{n}{2} - 1)!} Q$$

where  $Q$  is Branson  $Q$ -curvature. So  $Q$  can be seen as an obstruction to find a harmonic 1-form  $\omega$  with  $x\omega$  having a high regularity at the boundary and value  $dx$  at the boundary.

In addition, this method allows to obtain the conformal change law of  $L_k, G_k, Q_k$ , the relations between these operators, and some of their analytic properties (e.g. symmetry of  $L_k$  and  $Q_k$ ) see Subsection 4.4 and Section 4.6.

Next, we analyze the set of regular closed and coclosed forms on  $\bar{X}$ . Recall that on a compact manifold  $\bar{X}$  with boundary, equipped with a smooth metric  $\bar{g}$ , there is an isomorphism

$$H^k(\bar{X}) \simeq \{\omega \in C^\infty(\bar{X}, \Lambda^k(\bar{X})); d\omega = \delta_{\bar{g}}\omega = 0, i_{\partial_n}\omega|_{\partial\bar{X}} = 0\}$$

where  $\partial_n$  is a unit normal vector field to the boundary, and the absolute cohomology  $H^k(\bar{X})$  is  $\ker d / \text{Im } d$  where  $d$  acts on smooth forms. Moreover, one has the long exact sequence in cohomology

$$(1.10) \quad \dots \rightarrow H^{k-1}(\partial\bar{X}) \rightarrow H^k(\bar{X}, \partial\bar{X}) \rightarrow H^k(\bar{X}) \rightarrow H^k(\partial\bar{X}) \rightarrow H^{k+1}(\bar{X}, \partial\bar{X}) \rightarrow \dots$$

and all these spaces are represented by forms which are closed and coclosed, the maps in the sequence are canonical with respect to  $\bar{g}$ . In our Poincaré-Einstein case  $(X, g)$ , say when  $k < \frac{n}{2}$ , only the space  $H^k(\bar{X}, \partial\bar{X})$  in the long exact sequence has a canonical basis of closed and coclosed representatives with respect to  $g$  (the  $L^2$  harmonic forms), in particular there is no canonical metric on the boundary induced by  $g$  but only a canonical conformal class. We prove

**Theorem 1.3.** *Let  $(X^{n+1}, g)$  be an odd dimensional Poincaré-Einstein manifold with conformal infinity  $(M, [h_0])$  and let  $k \leq \frac{n}{2}$ . Then the spaces*

$$Z^k(\bar{X}) := \{\omega \in C^{n-2k+1}(\bar{X}, \Lambda^k(\bar{X})); d\omega = \delta_g\omega = 0\}$$

*are finite dimensional and, if the obstruction tensor of  $[h_0]$  vanishes, they are equal to  $\{\omega \in C^\infty(\bar{X}, \Lambda^k(\bar{X})); d\omega = \delta_g\omega = 0\}$ . Then, we have*

*(i) For  $k < \frac{n}{2}$  there is a canonical exact sequence*

$$0 \rightarrow H^k(\bar{X}, M) \rightarrow Z^k(\bar{X}) \rightarrow \mathcal{H}^k(M) \rightarrow H^{k+1}(\bar{X}, M)$$

*where  $\mathcal{H}^k(M)$  is the set of conformal harmonics defined in (1.4).*

*(ii) Let  $[Z^k(\bar{X})]$  and  $[\mathcal{H}^k(M)]$  be respectively the image of  $Z^k(\bar{X})$  and  $\mathcal{H}^k(M)$  by the*

natural cohomology maps  $Z^k(\bar{X}) \rightarrow H^k(\bar{X})$  and  $\mathcal{H}^k(M) \rightarrow H^k(M)$ . Then there exists a canonical complex with respect to  $g$

(1.11)

$$0 \rightarrow \dots \xrightarrow{\iota^k} [Z^k(\bar{X})] \xrightarrow{r^k} [\mathcal{H}^k(M)] \xrightarrow{d_e^k} H^{k+1}(\bar{X}, M) \xrightarrow{\iota^{k+1}} [Z^{k+1}(\bar{X})] \rightarrow \dots \rightarrow H^{\frac{n}{2}}(\bar{X}, M)$$

whose cohomology vanishes except possibly the spaces  $\ker \iota^k / \text{Im } d_e^{k-1}$ . Here  $\iota^k, r^k$  and  $d_e^k$  denote respectively inclusion, restriction to the boundary and composition of  $d$  with harmonic extension (see Section 7).

(iii)  $[\mathcal{H}^k(M)] = H^k(M)$  if and only if  $[Z^k(\bar{X})] = H^k(\bar{X})$  and  $\ker \iota^{k+1} = \text{Im } d_e^k$ . If this holds for all  $k \leq \frac{n}{2}$  this is a canonical realization of (half of) the long exact sequence (1.10) with respect to  $g$ .

The surjectivity of the natural map  $\mathcal{H}^k(M) \rightarrow H^k(M)$  is named  $(k-1)$ -regularity by Branson and Gover, while  $(k-1)$ -strong regularity means that the map is an isomorphism, or equivalently  $\ker L_{k-1} = \ker d$  (see [3, Th.2.6]). Thus,  $(k-1)$  regularity means that the cohomology group can be represented by conformally invariant representatives. If  $H^{k+1}(X, M) = 0$ , our result implies that  $(k-1)$ -regularity means that the absolute cohomology group  $H^k(\bar{X})$  can be represented by  $C^{n-2k+1}(\bar{X}, \Lambda^k(\bar{X}))$  forms in  $\ker d + \delta_g$ . We give a criteria for  $(k-1)$ -regularity:

**Proposition 1.4.** *Let  $(M, [h_0])$  be a compact conformal manifold. If  $Q_k$  is a positive operator on closed forms in the sense that  $\langle Q_k \omega, \omega \rangle_{L^2} \geq 0$  for all  $\omega \in C^\infty(M, \Lambda^k(M)) \cap \ker d$ , then  $\mathcal{H}^k(M) \rightarrow H^k(M)$  is surjective.*

We should also remark that  $(k-1)$ -regularity holds for all  $k = 1, \dots, \frac{n}{2}$  if for instance  $(M, [h_0])$  contains an Einstein metric in  $[h_0]$ , this is a result of Gover and Silhan [7]. If  $n = 4$ ,  $L_{\frac{n}{2}-2} = L_0$  is the Paneitz operator (up to a constant factor) and using a result of Gursky [14], we deduce that if the Yamabe invariant  $Y(M, [h_0])$  is positive and

$$\int_M Q \text{dvol}_{h_0} + \frac{1}{6} Y(M, [h_0])^2 > 0$$

then  $\mathcal{H}^1(M) \simeq H^1(M)$  and there is a basis of conformal harmonics of  $H^1(M)$ .

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## 2. POINCARÉ-EINSTEIN MANIFOLDS AND LAPLACIAN ON FORMS

**2.1. Poincaré-Einstein manifolds.** Let  $(X, g)$  be a Poincaré-Einstein manifold with conformal infinity  $(M, [h])$ . Graham-Lee and Graham [10, 8] proved that for any conformal representative  $h_0 \in [h]$ , there exists a boundary defining function  $x$  of  $M = \partial \bar{X}$  in  $\bar{X}$  such that

$$|dx|_{x^2 g}^2 = 1 \text{ near } \partial \bar{X}, \quad x^2 g|_{TM} = h_0,$$

moreover  $x$  is the unique defining function near  $M$  satisfying these conditions. Such a function is called a *geodesic boundary defining function* and if  $\psi$  is the map  $\psi : [0, \epsilon] \times M \rightarrow \bar{X}$  defined by  $\psi(t, y) := \psi_t(y)$  where  $\psi_t$  is the flow of the gradient  $\nabla^{x^2 g} x$ , then  $\psi$  pulls the metric  $g$  back to

$$\psi^* g = \frac{dt^2 + h_t}{t^2}$$

for some one-parameter family of metrics on  $M$  with  $h_0 = x^2g|_{TM}$ . In other words the special form (1.5) of the metric near infinity is not unique and corresponds canonically to a geodesic boundary defining function, or equivalently to a conformal representative of  $[h_0]$ .

We now discuss the structure of the metric near the boundary, the reader can refer to Fefferman-Graham [6, Th 4.8] for proofs and details. Let us define the endomorphism  $A_x$  on  $TM$  corresponding to  $\partial_x h_x$  with respect to  $h_x$ , i.e. as matrices

$$A_x = h_x^{-1} \partial_x h_x.$$

Then the Einstein condition  $\text{Ric}(g) = -ng$  is equivalent to the following differential equations on  $A_x$

$$\begin{aligned} x\partial_x A_x + (1 - n + \frac{x}{2}\text{Tr}(A_x))A_x &= 2xh_x^{-1}\text{Ric}(h_x) + \text{Tr}(A_x)\text{Id} \\ \delta_{h_x}(\partial_x h_x) &= d\text{Tr}(A_x) \\ \partial_x \text{Tr}(A_x) + \frac{1}{2}|A_x|^2 &= \frac{1}{x}\text{Tr}(A_x) \end{aligned}$$

A consequence of these equations and (1.7) is that if  $\text{Ric}(g) = -ng + O(x^{n-2})$ , then  $h_x$  has an expansion at  $x = 0$  of the form

$$h_x = \begin{cases} h_0 + \sum_{j=1}^{\frac{n}{2}-1} x^{2j} h_{2j} + h_{n,1} x^n \log x + O(x^n) & \text{if } n \text{ is even} \\ h_0 + \sum_{j=1}^{(n-1)/2} x^{2j} h_{2j} + O(x^n) & \text{if } n \text{ is odd} \end{cases}$$

for some tensors  $h_{2j}$  and  $h_{n,1}$  on  $M$ , depending in a natural way on  $h_0$  and covariant derivatives of its Ricci tensor. When  $n$  is even, the tensor  $h_{n,1}$  is the *obstruction tensor* of  $h_0$  in the terminology of Fefferman-Graham [6], it is trace free (with respect to  $h_0$ ) and so the first log term in  $A_x$  is  $nh_0^{-1}h_{n,1}x^{n-1}\log(x)$ . A smooth Poincaré-Einstein manifold such that  $h_x$  has only even powers of  $x$  in the Taylor expansion at  $x = 0$  is called a *smooth even Poincaré-Einstein manifold*. If  $n$  is even and  $h_{n,1} = 0$ , the metric  $h_x$  is a smooth even Poincaré-Einstein manifold. When  $n$  is odd, the term  $\partial_x^n h_x|_{x=0}$  is trace free with respect to  $h_0$ , which implies that  $A_x$  has an even Taylor expansion at  $x = 0$  to order  $O(x^{n-1})$ . If  $\partial_x^n h_x|_{x=0} = 0$ , then  $h_x$  has an even Taylor expansion in powers of  $x$  at  $x = 0$  with all coefficients formally determined by  $h_0$ . The equations satisfied by  $A_x$  easily give (see [4]) the first terms in the expansion

$$(2.1) \quad h_x = h_0 - x^2 \frac{P_0}{2} + O(x^4), \quad \text{where } P_0 = \frac{1}{n-2} \left( 2\text{Ric}_0 - \frac{\text{Scal}_0}{n-1} h_0 \right),$$

$P_0$  is the Schouten tensor of  $h_0$ ,  $\text{Ric}_0$  and  $\text{Scal}_0$  are the Ricci and scalar curvature of  $h_0$ .

**2.2. The Laplacian,  $d$  and  $\delta$ .** Let  $\Lambda^k(\bar{X})$  be the bundle of  $k$ -forms on  $\bar{X}$ . Since for the problem we consider it is somehow quite natural, we will also use along the paper the  $b$ -bundle of  $k$ -forms on  $\bar{X}$  in the sense of [19], it will be denoted  $\Lambda_b^k(\bar{X})$ . This is the exterior product of the  $b$  cotangent bundle  $T_b^* \bar{X}$ , which is canonically isomorphic to  $T^* \bar{X}$  over the interior  $X$  and whose local basis near a point of the boundary  $\partial \bar{X}$  is given by  $dy_1, \dots, dy_n, dx/x$  where  $y_1, \dots, y_n$  are local coordinates on  $\partial \bar{X}$  near this point. We refer the reader to Chapter 2 of [19] for a complete analysis about  $b$ -structures. Of course one can pass from  $\Lambda^k(\bar{X})$  to  $\Lambda_b^k(\bar{X})$  obviously when considering forms on  $X$ . The restriction  $\Lambda_b^k(U_\epsilon)$  of  $\Lambda^k(\bar{X})$  to the collar neighbourhood  $U_\epsilon := [0, \epsilon] \times M$  of  $M$  in  $\bar{X}$  can be decomposed as the direct sum

$$\Lambda_b^k(U_\epsilon) = \Lambda^k(M) \oplus (\Lambda^{k-1}(M) \wedge \frac{dx}{x}) =: \Lambda_t^k \oplus \Lambda_n^k.$$

In this splitting, the exterior derivative  $d$  and its adjoint  $\delta_g$  with respect to  $g$  have the form

$$(2.2) \quad d = \begin{pmatrix} d & 0 \\ (-1)^k x \partial_x & d \end{pmatrix}, \quad \delta = \begin{pmatrix} x^2 \delta_x & (-1)^k \star_x^{-1} x^{-2k+n+3} \partial_x x^{2k-n-2} \star_x \\ 0 & x^2 \delta_x \end{pmatrix}$$

and the Hodge Laplace operator is given by

$$(2.3) \quad \begin{aligned} \Delta_k &= \begin{pmatrix} -(x\partial_x)^2 + (n-2k)x\partial_x & 2(-1)^{k+1}d \\ 0 & -(x\partial_x)^2 + (n-2k+2)x\partial_x \end{pmatrix} \\ &+ \begin{pmatrix} x^2 \Delta_x - x \star_x^{-1} [\partial_x, \star_x] x \partial_x & (-1)^k x [d, \star_x^{-1} [\partial_x, \star_x]] \\ (-1)^{k+1} [x\partial_x, x^2 \delta_x] & x^2 \Delta_x - x \partial_x x \star_x^{-1} [\partial_x, \star_x] \end{pmatrix} \\ &= P + P'. \end{aligned}$$

where here, the subscript  $\cdot_x$  means “with respect to the metric  $h_x$  on  $M$ ” and  $d$  in the matrices is the exterior derivative on  $M$ . Note that  $P$  is the indicial operator of  $\Delta_k$  in the terminology of [19].

If  $H$  is an endomorphism of  $TM$ , we denote  $J(H)$  the operator on  $\Lambda^k(M)$

$$(2.4) \quad J(H)(\alpha_1 \wedge \cdots \wedge \alpha_k) := \sum_{i=1}^k \alpha_1 \wedge \cdots \wedge \alpha_i(H) \wedge \cdots \wedge \alpha_k.$$

When  $H$  is symmetric, a straightforward computation gives  $\star_0 J(H) + J(H) \star_0 = \text{Tr}(H) \star_0$  and so

$$(2.5) \quad [\star_0, J(H)] = 2 \star_0 J(H) - \text{Tr}(H) \star_0$$

Let us define the following operators on  $k$ -forms on  $M$

$$(2.6) \quad E_1 = J(h_0^{-1} P_0) - \frac{\text{Tr}(h_0^{-1} P_0)}{2} \text{Id} = \frac{2J(h_0^{-1} \text{Ric}_0)}{n-2} - \frac{n+2k-2}{2(n-1)(n-2)} \text{Scal}_0 \text{Id}.$$

Using the approximate Einstein equation for  $g$ , we obtain

**Lemma 2.1.** *The operator  $\Delta_k$  has a polyhomogeneous expansion at  $x = 0$  and the first terms in the expansion are given by*

$$(2.7) \quad \begin{aligned} \Delta_k &= P + \sum_{i=1}^{\lfloor \frac{n}{2} \rfloor} x^{2i} \begin{pmatrix} R_i + P_i x \partial_x & S_i \\ S'_i & R'_i + P'_i x \partial_x \end{pmatrix} \\ &+ n x^n \log(x) \begin{pmatrix} J(h_0^{-1} h_{n,1}) x \partial_x & (-1)^{k+1} [d, J(h_0^{-1} h_{n,1})] \\ 0 & J(h_0^{-1} h_{n,1})(n + x \partial_x) \end{pmatrix} + O(x^n) \end{aligned}$$

where the operators  $P_i, P'_i, S_i, S'_i, R_i$  and  $R'_i$  are universal differential operators on  $\Lambda(M)$  that can be expressed in terms of covariant derivatives of the Ricci tensor of  $h_0$ . Moreover the operators  $R_i$  and  $R'_i$  are of order at most 2, the  $S_i, S'_i$  are of order at most 1 and the  $P_i, P'_i$  are of order 0. For instance, we have

$$(2.8) \quad \begin{pmatrix} R_1 + P_1 x \partial_x & S_1 \\ S'_1 & R'_1 + P'_1 x \partial_x \end{pmatrix} = \begin{pmatrix} \Delta_0 - E_1 x \partial_x & (-1)^k [d, E_1] \\ 2(-1)^{k+1} \delta_0 & \Delta_0 - E_1(2 + x \partial_x) \end{pmatrix}$$

where  $A$  is defined in (2.6). If  $k = 0$ , the  $x^n \log(x)$  coefficient vanishes. Finally, if  $(X, g)$  is smooth Poincaré-Einstein, then  $\Delta_k$  is a smooth differential operator on  $\bar{X}$ , and if  $(X, g)$  is smooth even Poincaré-Einstein, then  $\Delta_k$  has an even expansion.

*Proof:* The polyhomogeneity comes from that of the metric  $g$ . It is moreover a smooth expansion if  $x^2 g$  is smooth on  $\bar{X}$ . A priori, by (2.3) the first  $\log x$  term in the expansion of  $\Delta$  at  $x = 0$  appear at order (at least)  $x^n \log x$  and it comes from the diagonal terms in  $P'$  in (2.3). Let us define  $p = \lfloor \frac{n}{2} \rfloor$  so that the metric  $h_x$  has even powers in its expansion at  $x = 0$  up to order  $x^{2p+1}$ . We set  $D$  the Levi-Civita connection of the metric  $x^2 g = dx^2 + h_x$ . Since  $D_{\partial_x} \partial_x = 0$  and  $D_{\partial_x} \partial_{y_i} = \frac{1}{2} \sum_{jk} \partial_x h_{ij} h^{kj} \partial_{y_k}$ , the matrix  $O_x$  of the

parallel transport along the geodesic  $x \mapsto (x, y)$  (with respect to the basis  $(\partial_{y_i})$ ) satisfies  $D_{\partial_x} O_x(\partial_{y_i}) = 0$ , hence  $\partial_x O_x = -\frac{1}{2} A_x \times O_x$  where  $A_x$  is the endomorphism  $h_x^{-1} \partial_x h_x$ . Note that  $A_x$  has a Taylor expansion with only odd powers of  $x$  up to  $x^{2p}$  and the first log term is  $nh_0^{-1} h_{n,1} x^{n-1} \log(x)$ . We infer that  $O_x$  is polyhomogeneous in the  $x$  variable and has only even powers of  $x$  in its Taylor expansion up to  $x^{2p}$ , the first log term is  $-\frac{h_0^{-1} h_{n,1}}{2} x^n \log(x)$ . By (2.1), we have  $\partial_x^2 h|_{x=0} = -P_0$ , hence

$$A_x = -xh_0^{-1}P_0 + O(x^2), \quad O_x = \text{Id} + \frac{1}{4}x^2h_0^{-1}P_0 + O(x^3).$$

We note also  $O_x$  the parallel transport map. Now the operator  $I_x(\alpha_1 \wedge \cdots \wedge \alpha_k) = \alpha_1(O_x) \wedge \cdots \wedge \alpha_k(O_x)$  is an isometry from  $\Lambda^k(M, h_x)$  to  $\Lambda^k(M, h_0)$ . So we have  $\star_x = I_x^{-1} \star_0 I_x$  and we infer that  $\star_x$  itself is an operator with a polyhomogeneous expansion in  $x$  and with only even powers of  $x$  in its Taylor expansion up to  $x^{2p}$ , the first log term being  $\frac{1}{2}x^n \log(x)[J(h_0^{-1}h_{n,1}), \star_0] = -x^n \log(x) \star_0 J(h_0^{-1}h_{n,1})$  by (2.5). Since we have

$$[\partial_x, \star_x] = \partial_x(\star_x), \quad \partial_x(\star_x)|_{x=0} = [\star_0, \partial_x I_x|_{x=0}] = 0 \quad \text{and} \quad \partial_x^2(\star_x)|_{x=0} = [\star_0, \partial_x^2 I_x|_{x=0}]$$

we get that  $[\partial_x, \star_x]$  is polyhomogeneous with only odd powers of  $x$  up to order  $x^{2p}$ , with first log term  $-nx^{n-1} \log(x) \star_0 J(h_0^{-1}h_{n,1})$ , and that

$$[\partial_x, \star_x] = \partial_x(\star_x) = x \star_0 (J(h_0^{-1}P_0) - \frac{\text{Scal}_0}{2(n-1)} \text{Id}) + O(x^2).$$

Since  $\delta_x = (-1)^k \star_x^{-1} d\star_x$ , the operators  $x\delta_x$  and  $x^2[\star_x^{-1}[\partial_x, \star_x], \delta_x]$  are odd in  $x$  up to  $O(x^{2p+2})$ . By the same way,  $x^2[d, \star_x^{-1}[\partial_x, \star_x]]$  is odd up to order  $x^{2p+2}$  and the operators  $\star_x^{-1}[\partial_x, \star_x]x(k-x\partial_x)$ ,  $x^2\Delta_x$  and  $(k-\partial_x x)x\star_x^{-1}[\star_x, \partial_x]$  are even in  $x$  up to  $O(x^{2p+1})$ . This achieves the proof by gathering all these facts.  $\square$

**2.3. Indicial equations.** We give the indicial equations satisfied by  $\Delta_k$ , which are essential to the construction of formal power series solutions of  $\Delta_k \omega = 0$ .

**Notation:** If  $f$  is a function on  $\bar{X}$  and  $\omega$  a  $k$ -form defined near the boundary, we will say that  $\omega$  is a  $O_n(f)$  (resp.  $O_t(f)$ ) if its  $\Lambda_n^k$  (resp.  $\Lambda_t^k$ ) components are  $O(f)$ .

For  $\lambda \in \mathbb{C}$ , the operator  $x^{-\lambda} \Delta_k x^\lambda$  can be considered near the boundary as a family of operators on  $\Lambda_t^k \oplus \Lambda_n^k$  depending on  $(x, \lambda)$ , and for any  $\omega \in C^\infty(U_\epsilon, \Lambda_t^k \oplus \Lambda_n^k)$  one has

$$(2.9) \quad x^{-\lambda} \Delta_k(x^\lambda \omega) = P_\lambda \left( \omega_0^{(t)} + \omega_0^{(n)} \wedge \frac{dx}{x} \right) + O(x)$$

where  $P_\lambda := x^{-\lambda} P x^\lambda$ ,  $\omega_0^{(t)} = (i_{x\partial_x}(\omega \wedge \frac{dx}{x}))|_{x=0}$  and  $\omega_0^{(n)} := (i_{x\partial_x} \omega)|_{x=0}$ . The operator  $P_\lambda$  is named indicial family and is a one-parameter family of operators on  $\Lambda_n^k \oplus \Lambda_t^k$  viewed as a bundle over  $M$ , its expression is

$$(2.10) \quad P_\lambda = \begin{pmatrix} -\lambda^2 + (n-2k)\lambda & 2(-1)^{k+1}d \\ 0 & -\lambda^2 + (n-2k+2)\lambda \end{pmatrix}$$

The *indicial roots* of  $\Delta_k$  are the  $\lambda \in \mathbb{C}$  such that  $P_\lambda$  is not invertible on the set of smooth sections of  $\Lambda_t^k \oplus \Lambda_n^k$  over  $M$ , i.e. on  $C^\infty(M, \Lambda^k(M) \oplus \Lambda^{k-1}(M))$ . In our case, a simple computation shows that these are given by  $0, n-2k, 0, n-2k+2$ . The first two roots are roots in the  $\Lambda_t^k$  component and the last two are roots in the  $\Lambda_n^k$  component. In particular, this proves that for  $j$  not a root, and  $(\omega_0^{(t)}, \omega_0^{(n)}) \in \Lambda^k(M) \oplus \Lambda^{k-1}(M)$ , there exists a unique pair  $(\alpha_0^{(t)}, \alpha_0^{(n)}) \in \Lambda^k(M) \oplus \Lambda^{k-1}(M)$  such that near  $M$

$$\Delta_k \left( x^j \alpha_0^{(t)} + x^j \alpha_0^{(n)} \wedge \frac{dx}{x} \right) = x^j \left( \omega_0^{(t)} + \omega_0^{(n)} \wedge \frac{dx}{x} \right) + O(x^{j+1})$$

More precisely, and including coefficients with log terms, we have for  $l \in \mathbb{N}^*$  (resp.  $l = 0$ )

$$(2.11) \quad \Delta_k x^j \log^l(x) \begin{pmatrix} \omega_0^{(t)} \\ \omega_0^{(n)} \end{pmatrix} = x^j \log^l(x) \begin{pmatrix} j(n-2k-j)\omega_0^{(t)} + 2(-1)^{k+1}d\omega_0^{(n)} \\ j(n-2k+2-j)\omega_0^{(n)} \end{pmatrix} \\ + O(x^j \log^{l-1}(x)) \quad (\text{resp. } + O(x^{j+1}))$$

if  $\omega_0^{(t)}, \omega_0^{(n)} \in C^\infty(M, \Lambda^k(M) \oplus \Lambda^{k-1}(M))$ , and in the critical cases, for any  $l \in \mathbb{N}_0 = \{0\} \cup \mathbb{N}$

$$(2.12) \quad \Delta_k(\log^l(x)\omega_0^{(t)}) = l(n-2k)\log^{l-1}(x)\omega_0^{(t)} - l(l-1)\log^{l-2}(x)\omega_0^{(t)} + O(x^2 \log x) \\ \Delta_k(x^{n-2k}\log^l(x)\omega_0^{(t)}) = l(2k-n)x^{n-2k}\log^{l-1}(x)\omega_0^{(t)} - l(l-1)x^{n-2k}\log^{l-2}(x)\omega_0^{(t)} \\ + O(x^{n-2k+2}\log^l(x)) \\ \Delta_k(x^{n-2k+2}\log^l(x)\omega^{(n)} \wedge \frac{dx}{x}) = l(2k-2-n)x^{n-2k+2}\log^{l-1}(x)\omega^{(n)} \wedge \frac{dx}{x} \\ - l(l-1)x^{n-2k+2}\log^{l-2}(x)\omega^{(n)} \wedge \frac{dx}{x} + O(x^{n-2k+3}\log^l(x)).$$

### 3. ABSOLUTE AND RELATIVE DIRICHLET PROBLEMS

The goal of this section is to solve the Dirichlet type problems for  $\Delta_k$  when  $k < \frac{n}{2}$  for the two natural boundary conditions. Note that the vector field  $x\partial_x$  can be seen as the unit, normal, inward vector field to  $M$  in  $\bar{X}$ . A  $k$ -form  $\omega \in \Lambda_b^k(\bar{X})$  is said to satisfy the absolute (resp. the relative) boundary condition if

$$\lim_{x \rightarrow 0} i_{x\partial_x} \omega = 0 \quad (\text{resp. } \lim_{x \rightarrow 0} i_{x\partial_x} (\frac{dx}{x} \wedge \omega) = 0).$$

We denote  $C^{p,\alpha}(\bar{X}, \Lambda_b^k(\bar{X}))$  the sections of  $\Lambda_b^k(\bar{X})$  which are  $C^{p,\alpha}$ , equivalently  $i_{x\partial_x} \omega$  and  $i_{x\partial_x} (\frac{dx}{x} \wedge \omega)$  are  $C^{p,\alpha}$  on  $\bar{X}$ .

#### 3.1. Absolute boundary condition.

**Proposition 3.1.** *Let  $k < n/2$ ,  $\alpha \in (0, 1)$  and  $\omega_0 \in C^\infty(M, \Lambda^k(M))$ .*

(i) *There exists a solution  $\omega$  to the following absolute Dirichlet problem:*

$$(3.1) \quad \begin{cases} \omega \in C^{n-2k-1,\alpha}(\bar{X}, \Lambda^k(\bar{X})), \\ \Delta_k \omega = 0 \text{ on } X, \\ \omega|_M = \omega_0, \lim_{x \rightarrow 0} i_{x\partial_x} \omega = 0. \end{cases}$$

Moreover, this solution is unique modulo the  $L^2$  kernel of  $\Delta_k$ .

(ii) *The solution  $\omega$  is smooth in  $\bar{X}$  when  $n$  is odd, while it is polyhomogeneous when  $n$  is even with an expansion at order  $x^n$  of the form*

$$(3.2) \quad \omega = \sum_{j=0}^{n-1} x^j \omega_j^{(t)} + \sum_{j=2}^{n-1} x^j \omega_j^{(n)} \wedge \frac{dx}{x} + \log x \left( \sum_{j=n-2k}^{n-1} x^j \omega_{j,1}^{(t)} + \sum_{j=n-2k+2}^n x^j \omega_{j,1}^{(n)} \wedge \frac{dx}{x} \right) \\ + \begin{cases} O_t(x^n \log x) + O_n(x^{n+1} \log x) & \text{if } k > 0 \\ O(x^n) & \text{if } k = 0 \end{cases}$$

as  $x \rightarrow 0$ , where  $\omega_j^{(\cdot)}, \omega_{j,1}^{(\cdot)}$  are smooth forms on  $M$ . Moreover, we have

$$\omega_j^{(t)} = P_j^{(t)} \omega_0 \text{ for } j < n-2k, \quad \omega_j^{(n)} = P_j^{(n)} \omega_0 \text{ for } j < n-2k+2 \\ \omega_{n-2k,1}^{(t)} = P_{n-2k,1}^{(t)} \omega_0$$

where  $P_j^{(t)}, P_j^{(n)}, P_{n-2k,1}^{(t)}$  are universal smooth differential operators on  $\Lambda(M)$  depending naturally on covariant derivatives of the curvature tensor of  $h_0$ .

(iii) *If  $n$  is even and  $(X, g)$  is a smooth Poincaré-Einstein manifold, then we have  $\omega = \omega_1 + x^{n-2k} \log(x) \omega_2$  for some forms  $\omega_1, \omega_2 \in C^\infty(\bar{X}, \Lambda_b^k(\bar{X}))$  with  $\omega_2 = O(x^\infty)$  if and only*

if  $\omega_{n-2k,1}^{(t)} = \omega_{n-2k+2,1}^{(n)} = 0$ .

(iv)  $\omega$  satisfies  $\delta_g \omega = 0$ . If in addition  $\omega_0$  is closed, then  $d\omega \in \ker_{L^2}(\Delta_{k+1})$  (and  $d\omega = 0$  when  $k = \frac{n-1}{2}$ ).

3.1.1. *Proof of Proposition 3.1.* To prove this Proposition, we first need a result of Mazzeo [16] (note that the ambient manifold has dimension  $n$  in [16] and  $n+1$  in this paper):

**Theorem 3.2 (Mazzeo).** *For  $k < n/2$ , the operator  $\Delta_k$  is Fredholm and there exists a pseudodifferential inverse  $E$ , bounded on  $L^2(X)$ , such that  $\Delta_k E = I - \Pi_0$  where  $\Pi_0$  is the projection on the finite dimensional space  $\ker_{L^2}(\Delta_k)$ . This implies an isomorphism between  $\ker_{L^2}(\Delta_k)$  and the relative cohomology  $H^k(\bar{X}, \partial\bar{X})$  of  $\bar{X}$ . Moreover any  $L^2$  harmonic form  $\alpha$  is polyhomogeneous with an expansion near  $\partial\bar{X}$  of the form*

$$(3.3) \quad \alpha \sim x^{n-2k} \sum_{j=0}^{\infty} \sum_{l=0}^{l(j)} (\alpha_{j,l}^{(t)} x^j \log(x)^l + x^{j+2} \log(x)^l \alpha_{j,l}^{(n)}) \wedge \frac{dx}{x}$$

for some  $\alpha_{j,l}^{(t)} \in C^\infty(M, \Lambda^k(M))$ ,  $\alpha_{j,l}^{(n)} \in C^\infty(M, \Lambda^{k-1}(M))$  and some sequence  $l : \mathbb{N}_0 \rightarrow \mathbb{N}_0$ . In addition  $E$  maps the space  $\{\omega \in C^\infty(\bar{X}, \Lambda^k(\bar{X})); \omega = O(x^\infty)\}$  into polyhomogeneous forms on  $\bar{X}$  with a behaviour like (3.3) near  $M$ .

**Remark:** By using duality through the Hodge star operator  $\star_g$ , one obtains trivially a corresponding result for the case  $k > \frac{n}{2} + 1$ . In particular, this gives  $\ker_{L^2}(\Delta_k) \simeq H^k(\bar{X})$  for  $k > \frac{n}{2} + 1$ . It should be noticed that in the case  $k = n/2, (n+1)/2$  and  $n/2 + 1$ , [16] does not give a bounded pseudodifferential inverse and actually the Laplacian is not Fredholm in these cases: for  $k = n/2$  or  $k = n/2 + 1$  the range is not closed while for  $k = (n+1)/2$  it has infinite dimensional kernel.

We can make the second part of this theorem more precise thanks to the indicial identities obtained by (2.3).

**Corollary 3.3.** *For  $k < n/2$ , any  $L^2$  harmonic  $k$ -form  $\alpha$  on  $(X, g)$  is polyhomogeneous and has an expansion at order  $x^n \log x$  of the form*

$$\alpha = x^{n-2k+2} \left( \sum_{j=0}^{n-1} x^j \alpha_j^{(t)} + \sum_{j=0}^{n-1} x^j \alpha_j^{(n)} \right) \wedge \frac{dx}{x} + O(x^n \log x)$$

where  $\alpha_j^{(\cdot)}$  are smooth forms on  $M$ . If in addition the metric  $(X, g)$  is a smooth Poincaré-Einstein manifold, then  $\alpha \in x^{n-2k+2} C^\infty(\bar{X}, \Lambda^k(\bar{X}))$  and  $E$  maps

$$E : \{\omega \in C^\infty(\bar{X}, \Lambda^k(\bar{X})); \omega = O(x^\infty), \Pi_0 \omega = 0\} \longrightarrow x^{n-2k} C^\infty(\bar{X}, \Lambda^k(\bar{X})).$$

*Proof:* Note that if

$$\alpha \sim x^{n-2k} \sum_{j=0}^{\infty} \sum_{l=0}^{l(j)} (\alpha_{j,l}^{(t)} x^j \log(x)^l + x^{j+2} \log(x)^l \alpha_{j,l}^{(n)}) \wedge \frac{dx}{x} \quad \text{and } \Delta_k \alpha = O(x^\infty),$$

then the indicial equations in Subsection 2.3 and Lemma 2.1 imply that  $l(0) = 0$  and  $l(j) \leq 1$  for all  $j = 1, \dots, n-1$  (and for all  $j > 0$  if  $h_x$  is smooth in  $x$ ). Moreover since  $d\alpha = 0$  for any  $\alpha \in \ker_{L^2}(\Delta_k)$ , we first obtain from (2.2) that  $\alpha_{0,0}^{(t)} = 0$  and so, by (2.12) that  $l(j) = 0$  for all  $j = 0, \dots, n-1$  (and for all  $j > 0$  if  $h_x$  is smooth). The mapping property of  $E$  is straightforward by the same type of arguments and the fact that  $\Delta_k E\omega = O(x^\infty)$  for  $\omega = O(x^\infty)$  such that  $\Pi_0 \omega = 0$ .

We will now use the relations (2.9), (2.11) and (2.12) to show that the jet of a solution  $\omega$  to the Dirichlet problem in Proposition 3.1 is partly determined. Let  $\omega_0 \in$

$C^\infty(M, \Lambda^k(M))$ . Using (2.9) and the form (2.7) of  $\Delta$ , we can construct a smooth form  $\omega_{F_1}$  on  $\bar{X}$ , solution to the problem

$$(3.4) \quad \begin{cases} \Delta_k \omega_{F_1} = O_t(x^{n-2k}) + O_n(x^{n-2k+2}) \\ \omega_{F_1}|_M = \omega_0 \end{cases}$$

it can be taken as a polynomial in  $x$

$$(3.5) \quad \omega_{F_1} = \sum_{2j=0}^{n-2k-1} x^{2j} \omega_{2j}^{(t)} + \sum_{2l=2}^{n-2k+1} x^{2l} \omega_{2l}^{(n)} \wedge \frac{dx}{x}$$

and it is the unique solution of (3.4) modulo  $O_t(x^{n-2k}) + O_n(x^{n-2k+2})$ . Moreover, by (2.7) and parity arguments, we see that when  $n$  is odd, the remaining term in (3.4) can be replaced by  $O_t(x^{n-2k+1}) + O_n(x^{n-2k+3})$  (recall also that  $h_x$  is smooth in that case). By construction, the  $\omega_j^{(t)}, \omega_n^{(n)}$  are forms on  $M$  which can be expressed as a differential operators  $P_j^{(t)}, P_j^{(n)}$  on  $M$  acting on  $\omega_0$ , determined by the expansion of  $P$  given in (2.7), i.e. by  $h_0$  and the covariant derivatives of its curvature tensor.

The indicial factor in (2.9) vanishes if and only if  $j = n - 2k, l = n - 2k + 2$  and  $n$  is even. Therefore, if  $n$  is odd, we can continue the construction and there is a formal series

$$\omega_\infty = \sum_{j=0}^{\infty} x^j (\omega_j^{(t)} + \omega_j^{(n)} \wedge dx)$$

such that  $\Delta_k \omega_\infty = O(x^\infty)$ . The formal form  $\omega_\infty$  can be realized by Borel's Lemma<sup>2</sup>, in the sense that there exists a form  $\omega'_\infty \in C^\infty(\bar{X}, \Lambda^k(\bar{X}))$  with the same asymptotic expansion than  $\omega_\infty$  at all order and then  $\Delta_k \omega'_\infty = O(x^\infty)$ .

Now for  $n$  even, we need to add log terms to continue the parametrix: by (2.12) one can modify  $\omega_{F_1}$  to

$$(3.6) \quad \omega_{F_2} = \omega_{F_1} + x^{n-2k} \log(x) \omega_{n-2k,1}^{(t)}, \quad \omega_{n-2k,1}^{(t)} = \frac{1}{n-2k} \left[ x^{-n+2k} \Delta_k \omega_{F_1} \right]_{|x=0}$$

such that  $\Delta_k \omega_{F_2} = O(x^{n-2k+2} \log x)$ . Actually, using (2.7) and parity arguments once more, we see that

$$(3.7) \quad \Delta_k \omega_{F_2} = 2(-1)^{k+1} x^{n-2k+2} \log x (\delta_0 \omega_{n-2k,1}^{(t)}) \wedge \frac{dx}{x} + O_t(x^{n-2k+2} \log x) + O_n(x^{n-2k+2}).$$

Now we want to show

**Lemma 3.4.** *The  $k$ -form  $\omega_{n-2k,1}^{(t)}$  on  $M$  satisfies  $\delta_0 \omega_{n-2k,1}^{(t)} = 0$ .*

*Proof:* From (3.7), and the expression of  $\delta$ , we obtain

$$\delta_g \Delta_k \omega_{F_2} = -2x^{n-2k+2} \delta_0 \omega_{n-2k,1}^{(t)} + O(x^{n-2k+3} \log x).$$

But  $\delta_g \Delta_k \omega_{F_2} = \Delta_{k-1} \delta_g \omega_{F_2}$  and

$$\delta_g \omega_{F_2} = \sum_{j=2}^{n-2k+2} x^j \omega_j^{(t)} + \sum_{l=3}^{n-2k+3} x^l \omega_l^{(n)} \wedge \frac{dx}{x} + x^{n-2k+2} \log(x) \delta_0 \omega_{n-2k,1}^{(t)} + O(x^{n-2k+3} \log x)$$

for some forms  $\omega_j^{(\cdot)}$  on  $M$ , so by uniqueness of (3.4) and the fact that  $\delta_g \omega_{F_2} = O(x^2)$  we deduce that

$$\delta_g \omega_{F_2} = x^{n-2k+2} \omega_{n-2k+2}^{(t)} + x^{n-2k+2} \log(x) \delta_0 \omega_{n-2k,1}^{(t)} + O(x^{n-2k+3} \log x).$$

<sup>2</sup>Borel's Lemma states that if  $(f_{k,l})_{l,k \in \mathbb{N}_0}$  is a given sequence of smooth functions on  $\partial \bar{X}$  such that, for each  $k$ ,  $f_{k,l}(y) = 0$  for all but finitely many  $l$ , then there exists a smooth function  $f$  in  $X$  with an asymptotic expansion at  $\partial \bar{X} = \{x = 0\}$  of the form  $f(x, y) \sim \sum_{k=0}^{\infty} \sum_{l=0}^{\infty} f_{k,l}(y) x^k (\log x)^l$ .

Using now (2.9) and (2.12), we obtain  $\Delta_{k-1}\delta_g\omega_{F_2} = (2k - n - 2)x^{n-2k+2}\delta_0\omega_{n-2k,1}^{(t)} + O(x^{n-2k+3}\log x)$ , and since  $k < \frac{n}{2}$  this implies  $\delta_0\omega_{n-2k,1}^{(t)} = 0$ .  $\square$

We infer that there is no term of order  $x^{n-2k+2}\log x$  in the  $\Lambda_n^k$  part of  $\Delta_k\omega_{F_2}$  and we can continue to solve the problem modulo  $O(x^\infty)$  using formal power series with log terms using the indicial equations. The formal solution when  $n$  is even will be given by

$$(3.8) \quad \begin{aligned} \omega_\infty &= \sum_{j=0}^{\frac{n}{2}-1} x^{2j}\omega_{2j}^{(t)} + \sum_{j=1}^{\frac{n}{2}} x^{2j}\omega_{2j}^{(n)} \wedge \frac{dx}{x} + \sum_{j=\frac{n}{2}-k}^{\frac{n}{2}-1} x^{2j}\log(x)\omega_{2j,1}^{(t)} \\ &+ \sum_{j=\frac{n}{2}-k+1}^{\frac{n}{2}} x^{2j}\log(x)\omega_{2j,1}^{(n)} \wedge \frac{dx}{x} + x^n \sum_{j=0}^{\infty} \sum_{l=0}^{j+1} (\omega_{n+j,l}^{(t)} + x\omega_{n+j,l}^{(n)} \wedge \frac{dx}{x}) x^j (\log x)^l \end{aligned}$$

which again is realized through Borel's Lemma to have  $\Delta_k\omega_\infty = O(x^\infty)$ . Notice that when the metric  $h_x$  is smooth, the second line in (3.8) has  $\omega_{j,l}^{(t)} = \omega_{j,l}^{(n)} = 0$  for  $l > 1$  since these terms come from the log terms of the expansion of  $h_x$  in (1.7) (and thus of  $\Delta_k$ ). The terms  $(\omega_j^{(t)})_{j < n-2k}$ ,  $(\omega_j^{(n)})_{j < n-2k+2}$  and  $\omega_{n-2k,1}^{(t)}$  are formally determined by  $\omega_0$  and are expressed as a differential operator on  $M$  acting on  $\omega_0$ , the terms  $\omega_{n-2k}^{(t)}$ ,  $\omega_{n-2k+2}^{(n)}$  are formally undetermined, the remaining terms are formally determined by  $\omega_0$ ,  $\omega_{n-2k}^{(t)}$  and  $\omega_{n-2k+2}^{(n)}$ .

So we have proved

**Proposition 3.5 (Formal solution).** *Let  $\omega_0, v^{(t)}, v^{(n)} \in C^\infty(M, \Lambda(M))$ , then there exists a form  $\omega_\infty \in C^{n-2k-1, \alpha}(\bar{X}, \Lambda_b^k(\bar{X}))$  with  $\alpha \in [0, 1)$ , unique modulo  $O(x^\infty)$ , which is smooth on  $\bar{X}$  when  $n$  is odd and with a polyhomogeneous expansion at  $\partial\bar{X}$  of the form (3.8) when  $n$  is even, such that  $\Delta_k\omega_\infty = O(x^\infty)$ ,  $\omega_\infty|_{\partial\bar{X}} = \omega_0$ ,  $\omega_{n-2k}^{(t)} = v^{(t)}$  and  $\omega_{n-2k+2}^{(n)} = v^{(n)}$  in the expansion (3.8).*

To correct the approximate solution and obtain a true harmonic form, we add  $-E\Delta_k(\omega_\infty)$  to  $\omega_\infty$  and so

$$\Delta_k(\omega_\infty - E\Delta_k\omega_\infty) = \Pi_0\Delta_k\omega_\infty.$$

We want to prove that  $\Pi_0\Delta_k\omega_\infty = 0$  or equivalently that  $\langle \Delta_k\omega_\infty, \alpha \rangle = 0$  for any  $\alpha \in \ker_{L^2}(\Delta_k)$ . For that, we use Green's formula on  $\{x \geq \varepsilon\}$  and let  $\varepsilon \rightarrow 0$ , together with the asymptotic  $\alpha = O(x^{n-2k+1})$  obtained from Theorem 3.2,  $d\alpha = 0$  and  $\delta\alpha = 0$ :

$$\int_{x \geq \varepsilon} \langle \Delta\omega_\infty, \alpha \rangle d\text{vol}_g = (-1)^n \int_{x=\varepsilon} (\star_g d\omega_\infty) \wedge \alpha - (\star_g \alpha) \wedge \delta\omega_\infty = O(\varepsilon) \rightarrow_{\varepsilon \rightarrow 0} 0.$$

In view of the mapping properties of  $E$  from Theorem 3.2, we have thus proved that  $\omega = \omega_\infty - E\Delta_k\omega_\infty$  is a harmonic  $k$ -form of  $X$  such that  $\omega|_M = \omega_0$ , with an asymptotic of the form (3.8) when  $n$  is even and smooth on  $\bar{X}$  when  $n$  is odd, such that

$$\omega - \omega_{F_2} = O_t(x^{n-2k}) + O_n(x^{n-2k+2})$$

and with  $C^{n-2k-1, \alpha}(\bar{X}, \Lambda^k(\bar{X}))$  regularity.

Let us now consider the problem of uniqueness. If one assumes polyhomogeneity of the solution of  $\Delta_k\omega = 0$  with boundary condition  $\omega = \omega_0 + o(x)$ , the construction above with formal series arguments and indicial equations shows that  $\omega$  is unique up to  $O_t(x^{n-2k}) + O_n(x^{n-2k+2})$ , i.e. the first positive indicial roots, then of course two such solutions would differ by an  $L^2$  harmonic form if  $k < \frac{n}{2}$ . Indeed, an easy computation shows that

**Remark 3.6.** *A polyhomeogeneous  $k$ -form in  $O_t(x^{\frac{n}{2}-k+\epsilon}) + O_n(x^{\frac{n}{2}-k+1+\epsilon})$  for some  $\epsilon > 0$  is in  $L^2(X)$ .*

This gives

**Lemma 3.7.** *Polyhomogeneous forms satisfying  $\Delta_k \omega$  and  $\omega = \omega_0 + o(x)$  are unique modulo the  $L^2$  kernel of  $\Delta_k$ .*

Here, since we want a sharp condition on regularity for uniqueness, i.e. we do not assume polyhomogeneity but  $C^{n-2k-1, \alpha}$  regularity, we first need a preliminary result. Let  $H^s(\Lambda^k(M))$  be the Sobolev space of order  $s \in \mathbb{Z}$  with  $k$ -forms values, which we will also denote by  $H^s(M)$  to simplify. The sections of the bundle  $\Lambda_t^k \oplus \Lambda_n^k$  over  $M$  are equipped with the natural Sobolev norm  $\|\cdot\|_{H^s(M)}$  induced by  $H^s(\Lambda^k(M) \oplus \Lambda^{k-1}(M))$ . The following property is proved by Mazzeo [18, Th. 7.3]<sup>3</sup>

**Lemma 3.8 (Mazzeo).** *Let  $k < n/2$  and let  $\omega \in x^\alpha L^2(\Lambda^k(X), dvol_g)$  with  $\alpha < -\frac{n}{2}$  such that  $\Delta_k \omega = 0$ , then for all  $N \in \mathbb{N}$ , there exist some forms  $\omega_{j,l}^{(t)}, \omega_{j,l}^{(n)} \in H^{-N}(M)$  for  $j, l \in \mathbb{N}_0$  and some sequence  $l: \mathbb{N}_0 \rightarrow \mathbb{N}_0$  such that*

$$(3.9) \quad \left\| \omega - \sum_{j=0}^{N-3} \sum_{l=0}^{l(j)} x^j (\log x)^l (\omega_{j,l}^{(t)} - \omega_{j,l}^{(n)}) \wedge \frac{dx}{x} \right\|_{H^{-N}(M)} = O(x^{N-2-\varepsilon})$$

for all  $\varepsilon > 0$ .

Let  $\omega, \omega'$  be two harmonic forms which are  $C^{n-2k-1, \alpha}(\bar{X}, \Lambda^k(\bar{X}))$  and which coincide on the boundary, we want to show that their Taylor expansions at  $x = 0$  coincide to order  $n - 2k - 1$ . Using Lemma 3.8 with  $N$  large enough, we see that the arguments used above on formal series (based on the indicial equations) also apply by considering norms  $\|\cdot\|_{H^{-N}(M)}$  on  $\Lambda_t^k \oplus \Lambda_n^k$ , in particular that  $l(j) = 0$  for  $j = 0, \dots, n - 2k - 1$  in (3.9) for both  $\omega$  and  $\omega'$ , and that their coefficients of  $x^j$  for  $j = 0, \dots, n - 2k - 1$  in the weak expansion (3.9) are the same for  $\omega$  and  $\omega'$ , these are given by  $\omega_{j,0}^t = P_j^{(t)} \omega_0$  and  $\omega_j^{(n)} = P_j^{(n)} \omega_0$  (and are then continuous on  $M$  since  $\omega \in C^{n-2k+1}(\bar{X}, \Lambda^k(\bar{X}))$ ). But by uniqueness of the expansion (3.9) and the regularity assumption on  $\omega, \omega'$ , this implies that  $\omega_j^{(t)}, \omega_j^{(n)}$  are the coefficients in the Taylor expansion of both  $\omega$  and  $\omega'$  to order  $n - 2k - 1$ . The extra Hölder regularity then gives that  $\|\omega - \omega'\|_{L^\infty(M)} = O(x^{n-2k-1+\alpha})$ , but then this implies that  $\omega - \omega' \in \ker_{L^2}(\Delta_k)$  thus it is in the  $L^2$  kernel of  $\Delta_k$ , so our construction is unique modulo  $\ker_{L^2}(\Delta_k)$ . This ends the proof of the solution of (3.1).

Now to deal with (iv), we notice that  $d\omega$  is solution of the problem (3.1) for  $(k+1)$ -forms with the additional condition that the boundary value is  $d\omega_0 = 0$ . When  $k+1 < \frac{n}{2}$ , we can then apply Proposition 3.1 (i), when  $k \geq \frac{n}{2} - 1$ , we have  $d\omega = O(x^2)$  and  $\Delta_{k+1} d\omega = 0$ . However, the discussion below in Subsections 3.1.2 and 3.1.3 about the solutions of  $\Delta_{k+1} \omega = O(x^\infty)$  gives the same result, namely that  $d\omega \in \ker_{L^2}(\Delta_{k+1})$  if  $\omega$  is a solution of (3.1) with  $k = \frac{n}{2} - 1$  or  $k = \frac{n-1}{2}$ .

We conclude the proof of (iv) Proposition 3.1 using

**Proposition 3.9.** *The forms  $\omega_{F_1}$  of (3.4) and  $\omega$  of Proposition 3.1 satisfy*

$$\delta_g \omega = 0, \quad \delta_g \omega_{F_1} = O_t(x^{n-2k+2}) + O_n(x^{n-2k+4}).$$

*Proof:* Let  $\omega$  be the exact solution of  $\Delta_k \omega = 0$ ,  $\omega|_{x=0} = \omega_0$  in Proposition 3.1. Since  $\delta_g \Delta_k = \Delta_{k-1} \delta_g$ , we deduce that  $\omega' := \delta_g \omega$  is solution of  $\Delta_k \omega' = 0$  with  $\omega'|_{x=0} = 0$  and moreover it is polyhomogeneous since  $\omega$  is polyhomogeneous, so Proposition 3.1 and Lemma 3.8 imply that  $\delta_g \omega \in \ker_{L^2}(\Delta_{k-1})$  and thus  $\delta_g \omega = O(x^{n-2k+3})$  by Corollary

<sup>3</sup>Notice that the result of Mazzeo is stated for 0-elliptic operators with smooth coefficients and acting on functions, but it is straightforward to check that it applies on bundles and with polyhomogeneous coefficients like this is the case for even-dimensionnal Poincaré-Einstein manifolds.

3.3. Since an  $L^2$  harmonic form is closed, then  $\delta_g\omega$  is closed and integration by parts on  $\{x \geq \epsilon\}$  shows, by letting  $\epsilon \rightarrow 0$  in

$$\int_{x \leq \epsilon} |\delta_g\omega|^2 \text{dvol}_g = - \int_{x=\epsilon} \langle \iota_{x\partial_x}\omega, \delta_g\omega \rangle \text{dvol}_g = O(\epsilon)$$

that  $\langle \delta_g\omega, \delta_g\omega \rangle = 0$ . The part with  $\omega_{F_1}$  is also based on  $\delta_g\Delta_k = \Delta_{k-1}\delta_g$  and the uniqueness of the solution of (3.4) up to  $O_t(x^{n-2k+2}) + O_n(x^{n-2k+4})$  on  $(k-1)$ -forms.  $\square$

3.1.2. *The case  $k = \frac{n}{2}$ .* In this case one only intends to solve the equation  $\Delta_k\omega = O(x^\infty)$ , say in the set of almost bounded forms (log  $x$  times bounded). The indicial equation tells us that 0 is a double indicial root for the  $\Lambda_t^k$  part, while 0, 2 are the two simple roots for the  $\Lambda_n$  part. By a double root, we mean a root  $\lambda = \lambda_0$  of order 2 of one of the eigenvalues of  $P_\lambda$  in (2.10). In this case, a straightforward inspection shows that an additional power of  $\log(x)$  must come in the formal expansion of solutions. Since the discussion of this case is not fundamental in our analysis, we prefer to give the result without details. For  $\omega_0, \omega_1 \in \Lambda^{\frac{n}{2}}(M)$  and  $\omega_2 \in \Lambda^{\frac{n}{2}-1}(M)$  one can construct, using (2.11) and (2.12), a polyhomogeneous form

$$\begin{aligned} \omega_F &= \omega_2 \wedge \frac{dx}{x} + \omega_1 \log(x) + \omega_0 \\ &+ x^2 \log x \left( -(\log x)^2 \frac{\delta_0 d\omega_2}{3} + \log x \left( \frac{\delta_0 d\omega_2}{2} + \frac{(-1)^{\frac{n}{2}+1}}{2} \delta_0 \omega_1 \right) + (-1)^{\frac{n}{2}+1} \delta_0 \omega_0 \right. \\ &\left. + (-1)^{\frac{n}{2}} \frac{1}{2} \delta_0 \omega_1 - \frac{1}{2} d\delta_0 \omega_2 - A(\omega_2) \right) \wedge \frac{dx}{x} + O_t(x^2(\log x)^2) + O_n(x^2) \end{aligned}$$

such that  $\Delta_k\omega_F = O(x^\infty)$  and it is unique modulo  $O(x^\infty)$  if the order  $x^2$  coefficient in the  $\Lambda_n$  component is assumed to be 0.

3.1.3. *The case  $k = (n+1)/2$ .* The indicial equation tells us that  $-1, 0$  are the roots for the  $\Lambda_t^k$  part, while 0, 1 are the roots for the  $\Lambda_n$  part. For  $\omega_0 \in \Lambda^k(M)$  and  $\omega_1, \omega_2 \in \Lambda^{k-1}(M)$ , one can construct, using (2.11) and (2.12), a polyhomogeneous form such that

$$\omega_F = \omega_1 \wedge \frac{dx}{x} + \omega_0 + x \omega_2 \wedge \frac{dx}{x} + O(x^2) \text{ and } \Delta_k\omega_F = O(x^\infty)$$

and it is unique modulo  $O(x^\infty)$ . So if  $\omega$  is a solution of Problem 3.1 with  $k = \frac{n-1}{2}$  and boundary value  $\omega_0$  closed, then  $d\omega = O(x^2)$  and  $\Delta_{\frac{n+1}{2}}d\omega = 0$  so  $d\omega = O(x^\infty)$ . But the unique continuation theorem of Mazzeo [17] implies that  $d\omega = 0$ .

We also recall a result proved by Yeganefar [21, Corollary 3.10].

**Proposition 3.10.** *For an odd dimensional Poincaré-Einstein manifold  $(X^{n+1}, g)$ , there is an isomorphism between  $\ker_{L^2}(\Delta_{\frac{n}{2}})$  and  $H^{\frac{n}{2}}(\bar{X}, \partial\bar{X})$  and between  $\ker_{L^2}(\Delta_{\frac{n}{2}+1})$  and  $H^{\frac{n}{2}+1}(\bar{X})$ .*

## 3.2. Relative boundary condition.

**Proposition 3.11.** *Let  $0 < k < \frac{n}{2}$ ,  $x$  be a geodesic boundary defining function and  $\omega_0 \in C^\infty(M, \Lambda^{k-1}(M))$  be a closed form. Then there exists a unique, modulo  $\ker_{L^2}(\Delta_k)$ , form  $\omega$  such that, for all  $\alpha \in [0, 1)$ ,*

$$(3.10) \quad \begin{cases} \omega \in C^{n-2k, \alpha}(\bar{X}, \Lambda_b^k(\bar{X})), \\ \Delta_k\omega = 0 \text{ on } X, \\ \omega|_M = 0, \lim_{x \rightarrow 0} i_{x\partial_x}\omega = \omega_0. \end{cases}$$

Moreover  $\omega$  is closed, smooth on  $\bar{X}$  when  $n$  is odd, while it is polyhomogeneous when  $n$  is even with an expansion at order  $O(x^{n-1} \log x)$  of the form

$$(3.11) \quad \begin{aligned} \omega = & \left( \sum_{j=0}^{n-1} x^j \omega_j^{(n)} \wedge \frac{dx}{x} + \sum_{j=1}^{n-2} x^j \omega_j^{(t)} \right) \\ & + \log x \left( \sum_{j=n-2k+2}^{n-1} x^j \omega_{j,1}^{(n)} \wedge \frac{dx}{x} + \sum_{j=n-2k+2}^{n-2} x^j \omega_{j,1}^{(t)} \right) + O(x^{n-1} \log x) \end{aligned}$$

for some forms  $\omega_j^{(\cdot)}, \omega_{j,1}^{(\cdot)}$  on  $M$ .

*Proof:* the proof is similar to that of Proposition 3.1, so we do not give the full details but we shall use the same notations. We search a formal solution  $\omega'_\infty$  of  $\Delta_k \omega'_\infty = 0$  with  $\omega'_\infty = \omega_0 \wedge \frac{dx}{x} + O(x)$ . Using the indicial equations in Subsection 2.3 and the form of  $\Delta_k$  in Lemma 2.1, we can construct the exponents in the formal series as long as the exponent is not a solution of the indicial equation. Since  $d\omega_0 = 0$  by assumption, we have

$$\Delta_k \left( \omega_0 \wedge \frac{dx}{x} \right) = 2(-1)^{k+1} d\omega_0 + O(x^2) = O(x^2)$$

and so we can continue the construction of  $\omega'_\infty$  until the power  $x^{n-2k}$  in the tangential part  $\Lambda_t^k$  and  $x^{n-2k+2}$  in the  $\Lambda_n^k$  part. At that point, since  $x^{n-2k}$  and  $x^{n-2k+2}$  are solutions of the indicial equation of  $\Delta_k$  in respectively the  $\Lambda_t^k$  and  $\Lambda_n^k$  components, there is a  $x^{n-2k} \log x$  term to include in the  $\Lambda_t^k$  part. Using in addition that  $\Delta_k$  begins with a sum of even powers of  $x$ , we see like in Proposition 3.1 that when  $n$  is odd, a formal series  $\omega'_\infty$  with no log terms can be constructed to solve  $\Delta_k \omega'_\infty = O(x^\infty)$ , while when  $n$  is even we can first construct

$$(3.12) \quad \omega'_{F_2} = \sum_{2j=0}^{n-2k} x^{2j} \omega_{2j}^{(n)} \wedge \frac{dx}{x} + \sum_{2j=2}^{n-2k-2} x^{2j} \omega_{2j}^{(t)} + x^{n-2k} \log(x) \omega_{n-2k,1}^{(t)}$$

with  $\omega_0^{(n)} = \omega_0$  so that  $\Delta_k \omega'_{F_2} = O(x^{n-2k+2} \log x)$ , and the coefficients are uniquely determined by  $\omega_0$ . First observe that  $d\omega'_{F_2} = O(x^2)$  satisfies  $\Delta_{k+1} d\omega'_{F_2} = O(x^{n-2k+2} \log x)$  and since the indicial root in  $[1, n-2k]$  for  $\Delta_{k+1}$  are  $n-2k-2$  in the  $\Lambda_t^{k+1}$  part and  $n-2k$  in the  $\Lambda_n^{k+1}$  part, we deduce that  $d\omega'_{F_2} = O_t(x^{n-2k-2}) + O_n(x^{n-2k})$  and so

$$(3.13) \quad \begin{aligned} d\omega'_{F_2} = & \sum_{2j=2}^{n-2k-2} d\omega_{2j}^{(t)} x^{2j} + \sum_{2j=2}^{n-2k-2} x^{2j} ((-1)^k 2j \omega_{2j}^{(t)} + d\omega_{2j}^{(n)}) \wedge \frac{dx}{x} + x^{n-2k} d\omega_{n-2k}^{(n)} \wedge \frac{dx}{x} \\ & + x^{n-2k} \log(x) (d\omega_{n-2k,1}^{(t)} + (-1)^k (n-2k) \omega_{n-2k,1}^{(t)} \wedge \frac{dx}{x}) + x^{n-2k} (-1)^k \omega_{n-2k,1}^{(t)} \wedge \frac{dx}{x} \\ = & x^{n-2k} \log(x) (d\omega_{n-2k,1}^{(t)} + (-1)^k (n-2k) \omega_{n-2k,1}^{(t)} \wedge \frac{dx}{x}) \\ & + x^{n-2k} (d\omega_{n-2k}^{(n)} + (-1)^k \omega_{n-2k,1}^{(t)}) \wedge \frac{dx}{x}. \end{aligned}$$

Note that we have used that  $(-1)^k (n-2k-2) \omega_{n-2k-2}^{(t)} = -d\omega_{n-2k-2}^{(n)}$ . With these simplifications, we get

$$\Delta_{k+1} d\omega'_{F_2} = (n-2k) x^{n-2k} ((-1)^{k+1} (n-2k) \omega_{n-2k,1}^{(t)} \wedge \frac{dx}{x} + \log(x) d\omega_{n-2k,1}^{(t)}) + O(x^{n-2k+1}).$$

But since  $d\Delta_k \omega'_{F_2} = O(x^{n-2k+2} \log(x))$ , we infer that  $\omega_{n-2k,1}^{(t)}$  must vanish, and we obtain

$$\Delta_k \omega'_{F_2} = O(x^{n-2k+2})$$

Since the order  $x^{n-2k+2}$  is a solution of the indicial equation in the normal part  $\Lambda_n^k$ , we need to add a  $x^{n-2k+2} \log(x)$  normal term to continue the construction of the formal

solution. Since all the subsequent orders are not solution of the indicial equation for  $\Delta_k$ , we can construct, using Borel lemma, a polyhomogeneous  $k$ -form on  $X$  with expansion to order  $x^{n-1} \log(x)$  of the form given by (3.11), which coincides with  $\omega_{F_2}$  at order  $O_n(x^{n-2k+2} \log x) + O_t(x^{n-2k})$ . To obtain an exact solution of (3.10), we can correct  $\omega'_\infty$  by setting  $\omega = \omega'_\infty - E\Delta_k \omega'_\infty$  where  $E$  is defined in Proposition 3.2.

The argument for the uniqueness modulo  $\ker_{L^2} \Delta_g$  is similar to that used in the proof of Proposition 3.1.

To prove that  $\omega$  is closed, it suffices to observe that  $\omega = \omega_{F_2} + O_n(x^{n-2k+2} \log x) + O_t(x^{n-2k})$  and so  $d\omega = x^{n-2k} d\omega_{n-2k}^{(n)} \wedge \frac{dx}{x} + O(x^{n-2k})$ . By remark 3.6, we have  $d\omega \in \ker_{L^2}(\Delta_{k+1})$ . Then  $\delta_g d\omega = 0$  and, considering the decay of  $d\omega$  and  $\omega$  at the boundary, we see by integration by parts that  $d\omega = 0$ .  $\square$

**Remarks:** it is important to remark that the solution  $\omega$  of the problem (3.10) depends on  $\omega_0$  but also on the choice of  $x$ . Note also that the form  $\omega$  solution of (3.10) satisfies  $x\omega \in C^{n-2k+1, \alpha}(\bar{X}, \Lambda^k(\bar{X}))$  for all  $\alpha \in (0, 1)$ .

#### 4. $L_k, G_k$ AND $Q_k$ OPERATORS

In this section we suppose that  $M$  has an even dimension  $n$ .

**4.1. Definitions.** The operators  $L_k, G_k$  derive from the solution of the absolute Dirichlet problem:

**Definition 4.1.** For  $k < \frac{n}{2}$ , the operators  $L_k : C^\infty(M, \Lambda^k(M)) \rightarrow C^\infty(M, \Lambda^k(M))$  and  $G_k : C^\infty(M, \Lambda^k(M)) \rightarrow C^\infty(M, \Lambda^{k-1}(M))$  are defined by  $L_k \omega_0 := \omega_{n-2k,1}^{(t)}$  and  $G_k \omega_0 := \omega_{n-2k+2,1}^{(n)}$  where  $\omega_{n-2k,1}^{(t)}, \omega_{n-2k+2,1}^{(n)}$  are given in the expansion (3.2). When  $k = \frac{n}{2}$ , we define  $G_{\frac{n}{2}} := (-1)^{\frac{n}{2}+1} \delta_0$ .

The operator  $Q_k$  derives from the solution of the relative Dirichlet problem:

**Definition 4.2.** Let  $n$  be even and  $k < \frac{n}{2}$ , the operator  $Q_{k-1} : (C^\infty(M, \Lambda^{k-1}(M)) \cap \ker d) \rightarrow C^\infty(M, \Lambda^{k-1}(M))$  is defined by  $Q_{k-1} \omega_0 := \omega_{n-2k+2,1}^{(n)}$  where  $\omega_{n-2k+2,1}^{(n)}$  is given in the expansion (3.11).

By Corollary 3.3,  $L_k, G_k$  and  $Q_k$  do not depend on the choice of the solution  $\omega$  in Propositions 3.1 or 3.11, though  $L_k$  depends only on the boundary  $(M, [h_0])$ , the operators  $G_k$  and  $Q_k$  may well depend on the whole manifold  $(X, g)$  and not only on the conformal boundary. We will see that they actually depend only on  $(M, [h_0])$  and that they are differential operators.

**4.2. A formal construction.** We show that the definition of  $L_k, G_k, Q_k$  can be done using only the formal series solutions. Let us first define

**Definition 4.3.** For  $k < \frac{n}{2}$ , the operators  $B_k, C_k : C^\infty(M, \Lambda^k(M)) \rightarrow C^\infty(M, \Lambda^{k-1}(M))$  and  $D_k : C^\infty(M, \Lambda^k(M)) \cap \ker d \rightarrow C^\infty(M, \Lambda^k(M))$  are defined by

$$(4.1) \quad \begin{aligned} B_k \omega_0 &:= \left( x^{-n+2k-2} i_{x\partial_x} \Delta_k \omega_{F_1} \right) |_{x=0}, \\ C_k \omega_0 &:= \left( x^{-n+2k-2} i_{x\partial_x} \left( \frac{dx}{x} \wedge \delta_g \omega_{F_1} \right) \right) |_{x=0} \\ D_k \omega_0 &:= \left( x^{-n+2k} i_{x\partial_x} d\omega_{F_1} \right) |_{x=0} \end{aligned}$$

where  $\omega_{F_1}$  solves (3.4).

**Remark:** from the indicial equations and Lemma 3.4,  $B_k\omega_0$  is  $(-1)^k(n-2k+2)$  times the  $x^{n-2k+2}\log(x)$  coefficient in the  $\Lambda_n^k$  part of  $\omega'_\infty$  defined in Proposition 3.5 when  $v^{(t)} = 0$ , this is a differential operator on  $M$  of order  $n-2k+1$  since by construction,  $\omega_{F_1}$  contains only derivatives of order at most  $n-2k-1$  with respect to  $\omega_0$ . The operator  $C_k$  is well defined thanks to Proposition 3.9, and it is a differential operator of order  $n-2k$ . As they come from the expansion of  $\Delta_k, \delta_g$ , they are natural differential operators depending only on  $h_0$  and the covariant derivatives of its curvature tensor.

4.2.1. *The case of  $L_k$ .* It is clear from the proof of Proposition 3.1 that  $L_k\omega_0$  is also the coefficient of the  $x^{n-2k}\log x$  term in the expansion of  $\omega_{F_2}$  defined in (3.6) and of the formal solution  $\omega_\infty$  defined in Proposition 3.8. The indicial equation shows that

$$(4.2) \quad L_k\omega_0 := \frac{1}{n-2k} \left( x^{2k-n} i_{x\partial_x} \left( \frac{dx}{x} \wedge \Delta_k\omega_{F_1} \right) \right) \Big|_{x=0}$$

where  $\omega_{F_1}$  solves (3.4).

4.2.2. *The case of  $G_k$ .* Let us return to the construction of the formal series solution in the proof of Proposition 3.1. Now let  $\omega_{F_2}$  defined in (3.6) and

$$\omega_{F_2} := \omega_{F_1} + x^{n-2k}v^{(t)} + x^{n-2k}\log(x)\omega_{n-2k,1}^{(t)}$$

where  $v^{(t)} \in C^\infty(M, \Lambda^k(M))$  is an arbitrary form. By construction of  $\omega_{F_1}, \omega_{F_2}$ , the fact that  $n-2k$  is an indicial root in the  $\Lambda_t^k$  component and Lemma 2.1, we have

$$\Delta_k\omega_{F_2} = (-1)^{k+1}x^{n-2k+2}(B_k\omega_0 + 2\delta_0v^{(t)}) \wedge \frac{dx}{x} + O_t(x^{n-2k+2}\log x) + O_n(x^{n-2k+4}\log x)$$

to solve away the  $x^{n-2k-2}$  term in  $\Lambda_n^k$  we need to define

$$(4.3) \quad \omega_{F_3} := \omega_{F_2} + \frac{(-1)^{k+1}}{n+2-2k} x^{n-2k}\log(x)(B_k\omega_0 + 2\delta_0v^{(t)}) \wedge \frac{dx}{x}$$

so that  $\Delta_k\omega_{F_3} = O_n(x^{n-2k+4}\log(x)) + O_t(x^{n-2k+2}\log(x))$ . Since  $v^{(t)}$  can be chosen arbitrarily, the coefficient of  $x^{n-2k+4}\log(x)$  in the  $\Lambda_n^k$  component of the formal solution  $\omega_{F_3}$  does not determine a natural operator in term of the initial data  $\omega_0$ , contrary to the  $x^{n-2k}\log(x)$  coefficient in  $\Lambda_t^k$ . In the definition of  $G_k$  above, we used an exact solution on  $X$  to fix the  $v^{(t)}$  term through the Green function, which a priori makes  $G_k$  depend on  $(X, g)$  and not only on  $(M, [h_0])$ . However there is an equivalent way of fixing  $\delta_0v^{(t)}$  without solving a global Dirichlet problem but by adding an additional condition:

**Proposition 4.4.** *Let  $\omega_0 \in C^\infty(M, \Lambda^k(M))$ , then there is a polyhomogeneous  $k$ -form  $\omega_F$  such that*

$$(4.4) \quad \begin{cases} \Delta_k\omega_F = O_t(x^{n-2k+1}) + O_n(x^{n-2k+3}) \\ \delta_g\omega_F = O(x^{n-2k+3}) \\ \omega = \omega_0 + O(x) \end{cases} .$$

*It is unique modulo  $O_t(x^{n-2k}) + O_n(x^{n-2k+2})$  and has an expansion of the form*

$$(4.5) \quad \begin{aligned} \omega_F = & \sum_{j=0}^{\frac{n}{2}-k-1} x^{2j}\omega_{2j}^{(t)} + \sum_{j=1}^{\frac{n}{2}-k} x^{2j}\omega_{2j}^{(n)} \wedge \frac{dx}{x} \\ & + x^{n-2k}\log(x) \left( L_k\omega_0 + x^2 \frac{(-1)^{k+1}}{n-2k} (B_k\omega_0 - 2C_k\omega_0) \wedge \frac{dx}{x} \right). \end{aligned}$$

*Proof:* First consider the uniqueness. By the discussion above, the condition on  $\Delta_k\omega_F$  implies that  $\omega_F$  is necessary of the form  $\omega_F = \omega_{F_3}$  defined in (4.3) for some  $v^{(t)}$ . Now we notice that  $\delta_g\omega_{F_3} = O(x^2)$  satisfies in particular  $\Delta_{k-1}\delta_g\omega_{F_3} = \delta_g\Delta_k\omega_{F_3} = O(x^{n-2k+3})$ , and again by the indicial equation this implies that  $\delta_g\omega_{F_3} = O(x^{n-2k+2})$  since the first

positive indicial root for  $\Delta_{k-1}$  is  $n - 2k + 2$ . Using that  $\delta_0 L_k \omega_0 = 0$  and the form of  $\delta_g$  we obtain

$$\delta_g \omega_{F_3} = \delta_g \omega_{F_1} + x^{n-2k+2} \left( \delta_0 v^{(t)} - \frac{1}{n+2-2k} (B_k \omega_0 + 2\delta_0 v^{(t)}) \right) + O(x^{n-2k+3}).$$

By Proposition 3.9,  $\delta_g \omega_{F_1} = O_t(x^{n-2k+2}) + O_n(x^{n-2k+4})$  and from the definition of  $C_k$ , a necessary condition to have  $\delta_g \omega_F = O(x^{n-2k+3})$  is

$$(n-2k)\delta_0 v^{(t)} = B_k \omega_0 - (n-2k+2)C_k \omega_0.$$

Writing now  $\delta_0 v^{(t)}$  in terms of  $B_k, C_k$  in (4.3) proves the uniqueness and the form of the expansion. Now for the existence, one can take the form in Proposition (3.1). Another way, which again is formal, is first to construct a polyhomogeneous  $(k+1)$ -form  $\omega'_F$  such that

$$\begin{cases} \Delta_{k+1} \omega'_F = O_t(x^{n-2k-1}) + O_n(x^{n-2k+1}) \\ \omega'_F = \frac{2(-1)^{k+1}}{n-2k} \log(x) d\omega_0 + \omega_0 \wedge \frac{dx}{x} + O(x), \end{cases}$$

which can be done as in Proposition 3.11 by using the indicial equations, and then to set  $\omega_F := \delta_g \omega'_F$ . It is easy to see that this form is a polyhomogeneous solution of (4.4).  $\square$

Since the exact solution in Proposition 3.1 is coclosed, we deduce from Proposition 4.4 the

**Corollary 4.5.** *The operator  $G_k$  is a natural differential operator of order  $n - 2k + 1$  which is given by*

$$G_k = (-1)^{k+1} \frac{B_k - 2C_k}{n-2k}$$

and depends only on  $h_0$  and the covariant derivatives of its curvature tensor.

**Remark 4.6.** *Problem (4.4) and Corollary 4.5 allow to define  $L_k$  and  $G_k$  on any even manifold  $M^{2m}$  with no need of cobordism assumption (as in Proposition 3.1). We just have to work on  $X = M \times [0, \epsilon[$ .*

4.2.3. *The operator  $Q_k$ .* Following the ideas used above for  $G_k$ , we shall show how to construct  $Q_k$  from a formal solution  $\omega_{F_1}$ . We start by

**Definition 4.7.** *For  $1 \leq k \leq \frac{n}{2}$ , define the operators  $B'_{k-1} : C^\infty(M, \Lambda^{k-1}(M)) \cap \ker d \rightarrow C^\infty(M, \Lambda^{k-1}(M))$  and  $D'_{k-1} : C^\infty(M, \Lambda^{k-1}(M)) \cap \ker d \rightarrow C^\infty(M, \Lambda^k(M))$  by*

$$(4.6) \quad \begin{aligned} B'_{k-1} \omega_0 &:= \left( x^{-n+2k-2} i_{x\partial_x} \Delta_k \omega'_{F_1} \right) |_{x=0}, \\ D'_{k-1} \omega_0 &:= \left( x^{-n+2k} i_{x\partial_x} d\omega'_{F_1} \right) |_{x=0} \end{aligned}$$

where  $\omega'_{F_1}$  is the form in (3.12) such that  $\Delta_k \omega'_{F_1} = O(x^{n-2k+2})$  and  $\omega'_{F_1} = \omega_0 \wedge \frac{dx}{x} + O(x^2)$ .

Let us now set  $\omega'_{F_2} := \omega'_{F_1} + v^{(t)} x^{n-2k}$  for some arbitrary smooth form  $v^{(t)}$  on  $M$ , we obtain

$$\Delta_k \omega'_{F_2} = (-1)^{k+1} x^{n-2k+2} (B'_{k-1} \omega_0 + 2\delta_0 v^{(t)}) \wedge \frac{dx}{x} + O_t(x^{n-2k+2}) + O_n(x^{n-2k+3}).$$

so to solve away the  $x^{n-2k+2}$  normal coefficient, we need to define

$$(4.7) \quad \omega'_{F_3} := \omega'_{F_2} + \frac{(-1)^{k+1}}{n-2k+2} x^{n-2k+2} \log(x) \left( B'_{k-1} \omega_0 + 2\delta_0 v^{(t)} \right) \wedge \frac{dx}{x}$$

which satisfies  $\Delta_k \omega'_{F_3} = O_t(x^{n-2k+2} \log(x)) + O_n(x^{n-2k+3})$ . Like for  $G_k$ , the term  $v^{(t)}$  is arbitrary and so we have to impose an additional condition to fix this term (or at least to fix  $\delta_0 v^{(t)}$ ).

**Proposition 4.8.** *Let  $\omega_0 \in C^\infty(M, \Lambda^{k-1}(M))$  be closed, then there is a polyhomogeneous  $k$ -form  $\omega'_F$  which satisfies*

$$(4.8) \quad \begin{cases} \Delta_k \omega'_F = O_t(x^{n-2k+1}) + O_n(x^{n-2k+3}) \\ d\omega'_F = O(x^{n-2k+1}) \\ \omega'_F = \omega_0 \wedge \frac{dx}{x} + O(x^2) \end{cases},$$

which is unique modulo  $O_t(x^{n-2k}) + O_n(x^{n-2k+2})$  and has an expansion of the form

$$(4.9) \quad \begin{aligned} \omega'_F &= \sum_{j=1}^{\frac{n}{2}-k-1} x^{2j} \omega_{2j}^{(t)} + \sum_{j=0}^{\frac{n}{2}-k} x^{2j} \omega_{2j}^{(n)} \wedge \frac{dx}{x} - x^{n-2k} \frac{1}{n-2k} D'_{k-1} \omega_0 \\ &+ \frac{(-1)^{k+1}}{n-2k+2} x^{n-2k+2} \log(x) \left( B'_{k-1} \omega_0 - \frac{2\delta_0 D'_{k-1} \omega_0}{n-2k} \right) \wedge \frac{dx}{x}. \end{aligned}$$

*Proof:* (i) Take  $\omega'_F = \omega'_{F_3}$  defined in (4.7), then  $\Delta_k \omega'_F = O_t(x^{n-2k+1}) + O_n(x^{n-2k+3})$  by construction. Moreover, since  $\omega_0$  is closed, one has  $d\omega'_F = O(x^2)$  and  $\Delta_{k+1} d\omega'_F = O(x^{n-2k+1})$ . Since the indicial roots for  $\Delta_{k+1}$  in  $[2, n-2k+1]$  are  $n-2k-2$  in the  $\Lambda_t^{k+1}$  part and  $n-2k$  in the  $\Lambda_n^{k+1}$  part, this implies that  $d\omega'_F = O_t(x^{n-2k-2}) + O_n(x^{n-2k})$ . Then, using (3.13), we obtain

$$\begin{aligned} d\omega'_F &= x^{n-2k} \left( dv^{(t)} + ((-1)^k (n-2k) v^{(t)} + d\omega_{n-2k}^{(n)}) \wedge \frac{dx}{x} \right) + O(x^{n-2k+1}) \\ &= x^{n-2k} \left( dv^{(t)} + ((-1)^k (n-2k) v^{(t)} + (-1)^k D'_{k-1} \omega_0 \right) \wedge \frac{dx}{x} + O(x^{n-2k+1}). \end{aligned}$$

So  $d\omega'_F = O(x^{n-2k+1})$  if and only if  $v^{(t)} = -D'_{k-1} \omega_0 / (n-2k)$ .  $\square$

The first corollary is

**Corollary 4.9.** *For  $k < \frac{n}{2}$ , the operator  $Q_k$  is a natural differential operator of order  $n-2k$  which is given by*

$$Q_k = \frac{(-1)^k}{n-2k} \left( B'_k - \frac{\delta_0 D'_k}{\frac{n}{2} - k - 1} \right)$$

and it depends only on  $h_0$  and the covariant derivatives of its curvature tensor.

**Remark 4.10.** *Here also this corollary allows to define the operator  $Q_k$  on any even manifold by considering Problem (4.8).*

As a corollary of Propositions 4.4 and 4.8, we also have

**Corollary 4.11.** *If  $\omega_0$  is a closed  $k$ -form on  $M$ , then there is a polyhomogeneous  $k$ -form  $\omega_F$  on  $\bar{X}$  such that*

$$(4.10) \quad \begin{cases} d\omega_F = O(x^{n-2k+1}) \\ \delta_g \omega_F = O(x^{n-2k+3}) \\ \omega_F = \omega_0 + O(x) \end{cases}$$

It is unique modulo  $O_t(x^{n-2k+1}) + O_n(x^{n-2k+2})$  and it has an expansion

$$\begin{aligned} \omega_F &= \sum_{j=0}^{\frac{n}{2}-k-1} x^{2j} \omega_{2j}^{(t)} + \sum_{j=1}^{\frac{n}{2}-k} x^{2j} \omega_{2j}^{(n)} \wedge \frac{dx}{x} - \frac{1}{n-2k} D_k \omega_0 x^{n-2k} \\ &+ x^{n-2k+2} \log(x) \left( \frac{(-1)^{k+1}}{n-2k} (B_k \omega_0 - 2C_k \omega_0) \wedge \frac{dx}{x} \right). \end{aligned}$$

*Proof:* for the existence, take  $\omega'_F$  in Proposition 4.8 ( $\omega'_F$  is  $k+1$  form now since  $\omega_0 \in \Lambda^k(M)$ ) and consider  $\omega_F := (-1)^{k+1} / (2k-n) \delta_g \omega'_F$ . It is easy to see that  $\omega_F = \omega_0 + O(x^2)$  and that  $\Delta_k \omega_F = O_t(x^{n-2k+1}) + O_n(x^{n-2k+3})$ . Since  $d\delta_g \omega'_F = -\delta_g d\omega'_F +$

$O_t(x^{n-2k-1}) + O_n(x^{n-2k+1})$ , we deduce that  $d\omega_F = O_t(x^{n-2k-1}) + O_n(x^{n-2k+1})$ . But from the Proposition 4.4,  $\omega_F = \omega_{F_1} + v^{(t)}x^{n-2k} + O(x^{n-2k+1})$  (note that  $L_k\omega_0 = 0$  by Proposition 4.12) for some  $k$ -form  $v^{(t)}$  on  $M$  and so we conclude that

$$\begin{aligned} d\omega_F &= \sum_{2j=2}^{n-2k-2} d\omega_{2j}^{(t)} x^{2j} + \sum_{2j=2}^{n-2k-2} x^{2j} ((-1)^k 2j \omega_{2j}^{(t)} + d\omega_{2j}^{(n)}) \wedge \frac{dx}{x} \\ &\quad + x^{n-2k} (d\omega_{n-2k}^{(n)} + (-1)^k (n-2k)v^{(t)}) \wedge \frac{dx}{x} + x^{n-2k} dv^{(t)} + O(x^{n-2k+1}) \\ &= O(x^{n-2k+1}) \end{aligned}$$

so  $v^{(t)}$  has to be  $(-1)^{k+1} d\omega_{n-2k}^{(n)} / (n-2k)$  to get  $d\omega_F = O_t(x^{n-2k-1}) + O_n(x^{n-2k+1})$ . But clearly this argument also implies that  $d\omega_{F_1} = x^{n-2k} d\omega_{n-2k}^{(n)} \wedge \frac{dx}{x}$  and the expansion of  $\omega_F$  is then a consequence of this fact together with the expansion (4.5) in Proposition 4.4 and the definition of  $D_k$ .  $\square$

**Remark:** in Proposition 4.4, 4.8 and Corollary 4.11, we do not really need to take  $\omega_0 \in C^\infty(M, \Lambda(M))$ . Indeed, for an  $\omega_0$  in  $L^2(\Lambda(M))$ , the arguments would work in a similar fashion except that the expansion in power of  $x$  and  $\log(x)$  have coefficients in some  $H^{-N}(\Lambda(M))$  with  $N$  large enough, as we discussed in the proof of Proposition 3.1.

### 4.3. Factorizations.

**Proposition 4.12.** *For any  $k < \frac{n}{2} - 1$ , the following identities hold*

$$(4.11) \quad \begin{aligned} G_k &= (-1)^k \frac{\delta_{h_0} Q_k}{n-2k} \quad \text{on closed forms,} \\ L_k &= \frac{(-1)^k}{(n-2k)} G_{k+1} d = -\frac{\delta_{h_0} Q_{k+1} d}{(n-2k)(n-2k-2)}. \end{aligned}$$

while for  $k = \frac{n}{2} - 1$

$$(4.12) \quad L_{\frac{n}{2}-1} = \frac{1}{2} \delta_{h_0} d.$$

*Proof:* Let  $\omega$  be a solution of Problem (4.8) with initial data  $\omega_0$  closed. Then its first log term is  $x^{n-2k+2} \log(x) Q_{k-1} \omega_0 \wedge \frac{dx}{x}$  and thus the first normal log term of  $\delta_g \omega$  is  $x^{n-2k+4} \log(x) (\delta_0 Q_{k-1} \omega_0) \wedge \frac{dx}{x}$ . But  $\delta_g \omega$  is a solution of Problem (4.4) with boundary term  $\delta_g \omega = (-1)^k (2k - n - 2) \omega_0 + O(x)$ . Thus, the form  $\delta_g \omega$  has for first normal log term  $(-1)^k (2k - n - 2) x^{n-2k+4} (G_{k-1} \omega_0) \wedge \frac{dx}{x}$ .

Let  $\omega$  be a solution of 3.1 with initial data  $\omega_0$ . Since  $\Delta_{k+1} d = d\Delta_k$  and  $\delta_g d\omega = \Delta_{k+1} d\omega - d\delta_g \omega$ , the form  $\omega' := d\omega$  is a solution Problem (4.4) with initial data  $d\omega_0$  and first log term  $(-1)^k (n-2k) x^{n-2k} \log(x) L_k \omega_0 \wedge \frac{dx}{x}$ , which gives (4.11).

To compute  $L_{\frac{n}{2}-1}$  we use Equation (4.2). Using Relations (2.9), we get  $\omega_{F_1} = \omega_0 - x \frac{(-1)^{\frac{n}{2}}}{2} \delta_0 \omega_0 \wedge dx$ , therefore  $\Delta_{\frac{n}{2}-1} \omega_{F_1} = x^2 \delta_0 d\omega_0 + o(x^2)$ .  $\square$

**Remark:** Note that this implies that  $L_k$  is zero on closed forms and  $G_k$  has its range in co-closed forms.

**4.4. Conformal properties.** A priori our construction of  $L_k, G_k, Q_k$  depends on the choice of geodesic boundary defining function  $x$ , i.e. on the choice of conformal representative in  $[h_0]$ . In order to study the conformal properties of these operators, we need to compare the splittings of the differential forms associated to different conformal representatives.

A system of coordinates  $y = (y_i)_{i=1, \dots, n}$  on  $M$  near a point  $p \in M$  give rise to a system of coordinates  $(x, y)$  in  $\bar{X}$  near the boundary point  $p$  through the diffeomorphism  $\psi : (x, y) \rightarrow \psi_x(y)$  where  $\psi_t$  is the flow of the gradient  $\nabla^{x^2 g} x$  of  $x$  with respect to  $x^2 g$ . Such a system  $(x, y)$  is called a system of *geodesic normal coordinates associated to  $h_0$* .

**Lemma 4.13.** *Let  $(x, y)$  and  $(\hat{x}, \hat{y})$  be two systems of geodesic normal coordinates associated respectively to  $h_0$  and  $\hat{h}_0 = e^{2\varphi_0} h_0$ . If  $\hat{\omega}$  (resp.  $\hat{\omega} \wedge d\hat{x}$ ) is a  $k$ -form tangential (resp. normal) in the coordinates  $(\hat{x}, \hat{y})$  with  $\hat{\omega}|_{\hat{x}=0} = \omega_0$ , then we have*

$$\begin{aligned}\hat{\omega} &= \omega_0 + (-1)^{k+1} x^2 (i_{\nabla\varphi_0} \omega_0) \wedge \frac{dx}{x} + O_t(x^2) + O_n(x^3), \\ \hat{\omega} \wedge \frac{d\hat{x}}{\hat{x}} &= \omega_0 \wedge \frac{dx}{x} + \omega_0 \wedge d\varphi_0 + O_t(x) + O_n(x^2).\end{aligned}$$

*Proof:* By the proof of Lemma 2.1 in [13], if  $\hat{h}_0 = e^{2\varphi_0} h_0$  is another conformal representative, a geodesic boundary defining function  $\hat{x}$  associated to  $\hat{h}_0$  satisfies  $\hat{x} = e^\varphi x$  with  $\varphi = \varphi_0 + O(x^2)$  at least  $C^{n-1}$  and  $\hat{y}_i(x, y) = y_i + \frac{x^2}{2} \sum_{j=1}^n h^{ij} \partial_{y_j} \varphi_0 + O(x^3)$ . Hence  $d\hat{y}_i = dy_i + x \sum_j h^{ij} \partial_{y_j} \varphi_0 dx$  and  $d\hat{x} = x e^{\varphi_0} d\varphi_0 + e^{\varphi_0} dx + O(x^2)$ , which gives the relations above.  $\square$

This implies the following corollary:

**Corollary 4.14.** *Under a conformal change  $\hat{h}_0 = e^{2\varphi_0} h_0$ , the associated operators  $\hat{L}_k$ ,  $\hat{H}_k$  and  $\hat{Q}_k$  are given by*

$$(4.13) \quad \begin{aligned}\hat{L}_k &= e^{(2k-n)\varphi_0} L_k, \quad \hat{G}_k = e^{(2k-2-n)\varphi_0} (G_k + (-1)^k i_{\nabla\varphi_0} L_k) \\ \hat{Q}_k \omega_0 &= e^{\varphi_0(2k-n)} \left( Q_k \omega_0 + (n-2k) L_k(\varphi_0 \omega_0) \right)\end{aligned}$$

where  $\omega_0 \in C^\infty(M, \Lambda^k(M))$  is any closed form. Thus  $L_k$  is conformally covariant and  $G_k$  is conformally covariant on the kernel of  $L_k$  (hence on closed forms).

*Proof:* Let  $\omega$  be a solution of Problem (4.4) with respect to  $(x, y)$  a system associated to  $h_0$ . Then by Lemma 4.13,  $\omega$  is also a solution of Problem (4.4) with respect to  $(\hat{x}, \hat{y})$  a system associated to  $\hat{h}_0$ . Now, when we change  $x$  to  $\hat{x}$  the first  $\log x$  term (i.e. the  $x^{n-2k} \log x$  term) in the expansion of  $\omega$  changes by a multiplication by  $e^{(2k-n)\varphi_0}$ . As for the  $x^{n-2k+2} \log x$  term in the normal part, we have a similar effect but the tangential  $x^{n-2k} \log x$  term gives rise to a  $\hat{x}^{n-2k+2} \log \hat{x}$  normal term which gives the term  $i_{\nabla\varphi_0} L_k$ .

Let  $\omega$  be a solution of Problem (4.8) in the variable  $x$  with initial data  $\omega_0 \wedge \frac{dx}{x}$ ,  $\hat{\omega}_1$  be a solution of Problem (4.8) in the variable  $\hat{x}$  with initial data  $\omega_0 \wedge \frac{d\hat{x}}{\hat{x}}$  and  $\hat{\omega}_2$  be a solution of Problem (4.10) in the variable  $\hat{x}$  with initial data  $-\omega_0 \wedge d\varphi_0$ . Using Lemma 4.13 get that  $\hat{\omega}_1 + \hat{\omega}_2$  satisfies Problem 4.8 in the variable  $x$  with initial data  $\omega_0$ . So  $\omega = \hat{\omega}_1 + \hat{\omega}_2$  modulo  $O_t(x^{n-2k}) + O_n(x^{n-2k+2})$  and the  $x^{n-2k+2} \log x$  normal terms must be the same. Using 4.13, we get

$$Q_{k-1} \omega_0 \wedge \frac{dx}{x} = e^{\varphi_0(n-2k+2)} \left( \hat{Q}_{k-1} \omega_0 - \hat{G}_k(\omega_0 \wedge d\varphi_0) \right) \wedge \frac{dx}{x}.$$

Now we use the transformation formula of  $\hat{G}_k$  and (4.11) with  $d\omega_0 = 0$  to see that

$$\begin{aligned}e^{\varphi_0(n-2k+2)} \hat{G}_k(\omega_0 \wedge d\varphi_0) &= (-1)^{k+1} G_k d(\varphi_0 \omega_0) + (-1)^k i_{\nabla\varphi_0} L_k d(\varphi_0 \omega_0) \\ &= (n-2k+2) L_{k-1}(\varphi_0 \omega_0)\end{aligned}$$

This ends the proof of the transformation law of  $Q_{k-1}$  by conformal change.  $\square$

**Remark:** while  $Q_k$  on  $\ker d$  is not conformally invariant (by Proposition 4.14), the pairing  $\langle Q_k u, u \rangle_{L^2(\text{dvol}_{h_0})}$  for the metric  $h_0$  is conformally invariant for  $u \in \ker d$ . Indeed, using

(4.13), a conformal change of metric  $\hat{h}_0 = e^{2\varphi_0} h_0$  gives

$$\int_M \langle \hat{Q}_k u, u \rangle_{\hat{h}_0} d\text{vol}_{h_0} = \int_M \langle Q_k u, u \rangle_{h_0} + \frac{\langle \delta_0 Q_{k+1} d(\varphi_0 u), u \rangle_{h_0}}{2k+2-n} d\text{vol}_{h_0}$$

which by integration by parts and  $du = 0$  gives the  $\langle \hat{Q}_k u, u \rangle_{L^2(\text{dvol}_{\hat{h}_0})} = \langle Q_k u, u \rangle_{L^2(\text{dvol}_{h_0})}$ . Of course, when we restrict this form to exact forms, this is given by

$$\langle Q_k du, du \rangle = \langle L_{k-1} u, u \rangle$$

which is real and conformally invariant.

#### 4.5. Analytical properties.

**Proposition 4.15.** *For any  $k < \frac{n}{2}$  we have*

$$\begin{aligned} Q_k &= \frac{(-1)^{\frac{n}{2}+k+1} (n-2k) (\Delta_0)^{\frac{n}{2}-k}}{2^{n-2k} [(\frac{n}{2}-k)!]^2} + \text{lower order terms in } \partial_{y_i}^j \\ L_k &= \frac{(-1)^{\frac{n}{2}+k+1} (n-2k) (\delta_0 d)^{\frac{n}{2}-k}}{2^{n-2k} [(\frac{n}{2}-k)!]^2} + \text{lower order terms in } \partial_{y_i}^j \\ G_k &= \frac{(-1)^{\frac{n}{2}+1} (\delta_0 d)^{\frac{n}{2}-k} \delta_0}{2^{n-2k} [(\frac{n}{2}-k)!]^2} + \text{lower order terms in } \partial_{y_i}^j \end{aligned}$$

*Proof:* We first review the computation of  $\omega_{F_1}$  which solves (3.4). By Lemma 2.1,  $\omega_{F_1}$  has the form  $\omega_{F_1} = \sum_{i=0}^{\frac{n}{2}-k-1} x^{2i} \omega_{2i}^{(t)} + \sum_{i=1}^{\frac{n}{2}-k} x^{2i} \omega_{2i}^{(n)} \wedge \frac{dx}{x}$ , where the  $\omega_i^{(*)}$  are images of  $\omega_0$  by differential operators on  $M$ . We compute the principal part of these operators by recurrence.

The decomposition (2.7) of  $\Delta_k$  and the identity  $\Delta_k \omega_{F_1} = O_t(x^{n-2k}) + O_n(x^{n-2k+1})$  give

$$\begin{aligned} &\sum_{i=1}^{\frac{n}{2}-k} x^{2i} \left( -4i(k+i-\frac{n}{2}-1) \omega_{2i}^{(n)} + \sum_{j=1}^i S'_j \omega_{2i-2j}^{(t)} + \sum_{j=1}^{i-1} (R'_j + 2(i-j)P'_j) \omega_{2i-2j}^{(n)} \right) \wedge \frac{dx}{x} + \\ &\sum_{i=1}^{\frac{n}{2}-k-1} x^{2i} \left( -4i(k+i-\frac{n}{2}) \omega_{2i}^{(t)} - (-1)^k d \omega_{2i}^{(n)} + \sum_{j=1}^i (R_j + 2(i-j)P_j) \omega_{2i-2j}^{(t)} + \sum_{j=1}^{i-1} S_j \omega_{2i-2j}^{(n)} \right) = 0 \end{aligned}$$

This determines uniquely the  $\omega_i^{(*)}$ .

Let us write LOT for lower order term operators on  $M$ . Then we get

$$\omega_2^{(n)} = \frac{(-1)^{k+1}}{2k-n} \delta_0 \omega_0, \quad \omega_2^{(t)} = \left( \frac{d\delta_0}{2(2k-n)} + \frac{\delta_0 d}{2(2k+2-n)} \right) \omega_0$$

and given the order in  $\partial_{y_i}$  of the  $R_i, R'_i, \bar{R}_i, \bar{R}'_i, Q_i$  and  $Q'_i$ , we have

$$\begin{aligned} \omega_{2i}^{(t)} &= \frac{1}{2i(2k+2i-n)} \left( 2(-1)^{k+1} d \omega_{2i}^{(n)} + \Delta_0 \omega_{2i-2}^{(t)} \right) + \text{LOT}(\omega_0) \\ \omega_{2i+2}^{(n)} &= \frac{1}{2(i+1)(2k+2i-n)} \left( 2(-1)^{k+1} \delta_0 \omega_{2i}^{(t)} + \Delta_0 \omega_{2i}^{(n)} \right) + \text{LOT}(\omega_0) \end{aligned}$$

So we have

$$\begin{aligned} \omega_{2i}^{(t)} &= \left( a_{2i} (\delta_0 d)^i + b_{2i} (d\delta_0)^i + \text{LOT} \right) \omega_0 \\ \omega_{2i+2}^{(n)} &= \left( a_{2i+1} (\delta_0 d)^i \delta_0 + \text{LOT} \right) \omega_0 \end{aligned}$$

where the sequences  $(a_i)$  and  $(b_{2i})$  satisfy the relations

$$a_{2i} = \frac{a_{2i-2}}{2i(2k+2i-n)}, \quad a_{2i+1} = \frac{2(-1)^{k+1}b_{2i}}{2(i+1)(2k+2i-n)} + \frac{a_{2i-1}}{2(i+1)(2k+2i-n)}$$

$$b_{2i} = \frac{2(-1)^{k+1}a_{2i-1}}{2i(2k+2i-n)} + \frac{b_{2i-2}}{2i(2k+2i-n)}$$

and  $a_1 = \frac{(-1)^{k+1}}{2k-n}$ ,  $a_2 = \frac{1}{2(2k+2-n)}$ ,  $b_2 = \frac{1}{2(2k-n)}$ . By uniqueness of the solution of this equation we find

$$a_{2i} = \frac{1}{2^i i! \prod_{j=1}^i (2k+2j-n)}, \quad a_{2i+1} = \frac{(-1)^{k+1}}{2^i i! \prod_{j=0}^i (2k+2j-n)},$$

$$b_{2i} = \frac{1}{2^i i! \prod_{j=0}^{i-1} (2k+2j-n)}$$

for all  $i \leq \frac{n}{2} - k - 1$ . We infer the equality

$$(4.14) \quad \begin{aligned} \Delta_k \omega_{F_1} &= x^{n-2k} \left( a_{n-2k-2} (\delta_0 d)^{\frac{n}{2}-k} + (b_{n-2k-2} + 2(-1)^{k+1} a_{n-2k-1}) (d\delta_0)^{\frac{n}{2}-k} + \text{LOT} \right) \omega_0 \\ &\quad + x^{2k-n+2} \left( a_{n-2k-1} (\delta_0 d)^{\frac{n}{2}-k} \delta_0 + \text{LOT} \right) \omega_0 \wedge \frac{dx}{x} + o(x^{n-2k+1}) \\ &= x^{n-2k} \left( \frac{(\delta_0 d)^{\frac{n}{2}-k}}{2^{\frac{n}{2}-k-1} (\frac{n}{2}-k-1)! \prod_{j=1}^{\frac{n}{2}-k-1} (2k+2j-n)} + \text{LOT} \right) \omega_0 \\ &\quad + x^{2k-n+2} \left( \frac{(-1)^{k+1} (\delta_0 d)^{\frac{n}{2}-k} \delta_0}{2^{\frac{n}{2}-k-1} (\frac{n}{2}-k-1)! \prod_{j=0}^{\frac{n}{2}-k-1} (2k+2j-n)} + \text{LOT} \right) \omega_0 \wedge \frac{dx}{x} \\ &\quad + o(x^{n-2k+1}) \end{aligned}$$

so we have

$$L_k = \frac{-(\delta_0 d)^{\frac{n}{2}-k}}{2^{\frac{n}{2}-k-1} (\frac{n}{2}-k-1)! \prod_{j=0}^{\frac{n}{2}-k-1} (2k+2j-n)} + \text{LOT}$$

$$B_k = \frac{(\delta_0 d)^{\frac{n}{2}-k} \delta_0}{2^{\frac{n}{2}-k-1} (\frac{n}{2}-k-1)! \prod_{j=0}^{\frac{n}{2}-k-1} (2k+2j-n)} + \text{LOT}$$

Note also that  $\delta_g$  is of order 1 so  $C_k$  has no contribution to the principal part of  $G_k$  and we get

$$G_k = \frac{(-1)^{k+1} (\delta_0 d)^{\frac{n}{2}-k} \delta_0}{2^{\frac{n}{2}-k} (\frac{n}{2}-k)! \prod_{j=0}^{\frac{n}{2}-k-1} (2k+2j-n)} + \text{LOT}.$$

The proof is the same (and even easier) for  $Q_k$ . We could have deduced the principal parts of  $L_k$  and  $G_k$  from the one of  $Q_k$ , but a slight generalization of the proof above will allow to compute the principal part of the non-critical  $L_k^l$  in the next section.  $\square$

We finally prove that the operators  $L_k$  and  $Q_k$  are symmetric on  $C^\infty(M, \Lambda(M))$ :

**Proposition 4.16.** *For  $k \leq \frac{n}{2} - 1$ , the operators  $L_k$  are symmetric on  $C^\infty(M, \Lambda^k(M))$  while for  $k < \frac{n}{2} - 1$ , the operators  $Q_k$  are symmetric on  $C^\infty(M, \Lambda^k(M)) \cap \ker d$ .*

*Proof:* The proof for  $L_k$  is done in Proposition 5.4 which covers the non-critical cases. The proof for  $Q_k$  is quite similar, we let  $\omega_0, \omega'_0$  be two closed  $k$ -forms on  $M$  and  $\omega, \omega'$  the forms constructed in the proof of Proposition 3.11 with respective initial conditions  $\omega_0$

and  $\omega'_0$ . Then integration by part and the fact that  $d\omega = d\omega' = 0$  gives

$$0 = \int_{x \geq \epsilon} (\langle \Delta_k \omega, \omega' \rangle_g - \langle \Delta_k \omega', \omega \rangle_g) d\text{vol}_g = \int_{x=\epsilon} (\langle i_{x\partial_x} \omega, \delta_g \omega' \rangle_{h_x} - \langle i_{x\partial_x} \omega', \delta_g \omega \rangle_{h_x}) x^{-n} d\text{vol}_{h_x}.$$

But a straightforward analysis and the fact that  $L_k(\omega_0) = L_k(\omega'_k) = 0$  give that the second line has an expansion of the form

$$a_{-2\ell} \epsilon^{-2\ell} + \dots + a_{-2} \epsilon^{-2} + L \log(\epsilon) + O(1)$$

with  $L := (-1)^{k+1} (2k - n) \left( \langle Q_k \omega_0, \omega'_0 \rangle_{L^2(\text{dvol}_{h_0})} - \langle \omega_0, Q_k \omega'_0 \rangle_{L^2(\text{dvol}_{h_0})} \right)$

This achieves the proof.  $\square$

**4.6. Branson  $Q$ -curvature.** We conclude this section by the observation that  $Q_0$  is the  $Q$ -curvature of Branson.

**Proposition 4.17.** *The operator  $Q_0$  of Definition 4.2 satisfies*

$$Q_0 1 = \frac{n(-1)^{\frac{n}{2}+1}}{2^{n-1} \frac{n!}{2} (\frac{n}{2} - 1)!} Q$$

where  $Q$  is Branson  $Q$ -curvature defined in [2].

*Proof:* Since the operator  $Q_0$  and the function  $Q$  are local on  $(M, [h_0])$  and do not depend on the chosen Poincaré-Einstein manifold with conformal infinity  $(M, [h_0])$ , it suffices to consider the cylinder  $X = (-1, 1) \times M$  equipped with a Poincaré-Einstein metric with conformal metric  $[h_0]$  on the boundary  $M \sqcup M$ . In [5], Fefferman and Graham showed that the  $Q$ -curvature of Branson is the function  $Q$  on  $M$  such that if  $U \in C^\infty(X)$  is a solution of

$$\begin{cases} \Delta_g U = n \\ U = \log(x) + A + x^n B \log(x) \text{ with } A, B \in C^\infty(\bar{X}) \\ A|_{x=0} = 0 \end{cases}$$

then  $B|_{x=0} = (-1)^{\frac{n}{2}+1} (2^{n-1} \frac{n!}{2} (\frac{n}{2} - 1)!)^{-1} Q$ . Consider  $dU$ , clearly it is a harmonic 1-form and it is given by

$$dU = \frac{dx}{x} + dA + nx^n B \log(x) \frac{dx}{x} + O(x^n)$$

and by uniqueness of the solution in Proposition 3.11 and the decay of  $L^2$  harmonic 1-forms (of order  $x^n$ ), we deduce that  $Q_0 1 = nB|_{x=0}$ , this proves the claim (note that the log term in the development of  $\Delta_k$  does not interfere since it acts trivially on normal zero forms).  $\square$

## 5. THE NON-CRITICAL CASE

Let  $(X, g)$  be a Poincaré-Einstein manifold with conformal infinity  $(M, [h_0])$ . We assume  $k \leq (n+1)/2$  and  $n$  may be odd or even in this section, and we let  $\ell$  be an integer in  $[1, \frac{n}{2} - k]$  in general, and  $\ell \in \mathbb{N}$  if  $n$  is odd and  $(X, g)$  is an even Poincaré-Einstein manifold. We want to construct the operators  $L_k^\ell$  of [3] by solving the following equation

$$(5.1) \quad \left( \Delta_k - \left( \frac{n}{2} - k + \ell \right) \left( \frac{n}{2} - k - \ell \right) \right) \omega = O_t(x^{\frac{n}{2}-k+\ell}) + O_n(x^{\frac{n}{2}-k+\ell+1})$$

with  $\omega = x^{\frac{n}{2}-k-\ell} \omega_0 + o(x^{\frac{n}{2}-k-\ell})$  as  $x \rightarrow 0$ .

where  $O_n, O_t$  are defined in the proof of Proposition 3.1 and where  $\omega_0 \in C^\infty(M, \Lambda^k(M))$ . This can be done essentially like in the critical case, using the indicial equations of Subsection 2.3. Indeed, the indicial roots of  $\Delta_k - (\frac{n}{2} - k + \ell)(\frac{n}{2} - k - \ell)$  can be computed

rather easily, these are

$$\begin{aligned} & \frac{n}{2} - k \pm \ell \quad \text{in the } \Lambda_t^k \text{ component} \\ & \frac{n}{2} - k + 1 \pm \sqrt{\ell^2 + n + 1 - 2k} \quad \text{in the } \Lambda_n^k \text{ component.} \end{aligned}$$

Notice that there are no indicial roots in  $(\frac{n}{2} - k - \ell, \frac{n}{2} - k + \ell)$  when  $\ell \leq n/2 - k$ , but there is one root in this interval when  $\ell > n/2 - k$  and  $n$  odd, it is given by

$$(5.2) \quad \frac{n}{2} - k + 1 - \sqrt{\ell^2 + n + 1 - 2k} \in \left(\frac{n}{2} - k - \ell, \frac{n}{2} - k - \ell + 1\right]$$

and thus is not in  $n/2 - k - \ell + 2\mathbb{N}_0$ . We obtain

**Lemma 5.1.** *For  $\omega_0 \in C^\infty(M, \Lambda^k(M))$  fixed, there exists a series*

$$(5.3) \quad \omega_{F_1} = x^{\frac{n}{2} - k - \ell} \left( \sum_{2j=0}^{2l-2} x^{2j} \omega_{2j}^{(t)} + \sum_{2j=2}^{2l} x^{2j} (\omega_{2j}^{(n)} \wedge \frac{dx}{x}) \right)$$

such that  $\omega_0^{(t)} = \omega_0$  and

$$(5.4) \quad \left( \Delta_k - \left(\frac{n}{2} - k + \ell\right) \left(\frac{n}{2} - k - \ell\right) \right) \omega_{F_1} = O_t(x^{\frac{n}{2} - k + \ell}) + O_n(x^{\frac{n}{2} - k + \ell + 2})$$

where the forms  $\omega_j^{(t)}$  on  $M$  are uniquely determined by  $\omega_0$  and the expansion of  $\Delta_k$  in powers of  $x$  given by Lemma 2.1.

Note that the condition  $\ell \leq \frac{n}{2}$  when  $n$  is even insures that that the first  $\log(x)$  coefficient coming from the metric does not show up in (5.1). Remark that when  $\ell > n/2 - k$  and  $n$  odd, the fact that the indicial root (5.2) is not in  $n/2 - k - \ell + 2\mathbb{N}_0$  does not affect the construction. Since  $(\frac{n}{2} - k + \ell)$  is an indicial root in the  $\Lambda_t^k$  component, we can then define

$$(5.5) \quad \begin{aligned} & \omega_{F_2} = \omega_{F_1} + x^{\frac{n}{2} - k + \ell} \log(x) \omega_{n-k+\ell,1}^{(t)}, \\ & \text{with } \omega_{n-k-\ell,1}^{(t)} = \frac{1}{2\ell} \left[ x^{-\frac{n}{2} + k - \ell} \left( \Delta_k - \left(\frac{n}{2} - k + \ell\right) \left(\frac{n}{2} - k - \ell\right) \right) \omega_{F_1} \right]_{|x=0} \end{aligned}$$

which satisfies

$$(5.6) \quad \left( \Delta_k - \left(\frac{n}{2} - k + \ell\right) \left(\frac{n}{2} - k - \ell\right) \right) \omega_{F_2} = O(x^{\frac{n}{2} - k + \ell + 2} \log x).$$

**Remark:** we could continue the construction to get a solution  $\omega$  of

$$\left( \Delta_k - \left(\frac{n}{2} - k - \ell\right) \left(\frac{n}{2} - k + \ell\right) \right) \omega = O(x^\infty)$$

and even an exact solution (with no  $O(x^\infty)$ ) using the resolvent of  $\Delta_k$ . However, since the mapping properties of  $(\Delta_k - (\frac{n}{2} - k - \ell)(\frac{n}{2} - k + \ell))^{-1}$  is not explicitly available in the literature when  $\ell \neq \frac{n}{2} - k$ , we do not discuss this case further.

Like we did for  $L_k$ , we can then define an operator on  $M$  as follows:

**Definition 5.2.** *For  $k \leq (n+1)/2$ , we let  $\ell$  be an integer in  $[1, \frac{n}{2} - k]$  if  $n$  is even and in  $\mathbb{N}$  if  $n$  is odd. The operator  $L_k^\ell : C^\infty(M, \Lambda^k(M)) \rightarrow C^\infty(M, \Lambda^k(M))$  is defined by  $L_k^\ell \omega_0 := \omega_{n-k-\ell,1}^{(t)}$  where  $\omega_{n-k-\ell,1}^{(t)}$  is given in (5.5).*

**Remark:** clearly, we have  $L_k^{\frac{n}{2} - k} = L_k$  when  $n$  is even.

**Lemma 5.3.** *The form  $\omega_{F_1}$  of (5.3) satisfies  $\delta_g \omega_{F_1} = O(x^{\frac{n}{2} - k + \ell + 2})$ .*

*Proof:* by (5.4) and  $\delta_g \Delta_k = \Delta_{k-1} \delta_g$ , the form  $\delta_g \omega_{F_1}$  solves

$$(5.7) \quad \left( \Delta_{k-1} - \left(\frac{n}{2} - k + \ell\right) \left(\frac{n}{2} - k - \ell\right) \right) \delta_g \omega_{F_1} = O(x^{\frac{n}{2} - k + \ell + 2})$$

and with  $\delta_g \omega_{F_1} = O(x^{\frac{n}{2}-k-\ell+2})$ . The Taylor series  $T$  of  $x^{-\frac{n}{2}+k+\ell} \delta_g \omega_{F_1}$  to order  $O(x^{2\ell+2})$  is such that  $x^{\frac{n}{2}-k-\ell} T$  solves (5.7), and moreover  $T$  is even by Lemma 2.1. A short computation shows that the indicial roots of  $(\Delta_{k-1} - (\frac{n}{2} - k + \ell)(\frac{n}{2} - k - \ell))$  are

$$\begin{aligned} \frac{n}{2} - k + 1 \pm \sqrt{\ell^2 + n + 1 - 2k} & \text{ in } \Lambda_t^{k-1} \text{ component ,} \\ \frac{n}{2} - k + 2 \pm \sqrt{\ell^2 + 2(n - 2k + 2)} & \text{ in } \Lambda_n^{k-1} \text{ component .} \end{aligned}$$

Thus if  $\ell \leq n/2 - k$ , the indicial roots are not contained in  $[\frac{n}{2} - k - \ell + 2, \frac{n}{2} - k + \ell + 2]$  except when  $2k = n + 1$  where  $\frac{n}{2} - k + \ell + 1$  is a root in the  $\Lambda_t$  component, this implies that the Taylor series of  $\delta_g \omega_{F_1}$  vanishes to order  $O(x^{\frac{n}{2}-k+\ell+1})$  except maybe when  $2k = n + 1$ . However in the last case, by parity of  $T$ , we see that there is no  $\frac{n}{2} - k + \ell + 1$  term in the expansion of  $\delta_g \omega_{F_1}$ , this ends the proof for  $n$  even. When  $n$  is odd and  $\ell > n/2 - k$ , the only indicial root in the interval of interest is  $n/2 - k + 1 + \sqrt{\ell^2 + n + 1 - 2k}$  and it is in  $(\frac{n}{2} - k + \ell + 1, \frac{n}{2} - k + \ell + 2)$  thus not in  $n/2 - k - \ell + \mathbb{N}_0$ , which shows that the argument used for  $\ell \leq n/2 - k$  applies similarly.  $\square$

By an obvious integration by part, we have the

**Proposition 5.4.** *The operators  $L_k^\ell$  are symmetric on  $C^\infty(M, \Lambda^k(M))$ .*

*Proof:* Consider  $\omega_{F_2}^1$  and  $\omega_{F_2}^2$  like in (5.5) with respective boundary values  $\omega_0^1$  and  $\omega_0^2$ , they are well defined form in some collar neighbourhood  $X_1 := (0, \epsilon_0)_x \times M$  of  $M$  in  $\bar{X}$ . Let  $\varphi \in C_0^\infty((-\epsilon_0, \epsilon_0))$  be a cut-off function which equals 1 near 0 and  $\tilde{\omega}^i := \varphi(x) \omega_{F_2}^i$  for  $i = 1, 2$ . Then using Lemma 5.3 we have  $\delta_g \tilde{\omega}^i = O(x^{\frac{n}{2}-k+\ell+1})$ , but since  $i_x \partial_x \tilde{\omega}^i = O(x^{\frac{n}{2}-k-\ell+2})$ , the Green formula gives for small  $\epsilon > 0$

$$\int_{x \geq \epsilon} (\langle \Delta_k \tilde{\omega}^1, \tilde{\omega}^2 \rangle_g - \langle \Delta_k \tilde{\omega}^2, \tilde{\omega}^1 \rangle_g) d\text{vol}_g = (-1)^n \int_{x=\epsilon} (\star_g d\tilde{\omega}^1) \wedge \tilde{\omega}^2 - (\star_g d\tilde{\omega}^2) \wedge \tilde{\omega}^1 + O(\epsilon).$$

But the first line is a  $O(1)$  as  $\epsilon \rightarrow 0$  by (5.6), and a straightforward analysis gives that the second line has an expansion of the form

$$\begin{aligned} a_{-2\ell-1} \epsilon^{-2\ell-1} + \dots + a_{-1} \epsilon^{-1} + L \log(\epsilon) + O(1) \\ \text{with } L := (-1)^n \left( \frac{n}{2} - k + \ell \right) \int_M (\star_0 L_k^\ell \omega_0^1) \wedge \omega_0^2 - (\star_0 L_k^\ell \omega_0^2) \wedge \omega_0^1 \end{aligned}$$

and this implies  $L = 0$  by comparing the  $\log(\epsilon)$  terms.  $\square$

**Proposition 5.5.** *We have  $L_k^l = \frac{(-1)^{l+1}}{2^{2l-1}(l!)^2} [(\delta_0 d)^l + \frac{n-2k-2l}{n-2k+2l} (d\delta_0)^l] + \text{LOT}$ .*

*Proof:* We define  $T$  by  $\omega_{F_1} = x^{\frac{n}{2}-k-l} T$ ,  $\lambda = (\frac{n}{2} - k + l)(\frac{n}{2} - k - l)$  and  $P = x^{k+l-\frac{n}{2}} (\Delta - \lambda) x^{\frac{n}{2}-k-l}$ .

Then we have  $T = \sum_{i=0}^{l-1} x^{2i} \omega_{2i}^{(t)} + \sum_{i=1}^l x^{2i} \omega_{2i}^{(n)} \wedge \frac{dx}{x}$  and  $P$  admits the same decomposition as  $\Delta_k$  in Lemma 2.1 but with indicial operator equal to

$$\begin{pmatrix} 2lx\partial_x - (x\partial_x)^2 & 2(-1)^{k+1}d \\ 0 & -(x\partial_x)^2 + 2(l+1)x\partial_x + n - 2k - 2l \end{pmatrix}$$

The equation  $PT = O_t(x^{2l}) + O_n(x^{2l+1})$  gives then

$$\omega_{2i+2}^{(n)} = (a_{2i+1}(\delta_0 d)^i \delta_0 + \text{LOT}) \omega_0 \quad \omega_{2i}^{(t)} = (a_{2i}(\delta_0 d)^i + b_{2i}(d\delta_0)^i + \text{LOT}) \omega_0$$

with

$$a_1 = \frac{(-1)^k}{\frac{n}{2} - k + l}, \quad a_2 = \frac{-1}{4(l-1)}, \quad b_2 = \frac{-(\frac{n}{2} - k + l - 2)}{4(l-1)(\frac{n}{2} - k + l)}$$

and

$$a_{2i+2} = \frac{a_{2i}}{4(i+1)(i+1-l)}, \quad a_{2i+1} = \frac{2(-1)^{k+1}b_{2i} + a_{2i-1}}{4(i+1)(i-l) + 2k - n + 2l},$$

$$b_{2i+2} = \frac{b_{2i} + 2(-1)^{k+1}a_{2i+1}}{4(i+1)(i+1-l)}.$$

The solutions of these equations are

$$b_{2i} = \frac{(-1)^i(n-2k+2l-4i)(l-i-1)!}{4^i i! (l-1)! (n-2k+2l)}, \quad a_{2i+1} = \frac{(-1)^{k+i}(l-i-1)!}{2^{2i-1} i! (l-1)! (n-2k+2l)}$$

$$a_{2i} = \frac{(-1)^i(l-i-1)!}{4^i i! (l-1)!}$$

Since the equation (5.5) reads

$$L_k^l = \left[ \frac{x^{-2l}}{2l} i_{x\partial_x} \left( \frac{dx}{x} \wedge PT \right) \right]_{|x=0}$$

we get the result.  $\square$

## 6. RELATION WITH BRANSON-GOVER OPERATORS

First we recall a few fact on the ambient metric of Fefferman-Graham, see [4, 6] for details. If  $(M, [h_0])$  is a compact manifold equipped with a conformal class, we call

$$\mathcal{Q} = \{t^2 h_0(m); t > 0, m \in M\} \subset S^2 T^* M$$

the conformal bundle, it is identified with  $(0, \infty)_t \times M$ . Let  $\tilde{\mathcal{Q}} = (-1, 1) \times \mathcal{Q}$  be the ambient space with the inclusion  $\iota : \mathcal{Q} \rightarrow \tilde{\mathcal{Q}}$  defined by  $z \rightarrow (0, z)$ . There are dilations  $\delta_s : (t, m) \rightarrow (st, m)$  of  $\mathcal{Q}$  which extends naturally to  $\tilde{\mathcal{Q}}$ . The functions on  $\mathcal{Q}$  which are  $w$ -homogeneous in the sense

$$f(st, m) = s^w f(t, m)$$

are the section of a bundle denoted  $E[w]$ , they extend naturally on  $\tilde{\mathcal{Q}}$ . We denote by  $\tilde{h}$  the ambient metric of Fefferman-Graham [4] on  $\tilde{\mathcal{Q}}$ . This is a smooth Lorentzian metric on  $\mathcal{Q}$  such that

- (1)  $\delta_s^* \tilde{h} = s^2 \tilde{h}, \forall s > 0$ ,
- (2)  $\iota^* \tilde{h}$  is the tautological tensor  $t^2 h_0$  on  $\mathcal{Q}$ ,
- (3\*)  $\text{Ric}(\tilde{h})$  vanishes to infinite order at  $\mathcal{Q}$  if  $n$  is odd,
- (3\*\*)  $\text{Ric}(\tilde{h})$  vanishes to order  $\frac{n}{2} - 1$  at  $\mathcal{Q}$  if  $n$  is even.

We let  $T$  be the vector field which generates the dilations  $\delta_s$ , and let

$$Q = \tilde{h}(T, T), \quad \rho := -t^{-2}Q/2, \quad x = \sqrt{2\rho}, \quad u = xt$$

so that  $Q$  is homogeneous of degree 2 with respect to  $\delta_s$ ,  $u$  and  $t$  are homogeneous of degree 1 and  $x$  of degree 0, moreover  $Q, \rho$  are smooth defining function of  $\mathcal{Q}$ ,  $x, u$  are defining function of  $\mathcal{Q}$  in  $\{Q \leq 0\}$  for some finer smooth structure on  $\{Q \leq 0\}$ . Let us define  $\mathcal{C} := \{Q = -1, \rho < \epsilon\}$  for some small fixed  $\epsilon$ , then  $\mathcal{C}$  can be identified with a collar  $(0, \epsilon)_\rho \times M$  and there is a system of coordinates  $(u, m) \in (0, 1] \times \mathcal{C}$  that covers the part  $\{0 > Q \geq -1, \epsilon > \rho > 0\}$  which is a neighbourhood of the cone  $\mathcal{Q}$  near  $t = \infty$ . The metric  $\tilde{h}$  has the model form (see [4]) in this neighbourhood

$$\tilde{h} = -du^2 + u^2 g$$

where  $g = (dx^2 + h_x)/x^2$  is a Poincaré-Einstein metric on the collar  $\mathcal{C}$ .

The space  $\mathcal{T}^k[s]$  is the space of  $k$ -form tractors which are homogeneous of degree  $s$ , i.e. these are restrictions to the null cone  $\mathcal{Q}$  of  $k$ -forms on  $\tilde{\mathcal{Q}}$  and such that  $\tilde{\nabla}_T F = sF$  where  $T = t\partial_t = u\partial_u$  is the generator of dilations in the cone fibers,  $\tilde{\nabla}$  is the Levi-Civita

connection on  $\tilde{\mathcal{Q}}$ . Since  $\tilde{\nabla}_T^* = *$  for  $*$  =  $T, \partial_x, \partial_{m_i}$ , we have  $\mathcal{L}_T = \tilde{\nabla}_T + k$ , on  $\mathcal{J}^k[s]$ , where  $\mathcal{L}$  denotes Lie derivative. The bundle  $\mathcal{E}^k[s]$  is the bundle which consists of the  $s$ -homogeneous  $k$  forms on  $M$ , in the sense that they are the sections of  $\Lambda^k T^*M \otimes E[s]$  and thus satisfy  $\mathcal{L}_T \omega = s\omega$ . We can view  $\mathcal{E}^k[s]$  as a subspace of  $\mathcal{J}^k[s-k]$ . We let  $\mathcal{G}_k[s]$  be the subbundle of  $\mathcal{J}^k[s+k-n]$  consisting of forms which are annihilated by the interior product  $i_T$ . It has a conformally invariant projection onto  $\mathcal{E}^k[s+2k-n]$  denoted by  $q^k$ , this is given for instance by  $i_{\partial_\rho} d\rho \wedge$ .

If  $\tilde{\Delta}$  is the ambient Laplacian on  $\tilde{\mathcal{Q}}$  associated to  $\tilde{h}$ , if  $\omega_0 \in \mathcal{E}^k[k+\ell-\frac{n}{2}]$  and  $\tilde{\omega}_0$  is an homogeneous extension of  $\omega_0$  to  $\tilde{\mathcal{Q}}$ , then it is proved in [3, Prop. 4.3] that the operator defined by the formula

$$(6.1) \quad \mathbf{L}_k^\ell \omega_0 = \left[ \iota_T \left( \tilde{d}(n+2\tilde{\nabla}_T-2) + \frac{1}{2} \tilde{d}Q \wedge \tilde{\Delta} \right) \tilde{\Delta}^\ell \tilde{\omega}_0 \right]_{|\Omega} = \left[ \iota_T \tilde{d}(n+2\tilde{\nabla}_T-2) \tilde{\Delta}^\ell \tilde{\omega}_0 \right]_{|\Omega}$$

can be viewed as a conformally invariant operator  $\mathbf{L}_k^\ell : \mathcal{E}^k[k+\ell-\frac{n}{2}] \rightarrow \mathcal{G}_k[\frac{n}{2}-k-\ell]$ . Here  $\tilde{d}$  denotes the exterior differential on  $\tilde{\mathcal{Q}}$ . They also define the operators (see Proposition 4.4 and Theorem 4.5 in [3])

$$(6.2) \quad \begin{aligned} L_k^{\text{BG},\ell} &:= q^k \mathbf{L}_k^\ell : \mathcal{E}^k[k+\ell-\frac{n}{2}] \rightarrow \mathcal{E}^k[k-\frac{n}{2}-\ell], \\ G_k^{\text{BG}} &:= q^{k-1} i_Y \mathbf{L}_k^{\frac{n}{2}-k} : \mathcal{E}^k[0] \rightarrow \mathcal{E}^{k-1}[2k-2-n] \end{aligned}$$

where  $Y = -\frac{\partial_\rho}{t^2}$  is a vector field dual to  $\tilde{d}t/t$  via  $\tilde{h}$ , it satisfies in particular  $\tilde{d}Q(Y) = 2$ . Finally the operator  $Q_k^{\text{BG}}$  acting on a closed  $k$ -form  $\omega_0$  is defined as follows

$$Q_k^{\text{BG}} \omega_0 := -2 \left( \frac{n}{2} - k + 1 \right) q^k \left[ i_Y i_T \tilde{\Delta}^{\frac{n}{2}-k} \left( \frac{dQ}{2} \wedge \frac{\tilde{d}t}{t} \wedge \tilde{\omega}_0 \right) \right]_{|\Omega}.$$

where  $\tilde{\omega}_0$  is any homogeneous extension of  $\omega_0$  to  $\tilde{\mathcal{Q}}$ .

We now prove a Lemma whose proof is essentially the same as the for functions in [11].

**Lemma 6.1.** *Let  $\omega \in \mathcal{J}^{k'}[-\alpha]$  and  $j \in \mathbb{N}$ , then we have*

$$\tilde{\Delta}(Q^j \omega) = 4j \left( \alpha - \frac{n}{2} - j \right) Q^{j-1} \omega + Q^j \tilde{\Delta} \omega.$$

*Proof:* Using  $\tilde{\nabla}Q = 2T$ , we have  $[\tilde{\Delta}, Q] = -2(2\tilde{\nabla}_T + n + 2)$  and so we can compute

$$\begin{aligned} \tilde{\Delta}(Q^j L \omega_0) &= \sum_{m=0}^{j-1} Q^m [\tilde{\Delta}, Q] Q^{j-1-m} \omega + Q^j \tilde{\Delta} \omega \\ &= -2Q^{j-1} \sum_{m=0}^{j-1} (2(-2m+2j-2-\alpha) + n+2) \omega + Q^j \tilde{\Delta} \omega \\ &= 4Q^{j-1} j \left( \alpha - \frac{n}{2} - j \right) \omega + Q^j \tilde{\Delta} \omega \end{aligned}$$

which achieves the proof. □

As a consequence, and using Lemma 6.1, we get the

**Theorem 6.2.** *(i) Let  $L_k^\ell, L_k$  and  $G_k$  be the operators of Definition 5.2 and 4.1, and let  $c_k^\ell := (-4)^\ell (\ell-1)! (\ell+1)! (k-\frac{n}{2}-\ell)$ . Then the following identity holds*

$$L_k^{\text{BG},\ell} = c_k^\ell L_k^\ell.$$

*In the critical case  $\ell = \frac{n}{2} - k$ , if  $G_k^{\text{BG}}$  is the Branson-Gover operator of (6.2) we have*

$$L_k^{\text{BG}} = c_k L_k, \quad G_k^{\text{BG}} = (-1)^k c_k G_k$$

with  $c_k := (-1)^{\frac{n}{2}-k-1} 2^{n-2k+1} ((\frac{n}{2}-k)!)^2 (\frac{n}{2}-k+1) = c_k^{\frac{n}{2}-k}$ .

(ii) Let  $Q_k$  be the operator of Definition 4.2, then

$$Q_k^{\text{BG}} = (2k - n - 2)c_{k+1}Q_k.$$

*Proof:* (i) For  $\omega_0 \in \Lambda^k(M)$ , we consider the form  $\omega_{F_1}$  of Lemma 5.1 of the previous section and we extend it homogeneously in a smooth  $k$ -form of degree  $k - \frac{n}{2} + \ell$  by

$$\begin{aligned} \tilde{\omega}_F &= u^{k-\frac{n}{2}+\ell} \omega_{F_1} = u^{k-\frac{n}{2}+\ell} x^{\frac{n}{2}-k-\ell} \sum_{i=0}^{\ell} x^{2i} (\omega_i^{(t)} + x^2 \omega_i^{(n)}) \wedge \frac{dx}{x} \\ &= t^{k-\frac{n}{2}+\ell} \sum_{i=0}^{\ell} (-Q)^i t^{-2i} (\omega_i^{(t)} + \omega_i^{(n)}) \wedge d\rho. \end{aligned}$$

In the coordinates  $u, x, y$  representing a neighbourhood  $\{-1 \leq Q < 0, \rho < \epsilon\}$  and in the  $k$ -form bundle decomposition  $\Lambda^k(\mathcal{C}) \oplus \Lambda^{k-1}(\mathcal{C}) \wedge \frac{du}{u}$ , the exterior derivative, its dual and the form Laplacian of  $\tilde{h}$  are given by

$$(6.3) \quad \tilde{d} = \begin{pmatrix} d & 0 \\ (-1)^k u \partial_u & d \end{pmatrix}, \quad \tilde{\delta} = u^{-2} \begin{pmatrix} \delta_g & (-1)^{k+1} (n+2-2k+u\partial_u) \\ 0 & \delta_g \end{pmatrix}$$

and

$$(6.4) \quad \tilde{\Delta} = u^{-2} \begin{pmatrix} (u\partial_u)(u\partial_u + n - 2k) + \Delta_k & 2(-1)^{k+1} d \\ 2(-1)^k \delta_g & (u\partial_u - 2)(u\partial_u + n - 2k + 2) + \Delta_{k-1} \end{pmatrix}.$$

So, using the properties of  $\omega_{F_1}$  in Lemma 5.1 and Lemma 5.3, we have (where  $s = k - \frac{n}{2} + \ell$ )

$$\begin{aligned} \tilde{\Delta} \tilde{\omega}_F &= u^{s-2} (\Delta_k + s(s+n-2k)) \omega_{F_1} + 2(-1)^k u^{s-3} \delta_g \omega_{F_1} \wedge du \\ &= 2\ell u^{s-2} x^{\ell-k+\frac{n}{2}} (L_k^\ell \omega_0 + O_t(x^2)) + u^{s-2} x^{\ell-k+\frac{n}{2}+2} (B \wedge \frac{dx}{x} + O_n(x^2)) \\ &\quad + 2(-1)^k u^{s-3} x^{\ell-k+\frac{n}{2}+2} (C + O(x^2)) \wedge du \\ &= (-Q)^{\ell-1} t^{k-\ell-\frac{n}{2}} (2\ell L_k^\ell \omega_0 + (B + 2(-1)^k C) \wedge d\rho) + O(Q^\ell) \end{aligned}$$

for some  $(k-1)$ -forms  $B, C$  on  $M$ . We can now apply  $\ell-1$  times Lemma 6.1 and get

$$\tilde{\Delta}^\ell \tilde{\omega}_F = \tilde{\Delta}^{\ell-1} \tilde{\Delta} \omega_F = (-4)^{\ell-1} [(\ell-1)!]^2 t^{k-\ell-\frac{n}{2}} (2\ell L_k^\ell \omega_0 + (B + 2(-1)^k C) \wedge d\rho) + O(Q).$$

Since  $(n + 2\tilde{\nabla}_T - 2)$  acts on homogeneous  $k$ -forms of degree  $k - \ell - \frac{n}{2}$  by multiplication by  $-2(\ell+1)$  and  $i_T \tilde{d} = \mathcal{L}_T$  on  $\mathcal{G}_k[\frac{n}{2} - k - \ell]$ , we get

$$\mathbf{L}_k^\ell = (\ell+1)(k-\ell-\frac{n}{2})(-4)^\ell [(\ell-1)!]^2 t^{k-\ell-\frac{n}{2}} (\ell L_k^\ell \omega_0 + (\frac{B}{2} + (-1)^k C) \wedge d\rho)$$

Note that by definition of  $B_k, C_k, G_k$  we have, in the case  $\ell = \frac{n}{2} - k$ ,

$$\frac{B}{2} + (-1)^k C = (-1)^{k-1} (\frac{B_k}{2} - C_k) \omega_0 = \ell G_k \omega_0.$$

(ii) Similarly, for  $\omega_0 \in \Lambda^k(M)$  closed, we set  $\tilde{\omega}_F := \omega'_{F_1} \wedge \frac{1}{2} \tilde{d}Q$  in  $\{Q < 0, \rho \leq \epsilon\}$  where the form  $\omega'_{F_1}$  is the 0-homogeneous expansion of  $\omega'_{F_1} \in \Lambda^{k+1}(\mathcal{C})$  given by (3.12). Since  $\frac{\tilde{d}Q}{2} = -t^2 d\rho + Q \frac{dt}{t}$ , we have

$$\begin{aligned} \tilde{\omega}_F &= \sum_{j=0}^{\frac{n}{2}-k} x^{2j} (\omega_{2j}^{(n)} \wedge \frac{dx}{x} + x^2 \omega_{2j}^{(t)}) \wedge \frac{\tilde{d}Q}{2} \\ &= \sum_{j=0}^{\frac{n}{2}-k} -(-Q)^j t^{-2(j-1)} \omega_{2j}^{(n)} \wedge d\rho \wedge \frac{dt}{t} + (-Q)^{j+1} t^{-2j-2} \omega_{2j}^{(t)} \wedge (-t^2 d\rho + Q \frac{dt}{t}) \end{aligned}$$

and so  $\tilde{\omega}_F$  is a smooth  $(k+2)$  form. By (6.4) and the definition of  $B'_k, D'_k$  we have

$$\begin{aligned}\tilde{\Delta}\tilde{\omega}_F &= \tilde{\Delta}(-u^2\omega'_{F_1} \wedge \frac{du}{u}) = 2(-1)^k d\omega'_{F_1} - (\Delta_{k+1}\omega'_{F_1}) \wedge \frac{du}{u} \\ &= -x^{n-2k-2}2D'_k\omega_0 \wedge \frac{dx}{x} + (-1)^{k+1}x^{n-2k}(B'_k\omega_0 \wedge \frac{dx}{x} + \omega_1) \wedge \frac{du}{u} + O(Q^{\frac{n}{2}-k}) \\ &= (-1)^{\frac{n}{2}-k-1}2Q^{\frac{n}{2}-k-2}t^{2k-n+4}D'_k\omega_0 \wedge d\rho \\ &\quad + (-1)^{\frac{n}{2}+1}Q^{\frac{n}{2}-k-1}t^{2k-n+2}(B'_k\omega_0 \wedge \frac{dt}{t} + (-1)^k\omega_1) \wedge d\rho + O(Q^{\frac{n}{2}-k})\end{aligned}$$

for some form  $\omega_1$  on  $M$ , the value of which is not important for our purpose. By Lemma 6.1, we have

$$\begin{aligned}\tilde{\Delta}^2\tilde{\omega}_F &= (-1)^{\frac{n}{2}-k-1}2Q^{\frac{n}{2}-k-2}t^{2k-n+4}\tilde{\Delta}(D'_k\omega_0 \wedge d\rho) \\ &\quad + 4\left(\frac{n}{2}-k-1\right)(-1)^{\frac{n}{2}+1}Q^{\frac{n}{2}-k-2}t^{2k-n+2}(B'_k\omega_0 \wedge \frac{dt}{t} + (-1)^k\omega_1) \wedge d\rho + O(Q^{\frac{n}{2}-k-1})\end{aligned}$$

and by (6.4), we have

$$\begin{aligned}\tilde{\Delta}(D'_k\omega_0 \wedge d\rho) &= \tilde{\Delta}(x^2D'_k\omega_0 \wedge \frac{dx}{x}) \\ &= u^{-2}\Delta_{k+2}(x^2D'_k\omega_0 \wedge \frac{dx}{x}) + 2u^{-2}(-1)^k\delta_g(x^2D'_k\omega_0 \wedge \frac{dx}{x}) \wedge \frac{du}{u} \\ &= 2(-1)^kt^{-2}\delta_0D'_k\omega_0 \wedge d\rho \wedge \frac{dt}{t} + 2(2k-n-4)t^{-2}D'_k\omega_0 \wedge \frac{dt}{t}\end{aligned}$$

where we have used (2.2), (2.3) and  $dD'_k\omega_0 = 0$ . We thus have

$$\begin{aligned}\tilde{\Delta}^2\tilde{\omega}_F &= Q^{\frac{n}{2}-k-2}t^{2k-n+2}\left[4(-1)^{\frac{n}{2}-1}\left(\delta_0D'_k\omega_0 - \left(\frac{n}{2}-k-1\right)B'_k\omega_0\right) \wedge d\rho \wedge \frac{dt}{t} + \omega'_1 \wedge \frac{dt}{t}\right. \\ &\quad \left.+ \omega'_2 \wedge d\rho\right] + O(Q^{\frac{n}{2}-k-1}) \\ &= Q^{\frac{n}{2}-k-2}t^{2k-n+2}\left[4(n-2k)\left(\frac{n}{2}-k-1\right)(-1)^{\frac{n}{2}+k}Q_k\omega_0 \wedge d\rho \wedge \frac{dt}{t} + \omega'_1 \wedge \frac{dt}{t}\right. \\ &\quad \left.+ \omega'_2 \wedge d\rho\right] + O(Q^{\frac{n}{2}-k-1})\end{aligned}$$

by Corollary 4.9, and where  $\omega'_1, \omega'_2$  are forms in  $\Lambda^{k+1}(M)$ . By iterative use of Lemma 6.1, we get

$$\begin{aligned}\tilde{\Delta}^{\frac{n}{2}-k}\tilde{\omega}_F &= \tilde{\Delta}^{\frac{n}{2}-k-2}\tilde{\Delta}^2\tilde{\omega}_F \\ &= t^{2k-n+2}\left[2^{n-2k-1}\left(\frac{n}{2}-k\right)\left[\left(\frac{n}{2}-k-1\right)!\right]^2(-1)^{\frac{n}{2}+k}Q_k\omega_0 \wedge d\rho \wedge \frac{dt}{t}\right. \\ &\quad \left.+ \omega'_1 \wedge \frac{dt}{t} + \omega'_2 \wedge d\rho\right] + O(Q)\end{aligned}$$

we infer from the definition of  $Q_k^{\text{BG}}$  that

$$Q_k^{\text{BG}} = (-1)^{\frac{n}{2}+k+1}2^{n-2k}\left(\frac{n}{2}-k+1\right)!\left(\frac{n}{2}-k-1\right)!Q_k$$

□

## 7. PROOF OF THE MAIN RESULTS

We start with the proof of Theorem 1.2.

*Proof of Theorem 1.2:* the existence of  $\omega$  in (i) is proved in Proposition 3.1. The fact that the log terms  $L_k, Q_k$  coincide with the Branson-Gover operators follows from Theorem 6.2. The uniqueness of the solution is rather clear by construction: using the arguments used in the proof of Proposition 3.1, a solution in  $C^{\frac{n}{2}-k, \alpha}(\bar{X}, \Lambda^k(\bar{X}))$  would have its first  $\frac{n}{2}-k$  Taylor coefficients uniquely (and locally) determined by the boundary

value  $\omega_0$  and then two such solutions with same boundary data would have a difference in  $O_t(x^{\frac{n}{2}-k+\alpha}) + O_n(x^{\frac{n}{2}-k+1+\alpha})$  and would then be in  $\ker_{L^2}(\Delta_k)$  (see Remark 3.6). We conclude that any  $C^{\frac{n}{2}-k,\alpha}$  solution is in fact a solution of Proposition 3.1. The proof of (ii) is similar and follows from Proposition 3.11 and Theorem 6.2.  $\square$

*Proof of Theorem 1.1:* The infinite dimensionality of  $K_m^k(\bar{X})$  for  $m < n - 2k + 1$  follows from Proposition 3.1. Indeed for  $m < n - 2k$  this is clear since the solutions of (3.1) are parametrized by  $C^\infty(M, \Lambda^k(M))$ . If  $m = n - 2k$ , one can use that there is an infinite set of  $\omega_0 \in C^\infty(M, \Lambda^k(M))$  such that  $G_k\omega_0 \neq 0$  and  $L_k\omega_0 = 0$  since  $\ker L_k$  is infinite dimensional and  $\ker G_k \cap \ker L_k$  is finite dimensional by ellipticity of  $dG_k + L_k$ . Solutions of (3.1) are then in  $C^{n-2k}(\bar{X}, \Lambda^k(\bar{X}))$ .

The finite dimensionality for  $m = n - 2k + 1$  is a little more involved. Let  $\omega$  be a harmonic form in  $C^{n-2k+1}(\bar{X}, \Lambda^k(\bar{X}))$ , then Taylor expanding, there exist some forms  $\omega_j^{(n)}, \omega_j^{(t)} \in C^{n-2k+1-j}(M, \Lambda(M))$  so that

$$\omega - \sum_{j=0}^{n-2k} x^j (\omega_j^{(t)} + \omega_j^{(n)}) \wedge dx \in x^{n-2k+1} L^\infty(\Lambda^k(\bar{X})),$$

and  $L_k\omega_0 = 0$ . Now by Lemma 3.8 we know that  $\omega$  has a weak expansion to order  $x^N$  with values in  $H^{-N}(M)$  like in (3.9) for any  $N > 0$  large. Moreover  $\delta_g\omega$  is also a harmonic form in  $C^{n-2k}(\bar{X}, \Lambda^{k-1}(M))$ , moreover using (2.2) after decomposing the form in  $\Lambda_t^{k-1} \oplus \Lambda_n^{k-1}$ , we see that it is a  $O(x)$  and has an expansion to order  $x^N$  with values in  $H^{-N-1}(M)$  for any  $N$ . Now, using the indicial equation like in the proof of Proposition 3.1, the weak expansion of  $\delta_g\omega$  vanishes to order  $x^{n-2k+2}$ , so in particular we obtain  $\delta_g\omega \in x^{n-2k} L^\infty(\Lambda^{k-1}(\bar{X}))$  from the regularity of  $\omega$ . Then  $\delta_g\omega \in L^2(\Lambda^{k-1}(X))$  for  $k < \frac{n}{2} - 1$ , while for  $k = \frac{n}{2} - 1$  it is in  $L^2$  if we assume in addition that  $\omega \in C^{n-2k+1,\alpha}(\bar{X}, \Lambda^k(\bar{X}))$  for some  $\alpha > 0$  (since then  $\delta_g\omega \in x^{n-2k+\alpha} L^\infty(\Lambda^k(\bar{X}))$ ). But as shown in Corollary 3.3,  $\delta_g\omega$  is then smooth and vanishing to order  $x^{n-2k+3}$ . Then, like in the proof of Proposition 3.9, an integration by part on  $\|\delta_g\omega\|_{L^2}^2$  shows that  $\delta_g\omega = 0$ . Now we can apply the result of Proposition 4.4 (see the Remark below Corollary 4.11), which gives  $G_k\omega_0 = 0$ . Since  $dG_k + L_k$  is elliptic,  $\ker L_k \cap \ker G_k$  is finite dimensional and contains only smooth forms, so  $\omega_0$  is smooth. Then  $\omega$  is polyhomogeneous and is the solution of Proposition 3.1, up to an element of  $\ker_{L^2}(\Delta_k)$ , it is then in  $C^{n-1}(\bar{X}, \Lambda^k(\bar{X}))$  in general and in  $C^\infty(\bar{X}, \Lambda^k(\bar{X}))$  if  $(X, g)$  smooth Poincaré-Einstein manifold.

Let  $m \in [n - 2k + 1, n - 1]$  be an integer. The exact sequence (1.9) is defined by inclusion of  $\iota : H^k(\bar{X}, \partial\bar{X}) \rightarrow K_m^k(\bar{X})$  and restriction to the boundary  $r : K_m^k(\bar{X}) \rightarrow \mathcal{H}_L^k(M)$ , here of course we use the identification  $H^k(\bar{X}, \partial\bar{X}) \simeq \ker_{L^2}(\Delta_k)$  and the regularity of harmonic  $L^2$  forms in Theorem 3.2. The injectivity of  $\iota$  is clear, the surjectivity of  $r$  comes from Proposition 3.1, the definition of  $\mathcal{H}_L^k$  and Theorem 6.2. The kernel of  $r$  is composed of those forms of  $K_m^k(\bar{X})$  which vanish at  $M$ , but by Proposition 3.1, these are  $L^2$ , and thus in the image of  $H^k(\bar{X}, \partial\bar{X})$  by the map  $\iota$ .  $\square$

*Proof of Theorem 1.3:* First note that the space  $Z^k(\bar{X})$  in Theorem 1.3 is included in  $K_{n-2k+1}^k(\bar{X})$ , and thus of finite dimension and composed of forms in  $C^{n-1}(\bar{X}, \Lambda^k(\bar{X}))$  (even in the case  $k = \frac{n}{2}$  by the arguments above).

(i) the maps in the complex

$$0 \rightarrow H^k(\bar{X}, \partial\bar{X}) \xrightarrow{\iota} Z^k(\bar{X}) \xrightarrow{r} \mathcal{H}^k(\partial\bar{X}) \xrightarrow{d_e} H^{k+1}(\bar{X}, \partial\bar{X})$$

are defined as follows:  $\iota$  is given by inclusion where  $H^k(\bar{X}, \partial\bar{X}) \simeq \ker_{L^2}(\Delta_k)$ , this is well defined since  $L^2$  harmonic forms are closed, coclosed and in  $C^{n-2k+1}(\bar{X}, \Lambda^k(\bar{X}))$ ;  $r$  is defined as restriction at the boundary and it maps in  $\mathcal{H}^k(M)$  since  $r(\omega) \in \ker L_k \cap \ker G_k$  by the discussion above and  $d\omega = 0$  implies  $dr(\omega) = 0$ ; the last map  $d_e$  is the composition

$d_e = d \circ \Phi$  where  $\Phi : C^\infty(M, \Lambda^k(M)) \rightarrow C^\infty(X, \Lambda^k(X)) / \ker_{L^2}(\Delta_k)$  is defined by  $\Phi(\omega_0) = \omega$  where  $\omega$  is the solution of (3.1) in Proposition 3.1. Note that  $\Phi$  is only defined modulo  $\ker_{L^2}(\Delta_k)$  and is linear by uniqueness of the solution in (3.1) modulo  $\ker_{L^2}(\Delta_k)$ . Applying  $d$  kills the indeterminacy with respect to  $\ker_{L^2}(\Delta_k)$  since  $L^2$  harmonic forms are closed. Then  $d\Phi(\omega_0)$  is harmonic and since the boundary value of  $\Phi(\omega_0)$  is closed, then  $d\Phi(\omega_0) = O(x)$ , and by Proposition 3.1 it is in  $L^2$ . For the exactness of the sequence, first note that  $\ker r$  is composed of closed and coclosed forms which are  $O(x)$ , this implies that those forms are  $L^2$  by Proposition 3.1, so  $\text{Im } \iota = \ker r$  since also  $L^2$  harmonic forms vanish at the boundary. Now  $\omega_0 \in \ker d_e$  if  $\Phi(\omega_0)$  is closed, but it is also coclosed and in  $C^{n-2k+1}(\bar{X}, \Lambda^k(\bar{X}))$  by Proposition 3.1 and the fact that  $\omega_0 \in \ker d \cap \ker G_k \subset \ker L_k \cap \ker G_k$ , therefore  $\Phi(\omega_0) \in Z^k(\bar{X})$  and  $\omega_0 \in \text{Im } r$ . Moreover by Proposition 3.1 we have  $\Phi(r(\omega)) - \omega \in \ker_{L^2}(\Delta_k)$ , this implies  $\text{Im } r \subset \ker d_e$ , this proves exactness of the sequence.

(ii) the map in the complex (1.11) are defined similarly: first  $\iota : H^k(\bar{X}, \partial\bar{X}) \rightarrow [Z^k(\bar{X})]$  is the composition of the inclusion  $\ker_{L^2}(\Delta_k) \rightarrow Z^k(\bar{X})$  with the natural map  $Z^k(\bar{X}) \rightarrow [Z^k(\bar{X})]$  obtained by taking cohomology class. The map  $r : [Z^k(\bar{X})] \rightarrow [\mathcal{H}^k(\partial\bar{X})]$  is the map induced by the restriction map  $Z^k(\bar{X}) \rightarrow \mathcal{H}^k(\partial\bar{X})$  used in (i). This is well defined since if  $d\alpha \in Z^k(\bar{X})$ , then  $r(d\alpha) = d\alpha_0$  where  $\alpha_0 = \alpha|_{\partial\bar{X}}$ , and so  $[r(d\alpha)] = 0$  if  $[\cdot]$  denotes cohomology class in  $H^k(\partial\bar{X})$ . The last map  $d_e : [Z^k(\partial\bar{X})] \rightarrow H^{k+1}(\bar{X}, \partial\bar{X})$  is the map induced by  $d_e$  defined in (i), i.e.  $d_e = d \circ \Phi$  where  $\Phi$  maps  $\omega_0$  to the solution of (3.1). Note that it is well defined since for  $d\alpha_0 \in \mathcal{H}^k(\partial\bar{X})$ , we have  $d_e(d\alpha_0) = d\Phi(d\alpha_0)$  and, by uniqueness of the solution of (3.1),  $\Phi(d\alpha_0) - d\Phi(\alpha_0) \in \ker_{L^2}(\Delta_{k+1})$  thus  $d\Phi(d\alpha_0) = 0$ .

To show that  $\ker r = \text{Im } \iota$ , we need to show that if  $\omega \in Z^k(\bar{X})$  is a representative in  $[Z^k(\bar{X})]$  such that  $r(\omega) = d\alpha_0$  for some smooth  $\alpha_0$ , then there is  $\omega' \in \ker_{L^2}(\Delta_k)$  such that  $\omega - \omega'$  is exact. But as said above, we have  $\Phi(d\alpha_0) - d\Phi(\alpha_0) \in \ker_{L^2}(\Delta_k)$  and  $\Phi(r(\omega)) - \omega \in \ker_{L^2}(\Delta_k)$  thus  $\omega - d\Phi(\alpha_0) \in \ker_{L^2}(\Delta_k)$  and we are done. To show that  $\ker d_e = \text{Im } r$ , we need to prove that for  $\omega_0 \in \mathcal{H}^k(\partial\bar{X})$  a representative in  $[\mathcal{H}^k(\partial\bar{X})]$  then  $\Phi(\omega_0)$  is closed if and only if there exists  $\omega \in Z^k(\bar{X})$  so that  $r(\omega) - \omega_0$  is exact. But  $\Phi(\omega_0)$  is in  $Z^k(\bar{X})$  if  $d\Phi(\omega_0) = 0$ , thus  $\ker d_e \subset \text{Im } r$ ; conversely if there is  $\omega \in Z^k(\bar{X})$  with  $\omega = \omega_0 + d\alpha_0 + O(x)$ , then  $\omega - \Phi(\omega_0 + d\alpha_0) \in \ker_{L^2}(\Delta_k)$  and so  $d\Phi(\omega_0) = 0$  since  $\Phi(d\alpha_0) - d\Phi(\alpha_0) \in \ker_{L^2}(\Delta_k)$ . To conclude, we need to prove that  $\text{Im } d_e \subset \ker \iota$ . But this is clear since  $d_e\omega_0 = d\Phi(\omega_0)$  is an exact  $(k+1)$ -form in  $L^2$  with  $\Phi(\omega_0) \in C^{n-2k+1}(\bar{X}, \Lambda^k(\bar{X}))$ . Note that in the case  $k = \frac{n}{2}$ , we make use of Proposition 3.10.

(iii) Suppose that  $[\mathcal{H}^k(\partial\bar{X})] = H^k(\partial\bar{X})$ . If  $\omega \in \ker \iota$ , it is a  $k$ -form in  $\ker_{L^2}(\Delta_k)$  which can be written  $\omega = d\alpha$  with  $\alpha$  smooth. Moreover if  $\alpha_0 = \alpha|_{\partial\bar{X}}$ , then  $d(\Phi(\alpha_0) - \alpha) \in \ker_{L^2}(\Delta_k)$  and  $\Phi(\alpha_0) - \alpha = O(x)$ , an easy integration by parts shows that  $d\Phi(\alpha_0) = d\alpha = \omega$ . Here  $\alpha_0$  is closed since  $\omega = O(x)$ , but by assumption there is a  $\alpha'_0 \in \mathcal{H}^k(\partial\bar{X})$  such that  $\alpha_0 - \alpha'_0 = d\beta$  for some smooth  $\beta$ . Since now  $d\Phi(d\beta) = d[\Phi, d]\beta = 0$ , we have  $d_e\alpha'_0 = \omega$  and  $\omega \in \text{Im } d_e$ , which gives  $\ker \iota = \text{Im } d_e$ . Eventually, the equality  $[Z^k(\bar{X})] = H^k(\bar{X})$  is clear from the discussion above since  $[Z^k(\bar{X})] \subset H^k(\bar{X})$  and

$$\begin{aligned} H^k(\bar{X}, \partial\bar{X}) &\xrightarrow{\iota} [Z^k(\bar{X})] \xrightarrow{r} H^k(\partial\bar{X}) \xrightarrow{d_e} H^{k+1}(\bar{X}, \partial\bar{X}) \\ H^k(\bar{X}, \partial\bar{X}) &\xrightarrow{\iota} H^k(\bar{X}) \xrightarrow{r} H^k(\partial\bar{X}) \xrightarrow{d_e} H^{k+1}(\bar{X}, \partial\bar{X}) \end{aligned}$$

are both exact sequences.

As for the converse, if  $\ker \iota^{k+1} = \text{Im } d_e^k$  and  $[Z^k(\bar{X})] = H^k(\bar{X})$ , then we have the exact sequences

$$\begin{aligned} H^k(\bar{X}) &\xrightarrow{r} [\mathcal{H}^k(M)] \xrightarrow{d_e} H^{k+1}(\bar{X}, M) \xrightarrow{\iota} [Z^{k+1}(\bar{X})] \\ H^k(\bar{X}) &\xrightarrow{r} H^k(M) \xrightarrow{d_e} H^{k+1}(\bar{X}, M) \xrightarrow{\iota'} H^{k+1}(\bar{X}) \end{aligned}$$

and since  $[Z^{k+1}(\bar{X})] \subset H^{k+1}(\bar{X})$ , we obviously have  $\ker \iota = \ker \iota' = \text{Im } d_e$  and so  $[\mathcal{H}^k(M)] = H^k(M)$  (recall  $[\mathcal{H}^k(M)] \subset H^k(M)$ ).  $\square$

*Proof of Proposition 1.4:* Assume  $\langle Q_k v, v \rangle \geq 0$ . To show surjectivity of  $\mathcal{H}^k(M) \rightarrow H^k(M)$ , we need to prove that for all  $\omega_0 \in C^\infty(M, \Lambda^k(M))$  closed, there exists an exact form  $d\alpha$  (with  $\alpha \in C^\infty(M, \Lambda^k(M))$ ) such that  $G_k(\omega_0 + d\alpha) = 0$ . Consider  $\square := \delta_0 Q_k d + (d\delta_0)^{\frac{n}{2}-k+1}$  which is elliptic, self-adjoint and non-negative if  $Q_k \geq 0$ . Its kernel is finite dimensional (containing  $\ker(d + \delta_0)$ ) and all  $v \in \ker \square$  are smooth by elliptic regularity, and satisfy  $\langle \delta_0 Q_k dv, v \rangle_{L^2} = 0$ , which implies  $\langle Q_k dv, dv \rangle_{L^2} = 0$ . Let  $\mathbf{H} \subset L^2(\Lambda^k(M))$  be the  $L^2$  completion of the set  $C^\infty(M, \Lambda^k(M)) \cap \ker d$  of smooth closed forms and let us define  $\mathbf{Q}$  the symmetric form  $\mathbf{Q}(v, v) := \langle Q_k v, v \rangle_{L^2}$  on  $\mathbf{H}$ , it is a non-negative form induced by  $\Pi_{\mathbf{H}} Q_k$  on  $\mathbf{H}$  where  $\Pi_{\mathbf{H}}$  denotes orthogonal projection from  $L^2(\Lambda^k(M))$  to  $\mathbf{H}$ . The form has a domain  $D(\mathbf{Q})$  and Friedrichs extension theorem implies that there exists a self adjoint operator  $Q_k^{\text{Fr}} : \mathbf{H} \rightarrow \mathbf{H}$  with domain  $D(Q_k^{\text{Fr}})$  such that  $\langle Q_k^{\text{Fr}} u, u \rangle = \mathbf{Q}(u, u)$  for  $u \in D(\mathbf{Q}) \cap D(Q_k^{\text{Fr}})$ . But clearly  $d(C^\infty(M, \Lambda^{k-1}(M))) \subset D(Q_k^{\text{Fr}})$  and so  $\Pi_{\mathbf{H}} Q_k dv = Q_k^{\text{Fr}} dv$  for  $v$  smooth. Using now the spectral theorem for  $Q_k^{\text{Fr}}$ , we see that  $Q_k^{\text{Fr}} dv = 0$  with  $v$  smooth if and only if  $\langle Q_k dv, dv \rangle = 0$  and  $v$  is smooth, thus in particular if  $v \in \ker \square$ . Thus  $Q_k dv \perp \omega$  for all  $\omega \in \mathbf{H}$  if  $v \in \ker \square$ . Now this implies that, with  $\omega$  closed and smooth, we have  $\langle v, G_k \omega \rangle = \langle Q_k dv, \omega \rangle = 0$  for  $v \in \ker \square$  since  $Q_k$  is symmetric on closed forms, and so  $G_k \omega$  is in the range of  $\square$  and there exists  $\alpha$  such that  $\square \alpha = -G_k \omega$ , but since  $\text{Im } G_k \subset \text{Im } \delta_0$  which is orthogonal to  $\text{Im } d$ , we deduce that  $(d\delta_0)^{\frac{n}{2}-k+1} \alpha = 0$  and this achieves the proof. Note in particular that in this case  $\{d\varphi; L_{k-1}\varphi = 0\} = \{d\varphi; Q_k d\varphi \in \text{Im } \delta_0\}$ , see Corollaries 2.12 and 2.13 of [3] for discussions about these spaces.  $\square$

## 8. COMPUTATIONS IN SOME SPECIAL CASES

In this section we compute the operator  $L_k, G_k$  and  $Q_k$  in dimension 4 and 6.

**Proposition 8.1.** *Let  $(M^4, h)$  a four dimensional Riemannian manifold and define for any symmetric 2-tensor  $H$  the map  $j(H) := J(h^{-1}H)$  where  $J$  is defined in (2.4). Then we have*

$$\begin{aligned} L_1 &= \frac{1}{2} \delta d, & G_1 &= -\frac{1}{4} \delta \left( \Delta - 2j(\text{Ric}) + \frac{2}{3} \text{Scal} \right), & Q_1 &= \frac{1}{2} \left( \Delta - 2j(\text{Ric}) + \frac{2}{3} \text{Scal} \right), \\ L_0 &= -\frac{1}{16} \delta \left( \Delta - 2j(\text{Ric}) + \frac{2}{3} \text{Scal} \right) d, & G_0 &= 0, & Q_0 &= -\frac{1}{24} \left( \Delta \text{Scal} - 3|\text{Ric}|^2 + \text{Scal}^2 \right) \end{aligned}$$

where  $\text{Ric}$  is the Ricci tensor of  $h$  and  $\text{Scal}$  its scalar curvature

**Remark:** If  $n = 4$ ,  $L_{\frac{n}{2}-2}$  is the Paneitz operator (up to a constant factor). The result of Gursky and Viaclovsky [14] says that if the Yamabe invariant  $Y(M, [h_0])$  is positive and

$$\int_M Q d\text{vol}_{h_0} + \frac{1}{6} Y(M, [h_0])^2 > 0$$

then  $L_0$  is a non-negative operator with kernel reduced to constants. Combining with Theorem 2.6 of Branson-Gover[3], we have that  $\mathcal{H}^1(M) \simeq H^1(M)$  and there is a conformally invariant basis of  $H^1(M)$  with respect to  $[h_0]$  made of conformal harmonics.

**Corollary 8.2.** *Let  $M^4$  be a four dimensional manifold and  $\lambda_1(x) \geq \dots \geq \lambda_4(x)$  the eigenvalues of its Ricci curvature at  $x$ . If  $\lambda_2(x) + \lambda_3(x) + \lambda_4(x) \geq 0$  for all  $x \in M$  then  $\mathcal{H}^1(M) \rightarrow H^1(M)$  is surjective.*

*Proof:* Let  $D$  be the Levi-Civita connection of the metric  $h$  and  $\omega$  be a closed 1-form. By decomposing orthogonally the bilinear tensor  $D\omega$  in antisymmetric part  $\frac{d\omega}{2}$ , symmetric trace free part  $S_0$  and trace part  $\frac{-\delta\omega}{4}h$ , we get (recall that  $\omega$  is closed)

$$|D\omega|^2 = \frac{|d\omega|^2}{2} + |S_0|^2 + \frac{(\delta\omega)^2}{4} \geq \frac{(\delta\omega)^2}{4}.$$

Now, by the Bochner formula, we get

$$\langle \Delta\omega, \omega \rangle = \|D\omega\|_2^2 + \int_M \text{Ric}(\omega, \omega) \geq \|\delta\omega\|^2/4 + \int_M \text{Ric}(\omega, \omega) = \frac{\langle \Delta\omega, \omega \rangle}{4} + \int_M \text{Ric}(\omega, \omega),$$

and so  $\langle \Delta\omega, \omega \rangle \geq \frac{4}{3} \int_M \text{Ric}(\omega, \omega)$ . Therefore,

$$\begin{aligned} \langle Q_1\omega, \omega \rangle &= \frac{1}{2} \langle \Delta\omega, \omega \rangle + \int_M \left[ -\text{Ric}(\omega, \omega) + \frac{\text{Scal}}{3} |\omega|^2 \right] \geq \frac{1}{3} \int_M \text{Scal} |\omega|^2 - \text{Ric}(\omega, \omega), \\ &\geq \int_M (\lambda_1 + \lambda_2 + \lambda_3 + \lambda_4) |\omega|^2 - \lambda_1 |\omega|^2 \geq \int_M (\lambda_2 + \lambda_3 + \lambda_4) |\omega|^2 \geq 0, \end{aligned}$$

and we conclude by using Proposition 1.4.  $\square$

**Proposition 8.3.** *Let  $(M^6, h)$  be a six dimensional manifold. Let  $j$  be defined like in Lemma 8.1 and  $\text{tr}(H)$  denotes the trace of a symmetric tensor  $H$  with respect to  $h$ , then we have*

$$\begin{aligned} L_2 &= \frac{1}{2} \delta d, \quad G_2 = \frac{1}{4} \delta \left( \Delta - j(\text{Ric}) + \frac{2}{5} \text{Scal} \right), \quad Q_2 = \frac{1}{2} \left( \Delta - j(\text{Ric}) + \frac{2}{5} \text{Scal} \right), \\ L_1 &= -\frac{1}{16} \delta \left( \Delta - j(\text{Ric}) + \frac{2}{5} \text{Scal} \right) d, \\ G_1 &= \frac{(\delta d)^2 \delta}{64} + \frac{\delta d E_1 \delta}{32} - \frac{\delta E_1 \delta d}{16} - \frac{\delta E_1 d \delta}{16} - \frac{\delta d \delta E_1}{16} + \frac{\delta E_2}{16} + \frac{\delta E_1^2}{8}, \\ Q_1 &= -\frac{\Delta^2}{16} + \frac{d \delta E_1}{4} + \frac{\delta d E_1}{4} - \frac{d E_1 \delta}{8} + \frac{E_1 d \delta}{4} - \frac{2E_1^2 + E_2}{4}, \\ L_0 &= \frac{1}{384} (\delta d)^3 + \frac{\delta d E_1 \delta d}{192} - \frac{\delta E_1 d \delta d}{96} - \frac{\delta d \delta E_1 d}{96} + \frac{\delta E_2 d}{96} + \frac{\delta E_1^2 d}{48}, \\ Q_{01} &= -\frac{\Delta^2 \text{Scal}}{800} + \frac{\delta_0(E_1(d\text{Scal}))}{160} + \frac{\Delta |P|^2}{48} - \frac{\Delta \text{Scal}^2}{800} \\ &\quad - \frac{\text{Scal}^3}{8000} - \frac{\text{tr}(P^3)}{32} - \frac{h_0(\text{Ric}, B)}{32}, \end{aligned}$$

where  $E_2 := J(P^2 + 2B) - |P|^2/2$  and  $E_1 := J(P) - \text{tr}(P)/2$ ,  $B$  denotes the Bach tensor of  $h$ ,  $P$  the Schouten tensor,  $\text{Ric}$  the Ricci tensor and  $\text{Scal}$  the scalar curvature.

**Proposition 8.4.** *For any  $n \geq 4$ , we have the identities*

$$\begin{aligned} G_{\frac{n}{2}-1} &= (-1)^{\frac{n}{2}+1} \left( \frac{\delta d \delta}{4} - \frac{\delta j(P)}{2} + \delta \frac{\text{Tr}(P)\text{Id}}{4} \right) = (-1)^{\frac{n}{2}+1} \left( \frac{\delta d \delta}{4} - \delta \left( \frac{j(\text{Ric})}{n-2} - \frac{\text{Scal Id}}{2(n-1)} \right) \right), \\ Q_{\frac{n}{2}-1} &= \left( \frac{\Delta}{2} - \frac{2j(\text{Ric})}{n-2} + \frac{\text{Scal Id}}{n-1} \right), \\ L_{\frac{n}{2}-2} &= -\delta \left( \frac{d \delta}{16} - \frac{j(\text{Ric})}{4(n-2)} + \frac{\text{Scal Id}}{8(n-1)} \right) d, \\ G_{\frac{n}{2}-2} &= (-1)^{\frac{n}{2}+1} \left( \frac{(\delta d)^2 \delta}{64} + \frac{\delta d E_1 \delta}{32} - \frac{\delta E_1 \delta d}{16} - \frac{\delta E_1 d \delta}{16} - \frac{\delta d \delta E_1}{16} + \frac{\delta E_2}{16} + \frac{\delta E_1^2}{8} \right), \\ Q_{\frac{n}{2}-2} &= -\frac{\Delta^2}{16} + \frac{d \delta E_1}{4} + \frac{\delta d E_1}{4} - \frac{d E_1 \delta}{8} + \frac{E_1 d \delta}{4} - \frac{2E_1^2 + E_2}{4}, \\ L_{\frac{n}{2}-3} &= \frac{(\delta d)^3}{384} + \frac{\delta d E_1 \delta d}{192} - \frac{\delta E_1 d \delta d}{96} - \frac{\delta d \delta E_1 d}{96} + \frac{\delta E_2 d}{96} + \frac{\delta E_1^2 d}{48}, \end{aligned}$$

where  $E_2 := J(P^2 + \frac{4B}{n-4}) - |P|^2/2$  and  $E_1 := J(P) - \text{tr}(P)/2$ ,  $B$  denotes the Bach tensor of  $h$ ,  $P$  the Schouten tensor,  $\text{Ric}$  the Ricci tensor and  $\text{Scal}$  the scalar curvature.

For the non critical case, we have

**Proposition 8.5.** *We set  $j^\sharp(H) = 2j(H) - \text{tr}(H)\text{Id}$ . For any  $n \geq 3$ , we have*

$$L_k^1 = \frac{\delta d}{2} + \frac{(n-2k-2)d\delta}{2(n-2k+2)} + \frac{(n+k-2)(n-2k-2)}{8(n-1)(n-2)}\text{Scal} - \frac{(n-2k-2)j(\text{Ric})}{2(n-2)}$$

which generalizes the conformal Laplacian on functions,

$$L_k^2 = -\frac{n-2k-4}{16} \left( \frac{(d\delta)^2}{n-2k+4} + \frac{(\delta d)^2}{n-2k-4} + \frac{2dj^\sharp(P)\delta}{n-2k+4} - \frac{2\delta j^\sharp(P)d}{n-2k-4} \right) - \frac{j(P)\Delta + \Delta j^\sharp(P)}{2} + j^\sharp(P^2 + \frac{B}{n-4}) + \frac{(n-2k)j^\sharp(P)^2}{4}$$

which generalizes the Paneitz-Branson operator on functions.

*Proofs of Propositions 8.1, 8.3, 8.4 and 8.5:* This is a quite tedious computation, therefore we do not give the full details. By [6, Eq. (3.18)], we have

$$h_0^{-1}h_x = (I - x^2\frac{P}{2} + x^4\frac{h_2}{8} - x^6\frac{h_3}{48} + o(x^6)),$$

with  $P = \frac{1}{n-2}(2h_0^{-1}\text{Ric} - \frac{\text{Scal}}{n-1}I)$ ,  $h_2 = \frac{P^2}{2}$  for  $n = 4$  and  $h_2 = -\frac{2h_0^{-1}B}{n-4} + \frac{P^2}{2}$  and  $\text{tr}(h_3) = -\frac{4\text{tr}(P(h_0^{-1}B))}{n-4}$  for  $n = 6$  (where  $B$  is the Bach tensor); note that we have ignored the first log term in the metric expansion (i.e. the obstruction tensor) in dimension 4 and 6 since, as it is clear from Lemma 2.1, they do not show up in the construction the  $L_k^\ell, G_k, Q_k$ . We set  $B' = 0$  for  $n = 4$  and  $B' = \frac{2h_0^{-1}B}{n-4}$  for  $n = 6$ . Using the relations

$$\begin{aligned} L^{-1} &= I - A_1x^2 - (A_2 - A_1^2)x^4 - (A_3 + A_1^3 - A_1A_2 - A_2A_1)x^6 + o(x^6), \\ L^{-1}QL &= I + A_1'x^2 + (A_2' + [A_1', A_1])x^4 + (A_3' + [A_1', A_2] + [A_2', A_1] + A_1[A_1, A_1'])x^6 + o(x^6), \end{aligned}$$

for  $L = (I + A_1x^2 + A_2x^4 + A_3x^6 + o(x^6))$  and  $Q = (I + A_1'x^2 + A_2'x^4 + A_3'x^6 + o(x^6))$ , and the notations of the proof of Lemma 2.1, we get

$$\begin{aligned} h_x^{-1}h_0 &= I + x^2\frac{P}{2} + x^4\frac{3P^2 + 2B'}{16} + x^6\frac{h_3 + 3P^2 + 3B'P + 3PB'}{48} + o(x^6), \\ A_x &= -xP - x^3\frac{P^2 + 2B'}{4} - x^5\frac{2h_3 + P^3 + 4PB' + 2B'P}{16} + o(x^6), \\ O_x &= I + x^2\frac{P}{4} + x^4\frac{P^2 + B'}{16} + x^6\frac{2h_3 + 3P^3 + 5PB' + 4B'P}{192} + o(x^6), \\ I_x &= I + x^2\frac{J(P)}{4} + x^4\frac{J(P^2 + 2B') + J(P)^2}{32}, \\ &\quad + x^6\left(\frac{J(4h_3 + 2P^3 + 10PB' + 2B'P) + 3J(P)J(2B' + P^2) + J(P)^3}{384}\right) + o(x^6), \\ I_x^{-1} &= I - x^2\frac{J(P)}{4} + x^4\frac{-J(P^2 + 2B') + J(P)^2}{32}, \\ &\quad + x^6\left(\frac{-J(4h_3 + 2P^3 + 10PB' + 2B'P) + 3J(2B' + P^2)J(P) - J(P)^3}{384}\right) + o(x^6). \end{aligned}$$

Then we obtain

$$\begin{aligned} \star_x &= \star_0 - x^2\frac{[J(P), \star_0]}{4} + x^4\left(\frac{[J(P), [J(P), \star_0]] - [J(P^2 + 2B'), \star_0]}{32}\right) \\ &\quad + x^6\left(-\frac{[J(P)^3 + J(4h_3 + 2P^3 + 10PB' + 2B'P), \star_0]}{384}\right. \\ &\quad \left. + \frac{J(P)[J(P), \star_0]J(P) + [J(2B' + P^2), [J(P), \star_0]]}{128}\right) + o(x^6) \end{aligned}$$

using the relations  $\star_0 J(H) + J(H)\star_0 = \text{tr}(H)\star_0$  and  $[J(H), J(H')] = -J([H, H'])$ , we get

$$\star_x^{-1}[\partial_x, \star_x] = xE_1 + x^3 \frac{E_2}{4} + x^5 \frac{C'}{32}$$

with

$$C'1 = -2\text{tr}(PB') - \text{tr}(P^3) - \frac{3}{2}\text{tr}(P)\text{tr}(P^2), \quad E_1 = J(P) - \frac{\text{tr}(P)}{2}, \quad E_2 := J(P^2 + 2B') - \frac{\text{tr}(P^2)}{2},$$

$$\delta_x = \delta_0 + x^2 \frac{[\delta_0, E_1]}{2} + x^4 \frac{D}{16}, \quad \text{where } D := [\delta_0, E_2] + 2[[\delta_0, E_1], E_1].$$

Therefore,

$$\begin{aligned} \Delta_k = & \begin{pmatrix} -(x\partial_x)^2 + (n-2k)x\partial_x & 2(-1)^{k+1}d \\ 0 & -(x\partial_x)^2 + (n-2k+2)x\partial_x \end{pmatrix} \\ & + x^2 \begin{pmatrix} \Delta_0 - E_1 x \partial_x & (-1)^k [d, E_1] \\ 2(-1)^{k+1} \delta_0 & \Delta_0 - E_1 (2 + x\partial_x) \end{pmatrix} \\ & + x^4 \begin{pmatrix} A_1 - A_2 x \partial_x & (-1)^k [d, A_2] \\ 2(-1)^{k+1} [\delta_0, E_1] & A_1 - A_2 (4 + x\partial_x) \end{pmatrix} \\ & + x^6 \begin{pmatrix} A_3 & A_4 \\ A_5 & A_6 \end{pmatrix} + o(x^6) \end{aligned}$$

where  $A_1 := \frac{d[\delta_0, E_1] + [\delta_0, E_1]d}{2}$ ,  $A_2 := \frac{E_2}{4}$  (since the Bach tensor is trace-free),  $A_5 := -(-1)^k \frac{3D}{8}$  and, when  $k = 1$ ,  $A_6 1 := \frac{9}{160}\text{Scal}|P|^2 + \frac{3}{16}\text{tr}(P^3) + \frac{3}{16}h_0(\text{Ric}, B)$ .

For  $n = 6$  and  $k = 1$  we follow the formal method of Subsection 4.2.3. We first have

$$\omega'_{F_1} = \frac{dx}{x} - x^2 \left( \frac{d\text{Scal}}{80} + \frac{\text{Scal}dx}{40x} \right) + x^4 \left( \frac{\Delta\text{Scal}}{160} + \frac{\text{Scal}^2}{800} - \frac{|P|^2}{16} \right) \frac{dx}{x}$$

and so, by computing  $d\omega'_{F_1}$  and  $\Delta_k \omega'_{F_1}$ , one finds

$$\begin{aligned} D'_0 1 &= -\frac{d\Delta\text{Scal}}{160} - \frac{d\text{Scal}^2}{800} + \frac{d|P|^2}{16} = -\frac{d\Delta\text{Scal}}{160} - \frac{11d\text{Scal}^2}{3200} + \frac{d|\text{Ric}|^2}{64} \\ -B'_0 1 &= A_6 1 + \frac{[E_1, \delta_0]d\text{Scal}}{40} - \frac{(A_1 - 6A_2)\text{Scal}}{40} + (\Delta_0 - 6E_1) \left( \frac{\Delta\text{Scal}}{160} + \frac{\text{Scal}^2}{800} - \frac{|P|^2}{16} \right) \\ &= \frac{3}{16}\text{tr}(P^3) + \frac{3}{16}h_0(\text{Ric}, B) - \frac{3\delta_0(E_1(d\text{Scal}))}{80} + \frac{\Delta^2\text{Scal}}{160} + \frac{\Delta\text{Scal}^2}{800} - \frac{\Delta|P|^2}{16} + \frac{3\text{Scal}^3}{4000} \end{aligned}$$

Hence we get  $Q_0 1$  in the case  $n = 6$  by Corollary 4.9. The other computations are made similarly. For instance, for  $k = n/2 - 1$ ,

$$\Delta\omega_{F_1} = x^2 \delta_0 d\omega_0 + x^3 (-1)^{\frac{n}{2}+1} \left( \frac{\delta_0 d\delta_0 \omega_0}{2} - 2\delta_0 E_1 \omega_0 \right) \wedge dx + O(x^4),$$

and so

$$B_{\frac{n}{2}-1} \omega_0 = -\frac{\delta_0 d\delta_0 \omega_0}{2} + 2\delta_0 E_1 \omega_0.$$

We have  $\delta\omega_{F_1} = \frac{x^4}{2}\delta_0 E_1 \omega_0 + O(x^5)$ , and so  $C_{\frac{n}{2}-1} = \frac{\delta_0 E_1}{2}$ . By Proposition 4.5,

$$G_{\frac{n}{2}-1} = (-1)^{\frac{n}{2}+1} \left( \frac{\delta_0 d\delta_0}{4} - \frac{\delta_0 E_1}{2} \right),$$

which implies the expression for  $L_{\frac{n}{2}-2}$  by (4.11). For  $k = \frac{n}{2} - 2$ , we have

$$\begin{aligned} \omega_{F_1} &= \omega_0 + x^2 \left( -\frac{d\delta\omega_0}{8} - \frac{\delta d\omega_0}{4} + \frac{(-1)^{\frac{n}{2}} \delta\omega_0}{4} \wedge \frac{dx}{x} \right) \\ &\quad + (-1)^{\frac{n}{2}} x^4 \left( \frac{\delta E_1 \omega_0}{4} - \frac{\delta d\delta\omega_0}{16} - \frac{E_1 \delta\omega_0}{8} \right) \wedge \frac{dx}{x} \\ \delta &= \begin{pmatrix} 0 & (-1)^{\frac{n}{2}} (x\partial_x - 6) \\ 0 & 0 \end{pmatrix} + x^2 \begin{pmatrix} \delta_0 & (-1)^{\frac{n}{2}} E_1 \\ 0 & \delta_0 \end{pmatrix} + x^4 \begin{pmatrix} \frac{[\delta_0, E_1]}{2} & (-1)^{\frac{n}{2}} A_2 \\ 0 & \frac{[\delta_0, E_1]}{2} \end{pmatrix} \\ &\quad + x^6 \begin{pmatrix} \frac{2[\delta_0, A_2] + [\delta_0, E_1], E_1}{8} & * \\ 0 & * \end{pmatrix} \end{aligned}$$

and so

$$\begin{aligned} C_{\frac{n}{2}-2} &= \frac{\delta_0 A_2}{4} + \frac{\delta_0 E_1^2}{8} - \frac{\delta_0 E_1 d\delta_0}{16} - \frac{\delta_0 E_1 \delta_0 d}{8}, \\ B_{\frac{n}{2}-2} &= \frac{3\delta_0 A_2}{2} - \frac{\delta E_1 \delta d}{2} - \frac{3\delta E_1 d\delta}{8} - \frac{\delta d\delta E_1}{4} + \frac{(\delta d)^2 \delta}{16} + \frac{\delta d E_1 \delta}{8} + \frac{3\delta E_1^2}{4}, \\ G_{\frac{n}{2}-2} &= (-1)^{\frac{n}{2}+1} \left( \frac{(\delta d)^2 \delta}{64} + \frac{\delta d E_1 \delta}{32} - \frac{\delta E_1 \delta d}{16} - \frac{\delta E_1 d\delta}{16} - \frac{\delta d\delta E_1}{16} + \frac{\delta E_2}{16} + \frac{\delta E_1^2}{8} \right). \end{aligned}$$

In the case  $k = \frac{n}{2} - 1$ , we have  $\omega'_{F_1} = \omega_0 \wedge \frac{dx}{x} + x^2 \left( -\frac{d\delta}{4} + \frac{E_1}{2} \right) \wedge \frac{dx}{x}$ ,  $D'_{\frac{n}{2}-2} = (-1)^{\frac{n}{2}+1} \frac{dE_1}{2}$  and  $B'_{\frac{n}{2}-2} = (-1)^{\frac{n}{2}} d\delta E_1 - \frac{dE_1 \delta}{2} - E_2 - \frac{(d\delta)^2}{4} + \delta dE_1 + E_1 d\delta - 2E_1^2 + \frac{\delta dE_1}{2}$ . So we get

$$Q_{\frac{n}{2}-2} = -\frac{\Delta^2}{16} + \frac{d\delta E_1}{4} + \frac{\delta dE_1}{4} - \frac{dE_1 \delta}{8} + \frac{E_1 d\delta}{4} - \frac{2E_1^2 + E_2}{4}.$$

□

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